Electronic Journal of Differential Equations, Vol. 2018 (2018), No. 22, pp. 1–12. ISSN: 1072-6691. URL: http://ejde.math.txstate.edu or http://ejde.math.unt.edu

FRACTIONAL ELLIPTIC PROBLEMS WITH TWO CRITICAL SOBOLEV-HARDY EXPONENTS

WENJING CHEN

Communicated by Giovanni Molica Bisci

ABSTRACT. By using the mountain pass lemma and a concentration compactness principle, we obtain the existence of positive solutions to the fractional elliptic problem with two critical Hardy-Sobolev exponents at the origin.

1. INTRODUCTION

In this article, we study the following doubly critical problem involving the fractional Laplacian

$$(-\Delta)^{s}u - \gamma \frac{u}{|x|^{2s}} = \frac{|u|^{2^{s}_{s}(\alpha)-2}u}{|x|^{\alpha}} + \frac{|u|^{2^{s}_{s}(\beta)-2}u}{|x|^{\beta}}, \quad u > 0, \quad \text{in } \mathbb{R}^{n}, \tag{1.1}$$

where $s \in (0, 1), 0 < \alpha, \beta < 2s < n$ with $\alpha \neq \beta, \gamma < \gamma_H$ with

$$\gamma_H = 4^s \frac{\Gamma^2(\frac{n+2s}{4})}{\Gamma^2(\frac{n-2s}{4})}$$

being the fractional best Hardy constant on \mathbb{R}^n , and $2^*_s(\alpha) = 2(n-\alpha)/(n-2s)$ is the fractional critical Hardy-Sobolev exponent. The operator $(-\Delta)^s$ is the fractional Laplacian defined as

$$(-\Delta)^{s} u(x) = c_{n,s} \operatorname{pv} \int_{\mathbb{R}^{n}} \frac{u(x) - u(y)}{|x - y|^{n + 2s}} dy, \quad s \in (0, 1),$$

where pv stands for the Cauchy principle value and

$$c_{n,s} = 2^{2s-1} \pi^{-\frac{n}{2}} \frac{\Gamma\left(\frac{n+2s}{2}\right)}{|\Gamma(-s)|}$$

is the normalization constant so that the identity

$$(-\Delta)^{s} u = \mathcal{F}^{-1}(|\xi|^{2s}(\mathcal{F}u)) \quad \forall \xi \in \mathbb{R}^{n}, \ s \in (0,1), \ u \in \mathcal{S}(\mathbb{R}^{n}),$$

holds, here $\mathcal{F}u$ denotes the Fourier transform of u, $\mathcal{F}u(\xi) = \int_{\mathbb{R}^n} e^{-2\pi i x \cdot \xi} u(x) dx$, and $\mathcal{S}(\mathbb{R}^n)$ the Schwartz class, see [14] and references therein for the basics on the fractional Laplacian.

²⁰¹⁰ Mathematics Subject Classification. 35J20, 35J60, 47G20.

Key words and phrases. Fractional elliptic problems; mountain pass lemma;

critical fractional Hardy-Sobolev exponent; concentration compactness principle. ©2018 Texas State University.

Submitted July 1, 2017. Published January 15, 2018.

W. CHEN

In previous twenty years, the nonlocal elliptic problems have been investigated by many researchers, for example, [18, 27, 29, 30, 31] for the subcritical case, [3, 8, 23, 19, 28, 32, 33] for the critical case, [9, 10, 11] for the existence of solutions to fractional Laplacian system. Moreover, a great attention has been devoted to study the existence of solutions for the nonlocal problems with Hardy potential or nonlinearity term, we refer to see [1, 2, 4, 13, 15, 16, 34, 35, 36] and the references therein. In particular, the existence of solutions to the problem

$$(-\Delta)^{s}u - \gamma \frac{u}{|x|^{2s}} = \frac{u^{2^{*}_{s}(\alpha)-1}}{|x|^{\alpha}}, \quad u > 0 \quad \text{in } \mathbb{R}^{n},$$
 (1.2)

corresponds to the minimization problem

$$\mu_{s,\gamma,\alpha}(\mathbb{R}^{n}) = \inf_{u \in H^{s}(\mathbb{R}^{n}) \setminus \{0\}} \frac{\int_{\mathbb{R}^{n}} |(-\Delta)^{s/2} u|^{2} dx - \gamma \int_{\mathbb{R}^{n}} \frac{|u|^{2}}{|x|^{2s}} dx}{\left(\int_{\mathbb{R}^{n}} \frac{|u|^{2s}(\alpha)}{|x|^{\alpha}} dx\right)^{\frac{2}{2s}(\alpha)}}.$$
 (1.3)

Fall et al. [16] proved the existence of extremals for $\mu_{s,0,\alpha}(\mathbb{R}^n)$ in the case $s = \frac{1}{2}$. Yang [35] proved that there exists a positive, radially symmetric and non-increasing extremal for $\mu_{s,0,\alpha}(\mathbb{R}^n)$ when $s \in (0,1)$. Asymptotic properties of the positive solutions was given by Lei [24] and Yang-Yu [37]. The existence of extremals for $\mu_{s,\gamma,\alpha}(\mathbb{R}^n)$ in (1.3), when $\alpha \in [0,2s)$ and $\gamma \in (-\infty,\gamma_H)$, was recently studied by Ghoussoub and Shakerian in [21]. Moreover, the authors in [21] used the mountain pass lemma to establish the existence of a nontrivial weak solution to the problem

$$(-\Delta)^{s}u - \gamma \frac{u}{|x|^{2s}} = |u|^{2^{*}_{s}-2}u + \frac{|u|^{2^{*}_{s}(\alpha)-2}u}{|x|^{\alpha}}, \quad u > 0, \quad \text{in } \mathbb{R}^{n}.$$

Furthermore, the authors in [36] showed the existence of nontrivial solutions for fractional elliptic problem in \mathbb{R}^n with the critical nonlocal Hartree term and critical fractional Hardy-Sobolev term.

It is worth pointing out that in the local case, i.e. s = 1, the existence and multiplicity of solutions for the Laplacian problems with Hardy terms have been extensively studied, we refer the reader to [5, 7, 12, 17, 22] and references therein.

The aim of this paper is to consider the existence of nontrivial weak solutions of (1.1), which has a single pole with different powers of singularity and fractional critical Hardy-Sobolev exponents. We get the existence of nontrivial weak solutions of our problem by the Mountain Pass Lemma with concentration-compactness principle. Our result can be stated as follows.

Theorem 1.1. Let 0 < s < 1, $0 < \alpha, \beta < 2s < n$ with $\alpha \neq \beta$, and $\gamma < \gamma_H$. Then problem (1.1) admits a nontrivial solution.

This article is organized as follows: in Section 2, we give some preliminaries about fractional Laplacian harmonic extension and function space, and also the fractional Hardy-Sobolev inequality. We prove the compactness of the energy in Section 3. Section 4 is concerned with the proof of our main result.

2. Preliminary results

In this section, we first introduce suitable function spaces for the variational principles that will be needed in the sequel. Caffarelli and Silvestre in [6] showed that the fractional Laplacian operator can be realized in a local way by using one more variable and the so-called s-harmonic extension, that is, for a function

 $u \in H^{s}(\mathbb{R}^{n})$, we say that $U = E_{s}(u)$ is its s-harmonic extension to the upper half-space, \mathbb{R}^{n+1}_{+} , i.e. it is a solution to the problem

$$\operatorname{div}(y^{1-2s}\nabla U) = 0 \quad \text{in } \mathbb{R}^{n+1}_+, U = u \quad \text{on } \mathbb{R}^n \times \{y = 0\}.$$

Define the space $X^s(\mathbb{R}^{n+1}_+)$ as the closure of $C_0^{\infty}(\overline{\mathbb{R}^{n+1}_+})$ with the norm

$$||U||_{X^{s}(\mathbb{R}^{n+1}_{+})} := \left(k_{s} \int_{\mathbb{R}^{n+1}_{+}} y^{1-2s} |\nabla U(x,y)|^{2} \, dx \, dy\right)^{1/2},$$

where $k_s = \frac{\Gamma(s)}{2^{1-2s}\Gamma(1-s)}$ is a normalization constant chosen in such a way that the extension operator $U : H^s(\mathbb{R}^n) \to X^s(\mathbb{R}^{n+1}_+)$ is an isometry, that is, for any $u \in H^s(\mathbb{R}^n)$, we have

$$||U||_{X^{s}(\mathbb{R}^{n+1}_{+})} = ||u||_{H^{s}(\mathbb{R}^{n})} = ||(-\Delta)^{s/2}u||_{L^{2}(\mathbb{R}^{n})}.$$
(2.1)

Conversely, for a function $U \in X^s(\mathbb{R}^{n+1}_+)$, we denote its trace on $\mathbb{R}^n \times \{y = 0\}$ as $u = \operatorname{Tr}(U) := U(\cdot, 0)$. This trace operator is also well defined and satisfies

$$||u||_{H^{s}(\mathbb{R}^{n})} = ||U(\cdot, 0)||_{H^{s}(\mathbb{R}^{n})} \le ||U||_{X^{s}(\mathbb{R}^{n+1}_{+})}.$$
(2.2)

Caffarelli and Silvestre [6] showed that the extension function $U := E_s(u)$ is related to the fractional Laplacian of the original function u in the following way:

$$(-\Delta)^{s}u(x) = \frac{\partial U}{\partial \nu^{s}} := -k_{s} \lim_{y \to 0^{+}} y^{1-2s} \frac{\partial U}{\partial y}(x,y).$$

Thus, problem (1.1) can be written as the local problem

$$-\operatorname{div}(y^{1-2s}\nabla U) = 0 \quad \text{in } \mathbb{R}^{n+1}_+$$
$$\frac{\partial U}{\partial \nu^s} = \gamma \frac{u}{|x|^{2s}} + \frac{|u|^{2^*_s(\alpha)-2}u}{|x|^{\alpha}} + \frac{|u|^{2^*_s(\beta)-2}u}{|x|^{\beta}} \quad \text{on } \mathbb{R}^n,$$
(2.3)

where and in the follows $u = U(\cdot, 0)$. A function $U \in X^{s}(\mathbb{R}^{n+1}_{+})$ is said to be a weak solution to (2.3), if for all $\Psi \in X^{s}(\mathbb{R}^{n+1}_{+})$,

$$k_s \int_{\mathbb{R}^{n+1}_+} y^{1-2s} \langle \nabla U, \nabla \Psi \rangle \, dx \, dy = \int_{\mathbb{R}^n} \gamma \frac{u}{|x|^{2s}} \psi \, dx + \int_{\mathbb{R}^n} \frac{|u|^{2^*_s(\alpha)-2}u}{|x|^{\alpha}} \psi \, dx$$
$$+ \int_{\mathbb{R}^n} \frac{|u|^{2^*_s(\beta)-2}u}{|x|^{\beta}} \psi \, dx,$$

where $\psi = \Psi(\cdot, 0)$. The energy functional corresponding to (2.3) is

$$J(U) = \frac{1}{2} \|U\|_{X^s(\mathbb{R}^{n+1}_+)}^2 - \frac{\gamma}{2} \int_{\mathbb{R}^n} \frac{|u|^2}{|x|^{2s}} dx - \frac{1}{2^*_s(\alpha)} \int_{\mathbb{R}^n} \frac{|u|^{2^*_s(\alpha)}}{|x|^s} dx - \frac{1}{2^*_s(\beta)} \int_{\mathbb{R}^n} \frac{|u|^{2^*_s(\beta)}}{|x|^s} dx.$$

We note that for any weak solution $U \in X^s(\mathbb{R}^{n+1}_+)$ to (2.3), the function $u = U(\cdot, 0)$ is in $H^s(\mathbb{R}^n)$ and is a weak solution to problem (1.1). Hence the associated trace of any critical point U of J in $X^s(\mathbb{R}^{n+1}_+)$ is a weak solution for (1.1). Let us recall the following results.

Lemma 2.1. Assume that 0 < s < 1.

(i) (The fractional Hardy inequality [20]) For all $u \in H^{s}(\mathbb{R}^{n})$, we have

$$\gamma_H \int_{\mathbb{R}^n} \frac{|u|^2}{|x|^{2s}} dx \le \int_{\mathbb{R}^n} |(-\Delta)^{s/2} u|^2 dx,$$
 (2.4)

where $\gamma_H = 4^s \frac{\Gamma^2(\frac{n+2s}{4})}{\Gamma^2(\frac{n-2s}{4})}$ is the best constant in the above inequality on \mathbb{R}^n .

(ii) (The fractional Hardy-Sobolev inequality [21]) Assume $0 \leq \alpha \leq 2s < n$. Then, there exist positive constants c and C, such that for all $u \in H^s(\mathbb{R}^n)$,

$$\left(\int_{\mathbb{R}^n} \frac{|u|^{2^*_s(\alpha)}}{|x|^{\alpha}} dx\right)^{\frac{2^*}{2^*_s(\alpha)}} \le c \int_{\mathbb{R}^n} |(-\Delta)^{s/2} u|^2 dx.$$
(2.5)

Moreover, if $\gamma < \gamma_H$, then

$$C\Big(\int_{\mathbb{R}^n} \frac{|u|^{2^*_s(\alpha)}}{|x|^{\alpha}} \, dx\Big)^{\frac{2}{2^*_s(\alpha)}} \le \int_{\mathbb{R}^n} |(-\Delta)^{s/2} u|^2 \, dx - \gamma \int_{\mathbb{R}^n} \frac{|u|^2}{|x|^{2s}} \, dx, \qquad (2.6)$$

for all $u \in H^s(\mathbb{R}^n)$.

Remark 2.2. One can use (2.1) to rewrite inequalities (2.4), (2.5) and (2.6) as the following trace class inequalities:

$$\gamma_H \int_{\mathbb{R}^n} \frac{|u|^2}{|x|^{2s}} \, dx \le \|U\|_{X^s(\mathbb{R}^{n+1}_+)}^2, \tag{2.7}$$

$$\left(\int_{\mathbb{R}^n} \frac{|u|^{2^*_s(\alpha)}}{|x|^{\alpha}} \, dx\right)^{\frac{2}{2^*_s(\alpha)}} \le c \|U\|^2_{X^s(\mathbb{R}^{n+1}_+)},\tag{2.8}$$

$$C\Big(\int_{\mathbb{R}^n} \frac{|u|^{2^*_s(\alpha)}}{|x|^{\alpha}} \, dx\Big)^{\frac{2}{2^*_s(\alpha)}} \le \|U\|^2_{X^s(\mathbb{R}^{n+1}_+)} - \gamma \int_{\mathbb{R}^n} \frac{|u|^2}{|x|^{2s}} \, dx.$$
(2.9)

In what follows, we will denote by $X^s(\mathbb{R}^{n+1}_+)$ the closure of $C_0^\infty(\overline{\mathbb{R}^{n+1}_+})$ for the following norm

$$||U|| := \left(k_s \int_{\mathbb{R}^{n+1}_+} y^{1-2s} |\nabla U|^2 \, dx \, dy - \gamma \int_{\mathbb{R}^n} \frac{|u|^2}{|x|^{2s}} \, dx\right)^{1/2} \quad \text{for all } \gamma < \gamma_H. \quad (2.10)$$

Note that inequality (2.7) asserts that $X^s(\mathbb{R}^{n+1}_+)$ is embedded in the weighted space $L^2(\mathbb{R}^n, |x|^{-2s})$ and this embedding is continuous. Set $\gamma_+ = \max\{\gamma, 0\}$ and $\gamma_- = -\max\{\gamma, 0\}$. The following inequalities hold for any $u \in X^s(\mathbb{R}^{n+1}_+)$,

$$(1 - \frac{\gamma_{+}}{\gamma_{H}}) \|U\|_{X^{s}(\mathbb{R}^{n+1}_{+})}^{2} \leq \|U\|^{2} \leq (1 + \frac{\gamma_{-}}{\gamma_{H}}) \|U\|_{X^{s}(\mathbb{R}^{n+1}_{+})}^{2}.$$
 (2.11)

Thus, $\|\cdot\|$ is equivalent to the norm $\|\cdot\|_{X^s(\mathbb{R}^{n+1}_+)}$.

The best constant $\mu_{s,\gamma,\alpha}(\mathbb{R}^n)$ in inequality (2.6) can be written as

$$S(n, s, \gamma, \alpha) = \inf_{\substack{U \in X^{s}(\mathbb{R}^{n+1}_{+}) \setminus \{0\}}} I_{\gamma, \alpha}(U), \quad \text{with}$$
$$I_{\gamma, \alpha}(U) = \frac{k_{s} \int_{\mathbb{R}^{n+1}_{+}} y^{1-2s} |\nabla U|^{2} \, dx \, dy - \gamma \int_{\mathbb{R}^{n}} \frac{|u|^{2}}{|x|^{2s}} \, dx}{\left(\int_{\mathbb{R}^{n}} \frac{|u|^{2^{s}_{s}(\alpha)}}{|x|^{\alpha}} \, dx\right)^{\frac{2}{2^{s}_{s}(\alpha)}}}$$

If $S(n, s, \gamma, \alpha)$ is attained at some function $U \in X^s(\mathbb{R}^{n+1}_+)$, then u = U(., 0) will be a function in $H^s(\mathbb{R}^n)$, where $\mu_{s,\gamma,\alpha}(\mathbb{R}^n)$ is attained. Recently, Ghoussoub and Shakerian [21] proved the extremal function of $S(n, s, \gamma, \alpha)$ is attained as following.

Lemma 2.3 ([21]). Suppose 0 < s < 1, $0 \le \alpha < 2s < n$, and $\gamma < \gamma_H$. Then

- (1) If $\{\alpha > 0\}$ or $\alpha = 0$ and $\gamma \ge 0$, then $S(n, s, \gamma, \alpha)$ is attained in $X^s(\mathbb{R}^{n+1}_+)$ by $W_{\gamma,\alpha}$.
- (2) If $\alpha = 0$ and $\gamma < 0$, then there are no extremals for $S(n, s, \gamma, \alpha)$ in $X^s(\mathbb{R}^{n+1}_+)$.

3. Compactness Lemmas

In this section, we study the compactness properties of the functional

$$J(U) = \frac{1}{2} \|U\|^2 - \frac{1}{2_s^*(\alpha)} \int_{\mathbb{R}^n} \frac{|u|^{2_s^*(\alpha)}}{|x|^{\alpha}} \, dx - \frac{1}{2_s^*(\beta)} \int_{\mathbb{R}^n} \frac{|u|^{2_s^*(\beta)}}{|x|^{\beta}} \, dx \tag{3.1}$$

for $U \in X^s(\mathbb{R}^{n+1}_+)$, where again $u := U(\cdot, 0)$. From Lemma 2.1, we have that $J \in C^1(X^s(\mathbb{R}^{n+1}_+))$.

Definition 3.1. Let $c \in \mathbb{R}$, E be a Banach space and $J \in C^1(E, \mathbb{R})$.

(i) $\{u_k\}$ is a $(PS)_c$ sequence in E for J if $J(u_k) = c + o(1)$ and $J'(u_k) = o(1)$ strongly in E^* as $k \to \infty$.

(ii) We say that J satisfies the $(PS)_c$ condition if any $(PS)_c$ sequence $\{u_k\}$ for J in E has a convergent subsequence.

Proposition 3.2. Suppose $0 < \alpha, \beta < 2s$ and $\gamma < \gamma_H$, then the functional J defined in (3.1) satisfies the Palais-Smale condition $(PS)_c$ for $c < c_*$, where

$$c_* := \min\left\{\frac{2s-\alpha}{2(n-\alpha)}S(n,s,\gamma,\alpha)^{\frac{n-\alpha}{2s-\alpha}}, \frac{2s-\beta}{2(n-\beta)}S(n,s,\gamma,\beta)^{\frac{n-\beta}{2s-\beta}}\right\}.$$
(3.2)

Proof. Let $\{U_k\}_{k\in\mathbb{N}}$ be the Palais-Smale sequence of the functional J, i.e.

$$J(U_k) \to c$$
, $J'(U_k) \to 0$ in $(X^s(\mathbb{R}^{n+1}_+))'$ as $k \to \infty$.

Then

$$J(U_k) = \frac{1}{2} \|U_k\|^2 - \frac{1}{2_s^*(\alpha)} \int_{\mathbb{R}^n} \frac{|u_k|^{2_s^*(\alpha)}}{|x|^{\alpha}} dx - \frac{1}{2_s^*(\beta)} \int_{\mathbb{R}^n} \frac{|u_k|^{2_s^*(\beta)}}{|x|^{\beta}} dx$$

= $c + o_k(1),$ (3.3)

and

c

$$\langle J'(U_k), U_k \rangle = \|U_k\|^2 - \int_{\mathbb{R}^n} \frac{|u_k|^{2^*_s(\alpha)}}{|x|^{\alpha}} \, dx - \int_{\mathbb{R}^n} \frac{|u_k|^{2^*_s(\beta)}}{|x|^{\beta}} \, dx = o_k(1) \|w_k\|, \quad (3.4)$$

where again $u_k = U_k(\cdot, 0)$ and $o_k(1) \to 0$ as $k \to \infty$. From (3.3) and (3.4), we have

$$+ o_{k}(1) ||U_{k}|| = J(U_{k}) - \frac{1}{2} \langle J'(U_{k}), U_{k} \rangle$$

= $\left(\frac{1}{2} - \frac{1}{2_{s}^{*}(\alpha)}\right) \int_{\mathbb{R}^{n}} \frac{|u_{k}|^{2_{s}^{*}(\alpha)}}{|x|^{\alpha}} dx + \left(\frac{1}{2} - \frac{1}{2_{s}^{*}(\beta)}\right) \int_{\mathbb{R}^{n}} \frac{|u_{k}|^{2_{s}^{*}(\beta)}}{|x|^{\beta}} dx.$

Since $2_{s}^{*}(\alpha) > 2, 2_{s}^{*}(\beta) > 2$, we have

$$\int_{\mathbb{R}^n} \frac{|u_k|^{2^*_s(\alpha)}}{|x|^{\alpha}} dx \le C + o_k(1) \|U_k\|, \qquad \int_{\mathbb{R}^n} \frac{|u_k|^{2^*_s(\beta)}}{|x|^{\beta}} dx \le C + o_k(1) \|U_k\|.$$
(3.5)

By (3.4) and (3.5), we obtain

$$||U_k||^2 + o_k(1)||U_k|| = \int_{\mathbb{R}^n} \frac{|u_k|^{2^*_s(\alpha)}}{|x|^{\alpha}} \, dx + \int_{\mathbb{R}^n} \frac{|u_k|^{2^*_s(\beta)}}{|x|^{\beta}} \, dx \le C + o_k(1)||U_k||, \quad (3.6)$$

which implies that $\{U_k\}_{k\in\mathbb{N}}$ is bounded in $X^s(\mathbb{R}^{n+1}_+)$. It follows that there exists a subsequence, still denote by U_k , such that $U_k \rightharpoonup U$ in $X^s(\mathbb{R}^{n+1}_+)$. For any $\Psi \in$ $C_0^{\infty}(\mathbb{R}^{n+1}_+)$, we have

$$o_{k}(1) = \langle J'(U_{k}), \Psi \rangle$$

$$= k_{s} \int_{\mathbb{R}^{n+1}_{+}} y^{1-2s} \langle \nabla U_{k}, \nabla \Psi \rangle \, dx \, dy - \gamma \int_{\mathbb{R}^{n}} \frac{u_{k}(x)}{|x|^{2s}} \psi(x) \, dx \qquad (3.7)$$

$$- \int_{\mathbb{R}^{n}} \frac{|u_{k}(x)|^{2^{*}_{s}(\alpha)-2} u_{k}(x)}{|x|^{\alpha}} \psi(x) \, dx - \int_{\mathbb{R}^{n}} \frac{|u_{k}(x)|^{2^{*}_{s}(\beta)-2} u_{k}(x)}{|x|^{\beta}} \psi(x) \, dx.$$

Since $U_k \rightharpoonup U$ in $X^s(\mathbb{R}^{n+1}_+)$ as $k \rightarrow \infty$, we have that

$$\int_{\mathbb{R}^{n+1}_+} y^{1-2s} \langle \nabla U_k, \nabla \Psi \rangle \, dx \, dy - \gamma \int_{\mathbb{R}^n} \frac{u_k(x)}{|x|^{2s}} \psi(x) \, dx$$
$$\to \int_{\mathbb{R}^{n+1}_+} y^{1-2s} \langle \nabla U, \nabla \Psi \rangle \, dx \, dy - \gamma \int_{\mathbb{R}^n} \frac{u(x)}{|x|^{2s}} \psi(x) \, dx,$$

for all $\Psi \in C_0^{\infty}(\mathbb{R}^{n+1}_+)$, where $u = U(\cdot, 0)$. Moreover, the boundedness of U_k in $X^s(\mathbb{R}^{n+1}_+)$ implies that $|u_k|^{2^*_s(\alpha)-2}u_k$ and $|u_k|^{2^*_s(\beta)-2}u_k$ are bounded in $L^{\frac{2^*_s(\alpha)}{2^*_s(\alpha)-1}}(\mathbb{R}^n, |x|^{-\alpha})$ and $L^{\frac{2^*_s(\beta)}{2^*_s(\beta)-1}}(\mathbb{R}^n, |x|^{-\beta})$ respectively. Therefore,

$$|u_k|^{2^*_s(\alpha)-2}u_k \rightharpoonup |u|^{2^*_s(\alpha)-2}u \quad \text{in } L^{\frac{2^*_s(\alpha)}{2^*_s(\alpha)-1}}(\mathbb{R}^n, |x|^{-\alpha}), |u_k|^{2^*_s(\beta)-2}u_k \rightharpoonup |u|^{2^*_s(\beta)-2}u \quad \text{in } L^{\frac{2^*_s(\beta)}{2^*_s(\beta)-1}}(\mathbb{R}^n, |x|^{-\beta}).$$

Thus, taking limits as $k \to \infty$ in (3.7), we obtain

$$0 = \langle J'(U), \Psi \rangle$$

= $k_s \int_{\mathbb{R}^{n+1}_+} y^{1-2s} \langle \nabla U, \nabla \Psi \rangle \, dx \, dy - \gamma \int_{\mathbb{R}^n} \frac{u(x)}{|x|^{2s}} \psi(x) \, dx$
$$- \int_{\mathbb{R}^n} \frac{|u(x)|^{2^*_s(\alpha)-2}u(x)}{|x|^{\alpha}} \psi(x) \, dx - \int_{\mathbb{R}^n} \frac{|u(x)|^{2^*_s(\beta)-2}u(x)}{|x|^{\beta}} \psi(x) \, dx.$$
 (3.8)

Hence U is a weak solution of (2.3).

The set $\mathbb{R}^n \cup \{\infty\}$ is compact for the standard topology which means that the measures can be identified as the dual space $C(\mathbb{R}^n \cup \{\infty\})$. For example, δ_{∞} is well defined and $\delta_{\infty} = \varphi(\infty)$. By the concentration compactness principle [25, 26], there exist a subsequence, still denoted by U_k and real numbers $\mu_0, \mu_\infty, \nu_0, \nu_\infty, \eta_0, \eta_\infty$ and ζ_0, ζ_∞ such that

$$\|U_k\|_{X^s(\mathbb{R}^{n+1}_+)}^2 \rightharpoonup d\mu \ge \|U\|_{X^s(\mathbb{R}^{n+1}_+)}^2 + \mu_0 \delta_0 + \mu_\infty \delta_\infty,$$
(3.9)

$$|u_k|^2 |x|^{-2s} \rightharpoonup d\nu = |u|^2 |x|^{-2s} + \nu_0 \delta_0 + \nu_\infty \delta_\infty, \tag{3.10}$$

$$|u_k|^{2^*_s(\alpha)}|x|^{-\alpha} \rightharpoonup d\eta = |u|^{2^*_s(\alpha)}|x|^{-\alpha} + \eta_0\delta_0 + \eta_\infty\delta_\infty, \tag{3.11}$$

$$|u_k|^{2^*_s(\beta)}|x|^{-\beta} \rightharpoonup d\zeta = |u|^{2^*_s(\beta)}|x|^{-\beta} + \zeta_0\delta_0 + \zeta_\infty\delta_\infty, \qquad (3.12)$$

where δ_0 and δ_{∞} are the Dirac mass at the origin and infinity respectively.

For $\varrho > 0$, define $B_{\varrho}^+ := \{(x, y) \in \mathbb{R}^{n+1}_+ : |(x, y)| < \varrho\}$, $B_{\varrho} := \{x \in \mathbb{R}^n : |x| < \varrho\}$ and let $\Phi \in C_0^{\infty}(\mathbb{R}^{n+1}_+)$ be a cut-off function such that $\Phi \equiv 1$ in $B_{\frac{1}{2}}^+$ and $0 \le \Phi \le 1$ in \mathbb{R}^{n+1}_+ . We use ΦU_k as test function, we have

$$\begin{split} \langle J'(U_{k}), \Phi U_{k} \rangle \\ &= k_{s} \int_{\mathbb{R}^{n+1}_{+}} y^{1-2s} \langle \nabla U_{k}, \nabla(\Phi U_{k}) \rangle \, dx \, dy - \gamma \int_{\mathbb{R}^{n}} \frac{u_{k}(x)^{2} \phi(x)}{|x|^{2s}} \, dx \\ &- \int_{\mathbb{R}^{n}} \frac{|u_{k}(x)|^{2^{*}_{s}(\alpha)} \phi(x)}{|x|^{\alpha}} \, dx - \int_{\mathbb{R}^{n}} \frac{|u_{k}(x)|^{2^{*}_{s}(\beta)} \phi(x)}{|x|^{\beta}} \, dx \\ &= k_{s} \int_{\mathbb{R}^{n+1}_{+}} y^{1-2s} |\nabla U_{k}|^{2} \Phi(x) \, dx \, dy - \gamma \int_{\mathbb{R}^{n}} \frac{u_{k}(x)^{2} \phi(x)}{|x|^{2s}} \, dx \\ &+ k_{s} \int_{\mathbb{R}^{n+1}_{+}} y^{1-2s} U_{k} \langle \nabla U_{k}, \nabla \Phi \rangle \, dx \, dy \\ &- \int_{\mathbb{R}^{n}} \frac{|u_{k}(x)|^{2^{*}_{s}(\alpha)} \phi(x)}{|x|^{\alpha}} \, dx - \int_{\mathbb{R}^{n}} \frac{|u_{k}(x)|^{2^{*}_{s}(\beta)} \phi(x)}{|x|^{\beta}} \, dx, \end{split}$$
(3.13)

where $\phi = \Phi(\cdot, 0)$. First, we have

$$\lim_{\varrho \to 0} \lim_{k \to \infty} \left(k_s \int_{\mathbb{R}^{n+1}_+} y^{1-2s} U_k \langle \nabla U_k, \nabla \Phi \rangle \, dx \, dy \right) = 0.$$

Moreover, from (3.9)-(3.12), we obtain

$$\lim_{\varrho \to 0} \lim_{k \to \infty} \left(k_s \int_{\mathbb{R}^{n+1}_+} y^{1-2s} |\nabla U_k|^2 \Phi \, dx \, dy \right) \ge \mu_0,$$
$$\lim_{\varrho \to 0} \lim_{k \to \infty} \int_{\mathbb{R}^n} \frac{u_k(x)^2 \phi(x)}{|x|^{2s}} \, dx = \nu_0,$$
$$\lim_{\varrho \to 0} \lim_{k \to \infty} \int_{\mathbb{R}^n} \frac{|u_k|^{2^*_s(\alpha)} \phi(x)}{|x|^{\alpha}} \, dx = \eta_0, \quad \lim_{\varrho \to 0} \lim_{k \to \infty} \int_{\mathbb{R}^n} \frac{|u_k|^{2^*_s(\beta)} \phi(x)}{|x|^{\beta}} \, dx = \zeta_0.$$

Thus we obtain

$$\lim_{\varrho \to 0} \lim_{k \to \infty} \langle J'(U_k), \Phi U_k \rangle \ge \mu_0 - \gamma \nu_0 - \eta_0 - \zeta_0.$$
(3.14)

By the fractional Hardy-Sobolev inequalities, we have

$$\eta_0^{\frac{2}{2^*_s(\alpha)}} S(n, s, \gamma, \alpha) \le \mu_0 - \gamma \nu_0, \quad \zeta_0^{\frac{2}{2^*_s(\beta)}} S(n, s, \gamma, \beta) \le \mu_0 - \gamma \nu_0.$$
(3.15)

By (3.14) and (3.15), we find

$$\eta_0^{\frac{2}{2_s^*(\alpha)}} S(n, s, \gamma, \alpha) \le \eta_0 + \zeta_0, \quad \zeta_0^{\frac{2}{2_s^*(\beta)}} S(n, s, \gamma, \beta) \le \eta_0 + \zeta_0.$$
(3.16)

 So

$$\begin{split} &\eta_0^{\frac{2}{2_s^*(\alpha)}} \left(1 - S(n, s, \gamma, \alpha)^{-1} \eta_0^{\frac{2_s^*(\alpha) - 2}{2_s^*(\alpha)}}\right) \le S(n, s, \gamma, \alpha)^{-1} \zeta_0, \\ &\zeta_0^{\frac{2}{2_s^*(\beta)}} \left(1 - S(n, s, \gamma, \beta)^{-1} \zeta_0^{\frac{2_s^*(\beta) - 2}{2_s^*(\beta)}}\right) \le S(n, s, \gamma, \beta)^{-1} \eta_0. \end{split}$$

Since $\{U_k\}_{k\in\mathbb{N}}$ is bounded in $X^s(\mathbb{R}^{n+1}_+)$, we have $\eta_0 \leq c_1$ and $\zeta_0 \leq c_2$ for positive constants c_1, c_2 , thus

$$\eta_0^{\frac{2}{2_s^*(\alpha)}} \left(1 - S(n, s, \gamma, \alpha)^{-1} c_1^{\frac{2_s^*(\alpha) - 2}{2_s^*(\alpha)}} \right) \le S(n, s, \gamma, \alpha)^{-1} \zeta_0,$$

$$\zeta_0^{\frac{2}{2_s^*(\beta)}} \left(1 - S(n, s, \gamma, \beta)^{-1} c_2^{\frac{2_s^*(\beta) - 2}{2_s^*(\beta)}} \right) \le S(n, s, \gamma, \beta)^{-1} \eta_0.$$

Therefore, there exist constants $A = A(\alpha, 2^*_s(\alpha), c_1)$ and $B = B(\beta, 2^*_s(\beta), c_2)$ such that

$$\eta_0^{\frac{2}{2s}(\alpha)} \le A\zeta_0, \quad \text{and} \quad \zeta_0^{\frac{2}{2s}(\beta)} \le B\eta_0.$$

In particular, we have that either $\eta_0 = 0$ and $\zeta_0 = 0$, or

$$\eta_0 \ge S(n, s, \gamma, \alpha)^{\frac{n-\alpha}{2s-\alpha}}, \quad \zeta_0 \ge S(n, s, \gamma, \beta)^{\frac{n-\beta}{2s-\beta}}.$$

On the other hand, we know that

$$c = J(U_k) - \frac{1}{2} \langle J'(U_k), U_k \rangle + o_k(1)$$

$$\geq \frac{2s - \alpha}{2(n - \alpha)} \Big(\int_{\mathbb{R}^n} \frac{|u_k(x)|^{2^*_s(\alpha)}}{|x|^{\alpha}} dx + \eta_0 \Big)$$

$$+ \frac{2s - \beta}{2(n - \beta)} \Big(\int_{\mathbb{R}^n} \frac{|u_k(x)|^{2^*_s(\beta)}}{|x|^{\beta}} dx + \zeta_0 \Big)$$

$$\geq \frac{2s - \alpha}{2(n - \alpha)} \eta_0 + \frac{2s - \beta}{2(n - \beta)} \zeta_0.$$
(3.17)

By the assumption that $c < c_*$, we obtain that $\eta_0 = 0$, $\zeta_0 = 0$.

For the concentration at infinity, we define $B_R^+ := \{(x,y) \in \mathbb{R}^{n+1}_+ : |(x,y)| < R\},\ B_R := \{x \in \mathbb{R}^n : |x| < R\}$ and let $\Psi \in C_0^\infty(\mathbb{R}^{n+1}_+)$ be a cut-off function such that $\Psi = 0$ in B_R^+ and $\Psi \equiv 1$ in $\mathbb{R}^{n+1}_+ \setminus B_{2R}^+$ and $0 \le \Psi \le 1$ in \mathbb{R}^{n+1}_+ . Consider

$$\mu_{\infty} = \lim_{R \to \infty} \lim_{k \to \infty} \sup \left(k_s \int_{\mathbb{R}^{n+1} \setminus B_{2R}^+} y^{1-2s} |\nabla U_k|^2 \Psi \, dx \, dy \right)$$
$$\nu_{\infty} = \lim_{R \to \infty} \lim_{k \to \infty} \sup \int_{\mathbb{R}^n \setminus B_{2R}} \frac{u_k(x)^2 \psi(x)}{|x|^{2s}} \, dx,$$
$$\eta_{\infty} = \lim_{R \to \infty} \lim_{k \to \infty} \sup \int_{\mathbb{R}^n \setminus B_{2R}} \frac{|u_k(x)|^{2^*_s(\alpha)} \psi(x)}{|x|^{\alpha}} \, dx,$$
$$\zeta_{\infty} = \lim_{R \to \infty} \lim_{k \to \infty} \sup \int_{\mathbb{R}^n \setminus B_{2R}} \frac{|u_k(x)|^{2^*_s(\beta)} \psi(x)}{|x|^{\beta}} \, dx.$$

By the same arguments as the concentration at the origin, we can get the following facts: either $\eta_{\infty} = 0$ and $\zeta_{\infty} = 0$, or

$$\eta_{\infty} \ge S(n, s, \gamma, \alpha)^{\frac{n-\alpha}{2s-\alpha}}, \quad \zeta_{\infty} \ge S(n, s, \gamma, \beta)^{\frac{n-\beta}{2s-\beta}}.$$

As for (3.17), we obtain

$$c \ge \frac{2s - \alpha}{2(n - \alpha)} \eta_{\infty} + \frac{2s - \beta}{2(n - \beta)} \zeta_{\infty}.$$
(3.18)

By the assumption that $c < c_*$, we obtain that $\eta_{\infty} = 0$, $\zeta_{\infty} = 0$. Therefore, up to a subsequence $\{U_k\}_k$ converges strongly to U in $X^s(\mathbb{R}^{n+1}_+)$.

10

9

Let $W_{\gamma,\alpha}$ be the extremal function of $S(n, s, \gamma, \alpha)$ in $X^s(\mathbb{R}^{n+1}_+)$, whose existence was obtained by Ghoussoub and Shakerian in [21] for $\alpha > 0$ or $\alpha = 0$ and $0 \le \gamma < \gamma_H$.

Lemma 3.3. Let 0 < s < 1, $0 < \alpha, \beta < 2s < n$, and $\gamma < \gamma_H$. Then

$$\sup_{t \ge 0} J(tW_{\gamma,\vartheta}) < c_* \quad for \ \vartheta = \alpha, \beta,$$

where c_* is defined in Proposition 3.2.

Proof. For $\vartheta = \alpha$, we have

$$J(tW_{\gamma,\alpha}) = \frac{t^2}{2} \|W_{\gamma,\alpha}\|^2 - \frac{t^{2^*_s(\alpha)}}{2^*_s(\alpha)} \int_{\mathbb{R}^n} \frac{|w_{\gamma,\alpha}|^{2^*_s(\alpha)}}{|x|^{\alpha}} \, dx - \frac{t^{2^*_s(\beta)}}{2^*_s(\beta)} \int_{\mathbb{R}^n} \frac{|w_{\gamma,\alpha}|^{2^*_s(\beta)}}{|x|^{\beta}} \, dx.$$

where $w_{\gamma,\alpha} := Tr(W_{\gamma,\alpha}) = W_{\gamma,\alpha}(\cdot, 0)$. By construction, we have that

$$J(tW_{\gamma,\alpha}) \le f_{\alpha}(t) := \frac{t^2}{2} \|W_{\gamma,\alpha}\|^2 - \frac{t^{2^*_s(\alpha)}}{2^*_s(\alpha)} \int_{\mathbb{R}^n} \frac{|w_{\gamma,\alpha}|^{2^*_s(\alpha)}}{|x|^{\alpha}} dx$$

Straightforward computations yield that $f_{\alpha}(t)$ attains its maximum at the point

$$\tilde{t} = \left(\frac{\|W_{\gamma,\alpha}\|^2}{\int_{\mathbb{R}^n} \frac{|w_{\gamma,\alpha}|^{2^*_s(\alpha)}}{|x|^{\alpha}} dx}\right)^{\frac{1}{2^*_s(\alpha)-2}}$$

It follows that

$$\sup_{t \ge 0} f_{\alpha}(t) = \frac{2s - \alpha}{2(n - \alpha)} \Big(\frac{\|W_{\gamma, \alpha}\|^2}{\left(\int_{\mathbb{R}^n} |w_{\gamma, \alpha}|^{2^*_s(\alpha)} / |x|^{\alpha} \, dx\right)^{\frac{2}{2^*_s(\alpha)}}} \Big)^{\frac{n - \alpha}{2s - \alpha}}$$

Since $W_{\gamma,\alpha}$ is an extremal for $S(n, s, \gamma, \alpha)$ on $X^s(\mathbb{R}^{n+1}_+)$, we obtain that

$$\sup_{t \ge 0} J(tW_{\gamma,\alpha}) \le \sup_{t \ge 0} f_{\alpha}(t) = \frac{2s - \alpha}{2(n - \alpha)} S(n, s, \gamma, \alpha)^{\frac{n - \alpha}{2s - \alpha}}.$$
(3.19)

We now need to show that equality does not hold in (3.19). Indeed, otherwise we would have that $\sup_{t\geq 0} J(tW_{\gamma,\alpha}) = \sup_{t\geq 0} f_{\alpha}(t)$. Consider t_1 (resp. $t_2 > 0$) where $\sup_{t\geq 0} J(tW_{\gamma,\alpha})$ (resp. $\sup_{t\geq 0} f_{\alpha}(t)$) is attained. We obtain

$$f_{\alpha}(t_1) - \frac{t_1^{2^*_s(\beta)}}{2^*_s(\beta)} \int_{\mathbb{R}^n} \frac{|w_{\gamma,\alpha}|^{2^*_s(\beta)}}{|x|^{\beta}} \, dx = f_{\alpha}(t_2),$$

which means that $f_{\alpha}(t_1) > f_{\alpha}(t_2)$ since $t_1 > 0$. This contradicts the fact that t_2 is a maximum point of $f_{\alpha}(t)$, hence the strict inequality holds in (3.19).

Similarly, for $\vartheta = \beta$, we obtain

$$\sup_{t\geq 0} J(tW_{\gamma,\beta}) < \sup_{t\geq 0} f_{\beta}(t) = \frac{2s-\beta}{2(n-\beta)} S(n,s,\gamma,\beta)^{\frac{n-\beta}{2s-\beta}}.$$

This completes the proof.

4. Proof of main result

Proof of Theorem 1.1. For any $U \in X^s(\mathbb{R}^{n+1}_+)$, the energy functional to problem (2.3) is

$$J(U) = \frac{1}{2} \|U\|^2 - \frac{1}{2_s^*(\alpha)} \int_{\mathbb{R}^n} \frac{|u|^{2_s^*(\alpha)}}{|x|^{\alpha}} \, dx - \frac{1}{2_s^*(\beta)} \int_{\mathbb{R}^n} \frac{|u|^{2_s^*(\beta)}}{|x|^{\beta}} \, dx.$$

where again $u := Tr(U) = U(\cdot, 0)$. By fractional Hardy-Sobolev inequality, we have

$$\begin{split} J(U) &\geq \frac{1}{2} \|U\|^2 - \frac{1}{2_s^*(\alpha)} S(n, s, \gamma, \alpha)^{-\frac{2_s^*(\alpha)}{2}} \|U\|^{2_s^*(\alpha)} \\ &- \frac{1}{2_s^*(\beta)} S(n, s, \gamma, \beta)^{-\frac{2_s^*(\beta)}{2}} \|U\|^{2_s^*(\beta)} \\ &= \left(\frac{1}{2} - \frac{1}{2_s^*(\alpha)} S(n, s, \gamma, \alpha)^{-\frac{2_s^*(\alpha)}{2}} \|U\|^{2_s^*(\alpha)-2} \\ &- \frac{1}{2_s^*(\beta)} S(n, s, \gamma, \beta)^{-\frac{2_s^*(\beta)}{2}} \|U\|^{2_s^*(\beta)-2} \right) \|U\|^2. \end{split}$$

Since $\alpha, \beta \in (0, 2s)$, we have that $2_s^*(\alpha) > 2, 2_s^*(\beta) > 2$. By (2.11), we then get that there exists R > 0 such that $J(U) \ge \rho$ for all $U \in X^s(\mathbb{R}^{n+1}_+)$ with $\|U\|_{X^{\alpha}(\mathbb{R}^{n+1}_+)} = R$. Moreover, for $\vartheta = \alpha$ or $\vartheta = \beta$,

$$J(tW_{\gamma,\vartheta}) = \frac{t^2}{2} \|W_{\gamma,\vartheta}\|^2 - \frac{t^{2^*_s(\alpha)}}{2^*_s(\alpha)} \int_{\mathbb{R}^n} \frac{|w_{\gamma,\vartheta}|^{2^*_s(\alpha)}}{|x|^{\alpha}} \, dx - \frac{t^{2^*_s(\beta)}}{2^*_s(\beta)} \int_{\mathbb{R}^n} \frac{|w_{\gamma,\vartheta}|^{2^*_s(\beta)}}{|x|^{\alpha}} \, dx,$$

hence $\lim_{t\to+\infty}J(tW_{\gamma,\vartheta})=-\infty$, then there exists $t_0>0$ such that $\|t_0W_{\gamma,\vartheta}\|>R$ and $J(t_0W_{\gamma,\vartheta})<0.$ Set

$$c_{\vartheta} := \inf_{g \in \Gamma_{\vartheta}} \max_{t \in [0,1]} J(g(t)),$$

where

$$\Gamma_{\vartheta} := \left\{ g \in C^0([0,1], X^s(\mathbb{R}^{n+1}_+)) : g(0) = 0, \ g(1) = t_0 W_{\gamma,\vartheta} \right\}.$$

Thus by Mountain Pass Lemma, there exists a sequence $\{U_k\}$ in $X^s(\mathbb{R}^{n+1}_+)$ such that

$$J(U_k) \to c$$
, $J'(U_k) \to 0$ in $(X^s(\mathbb{R}^{n+1}_+))'$ as $k \to \infty$.

By Lemma 3.3, we have

$$0 < c \le \sup_{t \in [0,1]} J(tt_0 W_{\gamma,\vartheta}) \le \sup_{t>0} J(tW_{\gamma,\vartheta}) < c_*.$$

By Proposition 3.2, we deduce that $\{U_k\}$ has a subsequence, still denote by $\{U_k\}$, such that $U_k \to U$ strongly in $X^s(\mathbb{R}^{n+1}_+)$. Thus U is a nontrivial solution of problem (2.3), and $u := Tr(U) = U(\cdot, 0)$ is a nontrivial solution of problem (1.1). \Box

Acknowledgments. This research was supported by the NSFC No. 11501468, by the Chongqing Research Program of Basic Research and Frontier Technology cstc2016jcyjA0323, by the Fundamental Research Funds for the Central Universities XDJK2017C049, and by the innovation support program of Chongqing cx2017070.

References

- B. Abdellaoui, I. Peral, A. Primo; A remark on the fractional Hardy inequality with a remainder term, C. R. Math. Acad. Sci. Paris, 352 (4) (2014), 299-303.
- [2] B. Abdellaoui, F. Mahmoudi; An improved Hardy inequality for a nonlocal operator, Discrete Contin. Dyn. Syst., 36 (2016), no. 3, 1143-1157.
- [3] C. Alves, O. Miyagaki; A critical nonlinear fractional elliptic equation with saddle-like potential in R^N, J. Math. Phys., 57 (2016), no. 8, 081501, 17 pp.
- [4] B. Barrios, M. Medina, I. Peral; Some remarks on the solvability of non-local elliptic problems with the Hardy potential, *Communications in Contemporary Mathematics*, 16.04 (2014).
- [5] M. Bouchekif, S. Messirdi; On elliptic problems with two critical Sobolev-Hardy exponents at the same pole, *Applied Mathematics Letters*, 42 (2015) 9-14.
- [6] L. Caffarelli, L. Silvestre; An extension problem related to the fractional Laplacian, Communications in Partial Differential Equations, 32 (2007), 1245-1260.
- [7] F. Catrina, Z. Wang; On the Caffarelli-Kohn-Nirenberg inequalities: sharp constants, existence (and nonexistence), and simmetry of extremal functions, *Comm. Pure Appl. Math.*, 54 (2001), 229-258.
- [8] W. Chen, S. Deng; The Nehari manifold for nonlocal elliptic operators involving concaveconvex nonlinearities, Z. Angew. Math. Phys., 66 (2015), no. 4, 1387-1400.
- [9] W. Chen, S. Deng; The Nehari manifold for a fractional p-Laplacian system involving concaveconvex nonlinearities, *Nonlinear Anal. Real World Appl.*, 27 (2016), 80-92.
- [10] W. Chen, S. Deng; Multiple solutions for a critical fractional elliptic system involving concaveconvex nonlinearities, Proc. Roy. Soc. Edinburgh Sect. A, 146 (2016), no. 6, 1167-1193.
- [11] W. Chen, M. Squassina; Critical nonlocal systems with concave-convex powers, Adv. Nonlinear Stud., 16 (2016), no. 4, 821-842.
- [12] J. L. Chern, C. S. Lin; Minimizers of Caffarelli-Kohn-Nirenberg inequalities with the singularity on the boundary, Arch. Ration. Mech. Anal., 197 (2010), no. 2, 401-432.
- [13] A. Cotsiolis, N. K. Tavoularis; Best constants for Sobolev inequalities for higher order fractional derivatives J. Math. Anal. Appl., 295 (2004), pp. 225-236.
- [14] E. Di Nezza, G. Palatucci, E. Valdinoci; Hitchhikers guide to the fractional Sobolev spaces, Bull. Sci. Math., 136 (2012), no. 5, 521-573.
- [15] M. M. Fall; Semilinear elliptic equations for the fractional Laplacian with Hardy potential, (2011); arXiv: 1109.5530v4.
- [16] M. M. Fall, I. A. Minlend, E. A. Thiam; The role of the mean curvature in a Sobolev-Hardy trace inequality, Nonlinear Differential Equations & Application, 22(5)(2014), 1-20.
- [17] R. Filippucci, P. Pucci, F. Robert; On a p-Laplace equation with multiple critical nonlinearities, J. Math. Pures Appl., 9(2) 91 (2009), 156-177.
- [18] A. Fiscella; Saddle point solutions for non-local elliptic operators, Topological Methods in Nonlinear Analysis, 44(2) (2014), 527-538.
- [19] A. Fiscella, G. Molica Bisci, R. Servadei; Bifurcation and multiplicity results for critical nonlocal fractional Laplacian problems, *Bull. Sci. Math.*, 140 (2016), no. 1, 14-35.
- [20] R. L. Frank, E. H. Lieb, R. Seiringer; Hardy-Lieb-Thirring inequalities for fractional Schrödinger operators, J. Amer. Math. Soc., 21(4) (2008) 925-950.
- [21] N. Ghoussoub, S. Shakerian; Borderline variational problems involving fractional Laplacians and critical singularities, Advanced Nonlinear Studies, 15(3)(2016), 93-100.
- [22] N. Ghoussoub, C. Yuan; Multiple solutions for quasi-linear PDEs involving the critical Sobolev and Hardy exponents, *Trans. Amer. Math. Soc.*, 12 (2000), 5703-5743.
- [23] H. Hajaiej, G. Molica Bisci, L. Vilasi; Existence results for a critical fractional equation, Asymptot. Anal., 100 (2016), no. 3-4, 209-225.
- [24] Y. Lei; Asymptotic properties of positive solutions of the Sobolev-Hardy type equations, J. Differential Equations, 253(2013), 1774-1799.
- [25] P. L. Lions; The concentration-compactness principle in the calculus of variations. The limit case II, Revista Mathematica Iberoamericana, 1(2)(1985), 45-121.
- [26] P. L. Lions; The concentration-compactness principle in the calculus of variations. The locally compact case, Ann. I. H. P., Anal. Nonlin., 1(1984), 273-288.
- [27] G. Molica Bisci, V. Radulescu, R. Servadei; Variational Methods for Nonlocal Fractional Problems, Encyclopedia of Mathematics and its Applications, 162 Cambride University Press - ISBN: 9781107111943. Foreword by J. Mawhin. (pp. 1-400).

- [28] G. Molica Bisci, R. Servadei; Lower semicontinuity of functionals of fractional type and applications to nonlocal equations with critical Sobolev exponent, Adv. Differential Equations, 20(2015), 635-660.
- [29] R. Servadei, E. Valdinoci; Lewy-Stampacchia type estimates for variational inequalities driven by (non)local operators, *Rev. Mat. Iberoam.*, 29(3)(2013), 1091-1126.
- [30] R. Servadei, E. Valdinoci; Mountain Pass solutions for non-local elliptic operators, J. Math. Anal. Appl., 389(2012), 887–898.
- [31] R. Servadei, E. Valdinoci; Variational methods for non-local operators of elliptic type, Discrete and Continuous Dynamical Systems, 5(2013), 2105-2137.
- [32] R. Servadei, E. Valdinoci; The Brezis-Nirenberg result for the fractional Laplacian, Trans. Amer. Math. Soc. 367(2015), 67-102.
- [33] R. Servadei, E. Valdinoci; A Brezis-Nirenberg result for non-local critical equations in low dimension, Commun. Pure Appl. Anal., 12(6)(2013), 2445-2464.
- [34] X. Wang, J. Yang; Singular critical elliptic problems with fractional Laplacian, *Electron. J. Diff. Equ.*, Vol. 2015 (2015), No. 297, pp. 1-12.
- [35] J. Yang; Fractional Hardy-Sobolev inequality in R^N, Nonlinear Analysis: Theory, Methods and Applications, (119)2015, 179-185.
- [36] J. Yang, F. Wu; Doubly Critical Problems Involving Fractional Laplacians in R^N, to appear in Adv. Nonlinear Stud., 2017,
- [37] J. Yang, X. Yu; Fractional Hardy-Sobolev elliptic problems, preprint, https://arxiv.org/abs/1503.00216.

Wenjing Chen

SCHOOL OF MATHEMATICS AND STATISTICS, SOUTHWEST UNIVERSITY, CHONGQING 400715, CHINA *E-mail address*: wjchen@swu.edu.cn