

**EXACT ASYMPTOTIC BEHAVIOR OF THE POSITIVE
 SOLUTIONS FOR SOME SINGULAR DIRICHLET PROBLEMS
 ON THE HALF LINE**

HABIB MÂAGLI, RAMZI ALSAEDI, NOUREDDINE ZEDDINI

ABSTRACT. In this article, we give an exact behavior at infinity of the unique solution to the following singular boundary value problem

$$-\frac{1}{A}(Au')' = q(t)g(u), \quad t \in (0, \infty),$$

$$u > 0, \quad \lim_{t \rightarrow 0} A(t)u'(t) = 0, \quad \lim_{t \rightarrow \infty} u(t) = 0.$$

Here A is a nonnegative continuous function on $[0, \infty)$, positive and differentiable on $(0, \infty)$ such that

$$\lim_{t \rightarrow \infty} \frac{tA'(t)}{A(t)} = \alpha > 1, \quad g \in C^1((0, \infty), (0, \infty))$$

is non-increasing on $(0, \infty)$ with $\lim_{t \rightarrow 0} g'(t) \int_0^t \frac{ds}{g(s)} = -C_g \leq 0$ and the function q is a nonnegative continuous, satisfying

$$0 < a_1 = \liminf_{t \rightarrow \infty} \frac{q(t)}{h(t)} \leq \limsup_{t \rightarrow \infty} \frac{q(t)}{h(t)} = a_2 < \infty,$$

where $h(t) = ct^{-\lambda} \exp(\int_1^t \frac{y(s)}{s} ds)$, $\lambda \geq 2$, $c > 0$ and y is continuous on $[1, \infty)$ such that $\lim_{t \rightarrow \infty} y(t) = 0$.

1. INTRODUCTION

In this article, we give the exact asymptotic behavior at infinity of the unique positive solution to the singular problem

$$\frac{1}{A}(Au')' = -q(t)g(u), \quad t \in (0, \infty),$$

$$u > 0, \quad \text{in } (0, \infty) \tag{1.1}$$

$$\lim_{t \rightarrow 0^+} A(t)u'(t) = 0, \quad \lim_{t \rightarrow \infty} u(t) = 0,$$

where the functions A , q and g satisfy the following assumptions.

(H1) A is a continuous function on $[0, \infty)$, positive and differentiable on $(0, \infty)$ such that

$$\lim_{t \rightarrow \infty} \frac{tA'(t)}{A(t)} = \alpha > 1.$$

2010 *Mathematics Subject Classification.* 34B16, 34B18, 34D05.

Key words and phrases. Singular nonlinear boundary value problems; positive solution; exact asymptotic behavior; Karamata regular variation theory.

©2016 Texas State University.

Submitted December 8, 2015. Published February 17, 2016.

(H2) q is a nonnegative continuous function on $(0, \infty)$ satisfying

$$0 < a_1 = \liminf_{t \rightarrow \infty} \frac{t^\lambda q(t)}{L(t)} \leq \limsup_{t \rightarrow \infty} \frac{t^\lambda q(t)}{L(t)} = a_2 < \infty,$$

where $\lambda \geq 2$ and $L \in \mathcal{K}$ (see (1.3) below), such that $\int_1^\infty s^{1-\lambda} L(s) ds < \infty$.

(H3) The function $g : (0, \infty) \rightarrow (0, \infty)$ is nonincreasing, continuously differentiable such that

$$\lim_{t \rightarrow 0^+} g'(t) \int_0^t \frac{1}{g(s)} ds = -C_g \quad \text{with } C_g \geq 0.$$

(H4) $\alpha + 1 - \lambda + (\lambda - 2)C_g > 0$.

Using that g is non-increasing, for $t > 0$, we obtain

$$0 < g(t) \int_0^t \frac{1}{g(s)} ds \leq t.$$

This implies $\lim_{t \rightarrow 0} g(t) \int_0^t \frac{1}{g(s)} ds = 0$. Now, since for $t > 0$,

$$\int_0^t g'(s) \int_0^s \frac{1}{g(r)} dr ds = g(t) \int_0^t \frac{1}{g(s)} ds - t,$$

we obtain

$$\lim_{t \rightarrow 0} \frac{g(t)}{t} \int_0^t \frac{1}{g(s)} ds = 1 - C_g. \quad (1.2)$$

This implies that $0 \leq C_g \leq 1$. The functions $t^{-1} \log(1 + t)$, $\log(\log(e + \frac{1}{t}))$, $t^{-\nu} \log(1 + \frac{1}{t})$, $\exp\{(\log(1 + \frac{1}{t}))^\nu\}$, $\nu \in (0, 1)$ satisfy the assumption (H3), as well as the function

$$\begin{aligned} t^2 e^{1/t}, & \quad \text{if } 0 < t < \frac{1}{2}, \\ \frac{1}{4} e^2, & \quad \text{if } t \geq \frac{1}{2}. \end{aligned}$$

Singular nonlinear boundary value problems appear in a variety of applications and often only positive solutions are important. When $A(t) = 1$, problems of type (1.1) with various boundary conditions arise in the study of boundary layer equations for the class of pseudoplastic fluids and have been studied for both bounded and unbounded intervals of \mathbb{R} (see [6, 11, 16, 22, 23, 29]) and the references therein. When $A(t) = t^{n-1}$ ($n \geq 1$), the operator $u \rightarrow \frac{1}{A}(Au)'$ appears as the radial part of the laplace operator Δ (see [10, 30]). Other results of existence and uniqueness of positive solutions were obtained by Agarwal and O'Regan in [1] on the interval $(0, 1)$ and in the case where A is continuous on $[0, 1]$, positive and differentiable on $(0, 1)$ and satisfying an integrability condition. In general the exact asymptotic behavior of the unique positive solution of (1.1) is extremely complex when the coefficients are in general continuous functions, even though upper and lower bounds for this solution are often given (see [1, 4, 10, 15]). Recent research (see [2, 8, 16]) show that these problems should be studied in the case of Karamata regularly varying functions. This approach was initiated by Avakumovic [3] and followed by Maric and Tomic (see [20, 21]). Our aim in this paper is to give a contribution to the qualitative analysis of problem (1.1) by giving the exact asymptotic behavior at infinity of the unique positive solution under the previous assumptions on A , q and g . We note that the existence and uniqueness of such a solution are established by Mâagli and Masmoudi in [17]. For related results, we refer to Barile and Salvatore

[5], Cencelj, Repovš, and Virk [7], Ghergu and Rădulescu [13, 14, 15, 16], Rădulescu and Repovš [24, 25], Repovš [26, 27].

To state our results, we denote by \mathcal{K} the set of Karamata functions L defined on $[1, \infty)$ by

$$L(t) := c \exp \left(\int_1^t \frac{y(s)}{s} ds \right), \quad (1.3)$$

where $c > 0$ and $y \in C([1, \infty))$ such that $\lim_{t \rightarrow \infty} y(t) = 0$.

Remark 1.1. It is clear that a function L is in \mathcal{K} if and only if L is a positive function in $C^1([1, \infty))$ such that

$$\lim_{t \rightarrow \infty} \frac{t L'(t)}{L(t)} = 0. \quad (1.4)$$

Throughout this paper, we denote by ψ_g the unique solution of the equation

$$\int_0^{\psi_g(t)} \frac{ds}{g(s)} = t, \quad \text{for } t \in [0, \infty), \quad (1.5)$$

and we mention that

$$\lim_{t \rightarrow 0} t g'(\psi_g(t)) = -C_g. \quad (1.6)$$

Theorem 1.2. Assume (H1)–(H4). Then problem (1.1) has a unique solution $u \in C^2((0, \infty)) \cap C([0, \infty))$ satisfying

(i) If $\lambda > 2$,

$$\left(\frac{\xi_1}{\lambda - 2} \right)^{1-C_g} \leq \liminf_{t \rightarrow \infty} \frac{u(x)}{\psi_g(t^{2-\lambda} L(t))} \leq \limsup_{t \rightarrow \infty} \frac{u(x)}{\psi_g(t^{2-\lambda} L(t))} \leq \left(\frac{\xi_2}{\lambda - 2} \right)^{1-C_g},$$

where $\xi_i = \frac{a_i}{\alpha+1-\lambda+(\lambda-2)C_g}$ for $i \in \{1, 2\}$.

(ii) If $\lambda = 2$,

$$\xi_1^{1-C_g} \leq \liminf_{t \rightarrow \infty} \frac{u(t)}{\psi_g(\int_t^\infty \frac{L(s)}{s} ds)} \leq \limsup_{t \rightarrow \infty} \frac{u(t)}{\psi_g(\int_t^\infty \frac{L(s)}{s} ds)} \leq \xi_2^{1-C_g}$$

An immediate consequence of Theorem 1.2 is the following result.

Corollary 1.3. Let u be the unique solution of (1.1). Then, we have the following exact asymptotic behavior:

(a) When $C_g = 1$, we have

$$(i) \lim_{t \rightarrow \infty} \frac{u(t)}{\psi_g(t^{2-\lambda} L(t))} = 1 \text{ if } \lambda > 2;$$

$$(ii) \lim_{t \rightarrow \infty} \frac{u(t)}{\psi_g(\int_t^\infty \frac{L(s)}{s} ds)} = 1 \text{ if } \lambda = 2.$$

(b) When $C_g < 1$ and $a_1 = a_2 = a_0$, we have

$$(i) \lim_{t \rightarrow \infty} \frac{u(t)}{\psi_g(t^{2-\lambda} L(t))} = \left(\frac{a_0}{(\lambda-2)(\alpha+1-\lambda+(\lambda-2)C_g)} \right)^{1-C_g} \text{ if } \lambda > 2;$$

$$(ii) \lim_{t \rightarrow \infty} \frac{u(t)}{\psi_g(\int_t^\infty \frac{L(s)}{s} ds)} = \left(\frac{a_0}{\alpha-1} \right)^{1-C_g} \text{ if } \lambda = 2.$$

Remark 1.4. In the hypothesis (H3), we do not need the monotonicity of the function g on $(0, \infty)$, but only the fact that g is non-increasing in a neighborhood of zero.

Example 1.5. Let g be the function

$$g(t) = \begin{cases} t^2 e^{1/t}, & \text{if } 0 < t < \frac{1}{2}, \\ \frac{1}{4} e^2, & \text{if } t \geq \frac{1}{2}. \end{cases}$$

and let q be a nonnegative function in $(0, \infty)$ such that

$$\lim_{t \rightarrow \infty} \frac{q(t)}{h(t)} = b_0 \in (0, \infty),$$

where $h(t) = t^{-\lambda} L(t)$, $\lambda \geq 2$ and $L \in \mathcal{K}$ such that $\int_1^\infty s^{1-\lambda} L(s) ds < \infty$. Then, we have $C_g = 1$ and $\psi_g(\xi) = \frac{-1}{\log(\xi)}$ for $\xi \in (0, e^{-2})$. Let u be the unique solution of (1.1), then we have the following exact behavior:

- (i) $\lim_{t \rightarrow \infty} u(t) \log\left(\frac{1}{t^{2-\lambda} L(t)}\right) = 1$ if $\lambda > 2$;
- (ii) $\lim_{t \rightarrow \infty} u(t) \log\left(\frac{1}{\int_t^\infty \frac{L(s)}{s} ds}\right) = 1$ if $\lambda = 2$.

To establish our second result, we consider the special case where $g(t) = t^{-\gamma}$ with $\gamma \geq 0$, and $\lambda = \alpha + 1 + \gamma(\alpha - 1)$. Note that in this case $C_g = \frac{\gamma}{\gamma+1}$ and $(\alpha + 1 - \lambda) + (\lambda - 2)C_g = 0$. We assume that A and q satisfy the following hypotheses:

- (H5) A is a continuous function on $(0, \infty)$ such that $A(t) = t^\alpha B(t)$ with $\alpha > 1$ and $\frac{t^\nu B'(t)}{B(t)}$ is bounded for t large and $\nu \in (0, 1)$.
- (H6) q is a nonnegative continuous function in $(0, \infty)$ and satisfies

$$0 < a_1 = \liminf_{t \rightarrow \infty} \frac{q(t)}{t^{\gamma-1-\alpha(\gamma+1)} L(t)} \leq \limsup_{t \rightarrow \infty} \frac{q(t)}{t^{\gamma-1-\alpha(\gamma+1)} L(t)} = a_2 < \infty,$$

where $L \in \mathcal{K}$ with $\int_1^\infty \frac{L(s)}{s} ds = \infty$.

Theorem 1.6. *Assume (H5), (H6) are satisfied. Then the Dirichlet problem*

$$\begin{aligned} -\frac{1}{A}(Au')' &= q(t)u^{-\gamma}, \quad t \in (0, \infty), \\ \lim_{t \rightarrow 0^+} A(t)u'(t) &= 0, \quad \lim_{t \rightarrow \infty} u(t) = 0, \end{aligned} \tag{1.7}$$

has a unique solution $u \in C([0, \infty)) \cap C^2((0, \infty))$, satisfying

$$\begin{aligned} \left(\frac{(\gamma+1)a_1}{\alpha-1}\right)^{\frac{1}{1+\gamma}} &\leq \liminf_{t \rightarrow \infty} \frac{u(t)}{t^{1-\alpha} \left(\int_1^t \frac{L(s)}{s} ds\right)^{\frac{1}{1+\gamma}}} \\ &\leq \limsup_{t \rightarrow \infty} \frac{u(t)}{t^{1-\alpha} \left(\int_1^t \frac{L(s)}{s} ds\right)^{\frac{1}{1+\gamma}}} \\ &\leq \left(\frac{(\gamma+1)a_2}{\alpha-1}\right)^{\frac{1}{1+\gamma}}, \end{aligned}$$

In particular if $a_1 = a_2$, then

$$\lim_{t \rightarrow \infty} \frac{u(t)}{t^{1-\alpha} \left(\int_1^t \frac{L(s)}{s} ds\right)^{\frac{1}{1+\gamma}}} = \left(\frac{(\gamma+1)a_1}{\alpha-1}\right)^{\frac{1}{1+\gamma}}.$$

2. ON THE KARAMATA CLASS

To make the paper self-contained, we begin this section by recapitulating some properties of Karamata regular variation theory. The following result is due to [19, 28].

Lemma 2.1. (i) Let $L \in \mathcal{K}$ and $\varepsilon > 0$, then

$$\lim_{t \rightarrow \infty} t^{-\varepsilon} L(t) = 0.$$

(ii) Let $L_1, L_2 \in \mathcal{K}$ and $p \in \mathbb{R}$. Then $L_1 + L_2 \in \mathcal{K}$, $L_1 L_2 \in \mathcal{K}$ and $L_1^p \in \mathcal{K}$.

Applying Karamata's theorem (see [19, 28]), we get the following result.

Lemma 2.2. Let $\gamma \in \mathbb{R}$, L be a function in \mathcal{K} defined on $[1, \infty)$. We have

(i) If $\gamma < -1$, then $\int_1^\infty s^\gamma L(s) ds$ converges. Moreover

$$\int_t^\infty s^\gamma L(s) ds \sim_{t \rightarrow \infty} -\frac{t^{1+\gamma} L(t)}{\gamma + 1}.$$

(ii) If $\gamma > -1$, then $\int_1^\infty s^\gamma L(s) ds$ diverges. Moreover

$$\int_1^t s^\gamma L(s) ds \sim_{t \rightarrow \infty} \frac{t^{1+\gamma} L(t)}{\gamma + 1}.$$

Lemma 2.3 ([8, 18]). Let $L \in \mathcal{K}$ be defined on $[1, \infty)$. Then

$$\lim_{t \rightarrow \infty} \frac{L(t)}{\int_1^t \frac{L(s)}{s} ds} = 0. \quad (2.1)$$

If further $\int_1^\infty \frac{L(s)}{s} ds$ converges, then

$$\lim_{t \rightarrow \infty} \frac{L(t)}{\int_t^\infty \frac{L(s)}{s} ds} = 0. \quad (2.2)$$

Remark 2.4. Let $L \in \mathcal{K}$, then using Remark 1.1 and (2.1), we deduce that

$$t \rightarrow \int_1^t \frac{L(s)}{s} ds \in \mathcal{K}.$$

If further $\int_1^\infty \frac{L(s)}{s} ds$ converges, then $t \rightarrow \int_t^\infty \frac{L(s)}{s} ds \in \mathcal{K}$.

Definition 2.5. A positive measurable function k is called normalized regularly varying at infinity with index $\rho \in \mathbb{R}$ and we write $k \in NRV I_\rho$ if $k(s) = s^\rho L(s)$ for $s \in [1, \infty)$ with $L \in \mathcal{K}$.

Using the definition of the class \mathcal{K} and the above Lemmas we obtain the following lemma.

Lemma 2.6 ([2]). (i) If $k \in NRV I_\rho$, then $\lim_{t \rightarrow \infty} \frac{k(\xi t)}{k(t)} = \xi^\rho$, uniformly for $\xi \in [c_1, c_2] \subset (0, \infty)$.

(ii) A positive measurable function k belongs to the class $NRV I_\rho$ if and only if $\lim_{t \rightarrow \infty} \frac{tk'(t)}{k(t)} = \rho$.

(iii) Let $L \in \mathcal{K}$ and assume that $\int_1^\infty s^{1-\lambda} L(s) ds < \infty$. Then the function $\theta(t) = \int_t^\infty s^{1-\lambda} L(s) ds$ belongs to $NRV I_{(2-\lambda)}$.

(iv) The function $\psi_g \circ \theta \in NRV I_{(2-\lambda)(1-C_g)}$.

(v) Let m_1, m_2 be positive functions on $(0, \infty)$ such that $\lim_{t \rightarrow \infty} m_1(t) = \lim_{t \rightarrow \infty} m_2(t) = 0$ and $\lim_{t \rightarrow \infty} \frac{m_1(t)}{m_2(t)} = 1$. Then $\lim_{t \rightarrow \infty} \frac{\psi_g(m_1(t))}{\psi_g(m_2(t))} = 1$.

3. PROOFS OF THEOREMS 1.2 AND 1.6

In the sequel, we denote by

$$v_0(t) = \int_t^\infty \frac{1}{A(s)} ds \quad \text{for } t \in (0, \infty).$$

Since the function A satisfies (H1), then using Definition 2.5 and assertion (ii) of Lemma 2.6, we deduce that there exists $L_0 \in \mathcal{K}$ such that $A(t) = t^\alpha L_0(t)$, for $t > 1$. Hence, using Lemma 2.1, we deduce that $1/A$ is integrable near infinity. So the function v_0 is well defined, and by Lemma 2.2 we have

$$v_0(t) = \int_t^\infty \frac{1}{A(s)} ds \sim \frac{t^{1-\alpha}}{(\alpha-1)L_0(t)} \quad \text{as } t \rightarrow \infty. \quad (3.1)$$

In the sequel, we denote also by $L_A u := \frac{1}{A}(Au)' = u'' + \frac{A'}{A}u'$ and we remark that $L_A v_0 = 0$.

Proof of Theorem 1.2. Let $\varepsilon \in (0, a_1/2)$. Put

$$\xi_i = \frac{a_i}{(\alpha+1-\lambda) + (\lambda-2)C_g} \quad \text{for } i \in \{1, 2\},$$

$\tau_1 = \xi_1 - \varepsilon \frac{\xi_1}{a_1}$ and $\tau_2 = \xi_2 + \varepsilon \frac{\xi_2}{a_2}$. Clearly, we have $\frac{\xi_1}{2} < \tau_1 < \tau_2 < \frac{3}{2}\xi_2$. Let $\theta(t) = \int_t^\infty s^{1-\lambda} L(s) ds$ and put

$$\omega_i(t) = \psi_g\left(\tau_i \int_t^\infty s^{1-\lambda} L(s) ds\right) = \psi_g(\tau_i \theta(t)), \quad \text{for } t > 0.$$

By a simple calculus, for $i \in \{1, 2\}$ we obtain

$$\begin{aligned} & L_A \omega_i(t) + q(t)g(\omega_i(t)) \\ &= g(\omega_i(t))t^{-\lambda} L(t) \left[\tau_i(\tau_i t^{2-\lambda} L(t)g'(\omega_i(t)) + (\lambda-2)C_g) \right. \\ & \quad \left. - \tau_i \left(\frac{tA'(t)}{A(t)} - \alpha + \frac{tL'(t)}{L(t)} \right) - \tau_i((\alpha+1-\lambda) + (\lambda-2)C_g) + a_i \right. \\ & \quad \left. + \left(\frac{q(t)}{t^{-\lambda} L(t)} - a_i \right) \right]. \end{aligned}$$

So, for the fixed $\varepsilon > 0$, there exists $M_\varepsilon > 1$ such that for $t > M_\varepsilon$ and $i \in \{1, 2\}$, we have

$$\begin{aligned} \tau_i \left| \frac{tA'(t)}{A(t)} - \alpha + \frac{tL'(t)}{L(t)} \right| &\leq \frac{3}{2}\xi_2 \left(\left| \frac{tA'(t)}{A(t)} - \alpha \right| + \left| \frac{tL'(t)}{L(t)} \right| \right) \leq \frac{\varepsilon}{4}, \\ a_1 - \frac{\varepsilon}{2} &\leq \frac{a(t)}{t^{-\mu} L(t)} \leq a_2 + \frac{\varepsilon}{2} \end{aligned}$$

$$|\tau_i(\tau_i t^{2-\lambda} L(t)g'(\omega_i(t)) + (\lambda-2)C_g)| \leq \frac{3}{2}\xi_2 |\tau_i t^{2-\lambda} L(t)g'(\omega_i(t)) + (\lambda-2)C_g| \leq \frac{\varepsilon}{4}.$$

Indeed, the last inequality follows from (1.6) and the fact that from Lemmas 2.2 and 2.3, we have

$$\lim_{t \rightarrow \infty} \frac{t^{2-\lambda} L(t)}{\int_t^\infty s^{1-\lambda} L(s) ds} = 2 - \lambda,$$

for all $\lambda \geq 2$. This implies that for each $t > M_\varepsilon$, we have

$$L_A \omega_1(t) + q(t)g(\omega_1(t)) \geq g(\omega_1(t))t^{-\lambda} L(t) [-\varepsilon + a_1 - \tau_1((\alpha+1-\lambda) + (\lambda-2)C_g)] = 0$$

and

$$L_A \omega_2(t) + q(t)g(\omega_2(t)) \leq g(\omega_2(t))t^{-\lambda} L(t)[\varepsilon + a_2 - \tau_2((\alpha + 1 - \lambda) + (\lambda - 2)C_g)] = 0.$$

Let $u \in C^2((0, \infty)) \cap C([0, \infty))$ be the unique solution of (1.1) (see [17]). Then, there exists $B > 0$ such that

$$\omega_1(M_\varepsilon) - B v_0(M_\varepsilon) \leq u(M_\varepsilon) \leq \omega_2(M_\varepsilon) + B v_0(M_\varepsilon). \tag{3.2}$$

We claim that

$$\omega_1(t) - B v_0(t) \leq u(t) \leq \omega_2(t) + B v_0(t) \quad \text{for all } t > M_\varepsilon. \tag{3.3}$$

Assume for instance that the right inequality of (3.3) is not true. Then the function $h(t) = u(t) - \omega_2(t) - B v_0(t)$ for $t > M_\varepsilon$ is not negative. Consequently, there exists $t_1 > M_\varepsilon$ such that $h(t_1) = \max_{M_\varepsilon \leq t < \infty} h(t) > 0$. Since h is continuous on $[M_\varepsilon, \infty)$, $h(M_\varepsilon) \leq 0$ and $\lim_{t \rightarrow \infty} h(t) = 0$, then $h'(t_1) = 0$ and $h(t) > 0$ for $t \in (t_1 - \delta, t_1 + \delta)$ for some $\delta > 0$, sufficiently small. Namely $u(t) > \omega_2(t) + B v_0(t)$ for $t \in (t_1 - \delta, t_1 + \delta)$. Since g is non-increasing on $(0, \infty)$, then

$$\frac{1}{A(t)}(A(t)h'(t))' = -q(t)g(u(t)) - \frac{1}{A(t)}(A(t)\omega_2'(t))' \geq q(t)(g(\omega_2(t)) - g(u(t))) \geq 0,$$

for $t \in (t_1 - \delta, t_1 + \delta)$. Which implies $h'(t) \leq h'(t_1) = 0$ for $t \in (t_1 - \delta, t_1)$ and $h'(t) \geq h'(t_1) = 0$ for $t \in (t_1, t_1 + \delta)$. This implies that h has a local minimum at t_1 . Which contradicts the fact that h a global maximum at t_1 on $[M_\varepsilon, \infty)$. This proves that

$$u(t) \leq \omega_2(t) + B v_0(t) \quad \text{for all } t > M_\varepsilon.$$

Similarly, we show that

$$\omega_1(t) - B v_0(t) \leq u(t) \quad \text{for all } t > M_\varepsilon.$$

This proves (3.3).

Now, since $\psi_g \circ \theta \in NRV I_{(2-\lambda)(1-C_g)}$, there exists $\hat{L} \in \mathcal{K}$ such that $\psi_g \circ \theta = t^{(2-\lambda)(1-C_g)} \hat{L}(t)$ for $t \in [1, \infty)$. Moreover since $(\alpha - 1) - (\lambda - 2)(1 - C_g) > 0$, it follows by Lemma 2.1 that

$$\lim_{t \rightarrow \infty} \frac{t^{1-\alpha}}{t^{(2-\lambda)(1-C_g)} \hat{L}(t)} = 0.$$

This implies that

$$\lim_{t \rightarrow \infty} \frac{t^{1-\alpha}}{\psi_g(\tau_i \int_t^\infty s^{1-\lambda} L(s) ds)} = \lim_{t \rightarrow \infty} \frac{t^{1-\alpha}}{\psi_g(\tau_i \theta(t))} = \lim_{t \rightarrow \infty} \frac{\psi_g(\theta(t))}{\psi_g(\tau_i \theta(t))} \frac{t^{1-\alpha}}{\psi_g(\theta(t))} = 0$$

uniformly in $\tau_i \in [\frac{\xi_1}{2}, \frac{3}{2}\xi_2] \subset (0, \infty)$. This together with (3.1) implies

$$\lim_{t \rightarrow \infty} \frac{v_0(t)}{\psi_g(\tau_1 \theta(t))} = \lim_{t \rightarrow \infty} \frac{v_0(t)}{\psi_g(\tau_2 \theta(t))} = 0.$$

So, we obtain

$$\limsup_{t \rightarrow \infty} \frac{u(t)}{\omega_2(t)} \leq 1 \leq \liminf_{t \rightarrow \infty} \frac{u(t)}{\omega_1(t)}.$$

Using this fact and assertions (i) and (iv) of Lemma 2.6, we deduce that

$$\liminf_{t \rightarrow \infty} \frac{u(t)}{\psi_g(\theta(t))} = \liminf_{t \rightarrow \infty} \frac{u(t)}{\omega_1(t)} \frac{\omega_1(t)}{\psi_g(\theta(t))} \geq \lim_{t \rightarrow \infty} \frac{\psi_g(\tau_1 \theta(t))}{\psi_g(\theta(t))} = \tau_1^{1-C_g}.$$

By letting ε approach zero, we obtain

$$\xi_1^{1-C_g} \leq \liminf_{t \rightarrow \infty} \frac{u(t)}{\psi_g(\theta(t))}.$$

Similarly, we obtain

$$\limsup_{t \rightarrow \infty} \frac{u(t)}{\psi_g(\theta(t))} \leq \xi_2^{1-C_g}.$$

This proves in particular the exact behavior at infinity in the case $\lambda = 2$. Now, for $\lambda > 2$, we have by Lemma 2.2 that $\theta(t) \sim_{t \rightarrow \infty} \frac{t^{2-\lambda}}{\lambda-2} L(t)$. Hence it follows by assertions (i), (iv) and (v) of Lemma 2.6 that for $\lambda > 2$, we have

$$\lim_{t \rightarrow \infty} \frac{\psi_g(\theta(t))}{\psi_g((t)^{2-\lambda} L(t))} = \lim_{t \rightarrow \infty} \frac{\psi_g(\theta(t))}{\psi_g((\lambda-2)\theta(t))} \frac{\psi_g((\lambda-2)\theta(t))}{\psi_g((t)^{2-\lambda} L(t))} = \frac{1}{(\lambda-2)^{1-C_g}}.$$

This achieves the proof of the Theorem. \square

Proof of Theorem 1.6. We recall that $g(t) = t^{-\gamma}$, $\lambda = \alpha + 1 + (\alpha - 1)\gamma$ and $C_g = \frac{\gamma}{1+\gamma}$. Let $\varepsilon \in (0, \frac{\alpha_1}{2})$ and put $\tau_1 = (\gamma + 1)(a_1 - \varepsilon)$ and $\tau_2 = (\gamma + 1)(a_2 + \varepsilon)$. Put $k(t) = \int_1^t \frac{L(s)}{s} ds$ and

$$\omega_i(t) = \left((1 + \gamma)\tau_i \int_t^\infty s^{1-\lambda} k(s) ds \right)^{\frac{1}{1+\gamma}} \quad \text{for } i \in \{1, 2\},$$

where L is the function given in hypothesis (H_6) . Then, by a simple computation, we have

$$\begin{aligned} & L_A \omega_i(t) + q(t)g(\omega_i(t)) \\ &= g(\omega_i(t))t^{-\lambda} L(t) \left[\tau_i \left(\frac{k(t)}{L(t)} (\tau_i t^{2-\lambda} k(t) g'(\omega_i) + (\lambda - 1 - \alpha)) - \frac{\gamma}{\gamma + 1} \right) \right. \\ & \quad \left. - \tau_i \frac{k(t)}{L(t)} \left(\frac{t A'(t)}{A(t)} - \alpha \right) + \frac{\gamma}{\gamma + 1} \tau_i - \tau_i + a_i + \left(\frac{q(t)}{t^{-\lambda} L(t)} - a_i \right) \right] \\ &= g(\omega_i(t))t^{-\lambda} L(t) \left[\tau_i \left(\frac{k(t)}{L(t)} (\tau_i t^{2-\lambda} k(t) g'(\omega_i) + (\alpha - 1)\gamma) - \frac{\gamma}{\gamma + 1} \right) \right. \\ & \quad \left. - \tau_i \frac{k(t)}{L(t)} \frac{t B'(t)}{B(t)} - \frac{\tau_i}{\gamma + 1} + a_i + \left(\frac{q(t)}{t^{-\lambda} L(t)} - a_i \right) \right]. \end{aligned}$$

Since $g(t) = t^{-\gamma}$ and $\lambda = \alpha + 1 + (\alpha - 1)\gamma$, integrating by parts, we obtain

$$\begin{aligned} & \tau_i t^{2-\lambda} k(t) g'(\omega_i(t)) + (\alpha - 1)\gamma \\ &= -\gamma \tau_i t^{2-\lambda} k(t) (\omega_i(t))^{-(1+\gamma)} + (\alpha - 1)\gamma \\ &= \gamma \left((\alpha - 1) - \frac{t^{2-\alpha} k(t)}{(\gamma + 1) \int_t^\infty s^{1-\lambda} k(s) ds} \right) \\ &= \gamma \left(\frac{(\alpha - 1)(1 + \gamma) \int_t^\infty s^{1-\lambda} k(s) ds - t^{2-\lambda} k(t)}{(1 + \gamma) \int_t^\infty s^{1-\lambda} k(s) ds} \right) \\ &= \frac{\gamma}{\gamma + 1} \frac{\int_t^\infty s^{1-\lambda} L(s) ds}{\int_t^\infty s^{1-\lambda} k(s) ds}. \end{aligned}$$

This gives

$$\frac{k(t)}{L(t)} (\tau_i t^{2-\lambda} k(t) g'(\omega_i(t)) + (\alpha - 1)\gamma) - \frac{\gamma}{\gamma + 1}$$

$$= \frac{\gamma}{\gamma + 1} \left[\frac{\int_t^\infty s^{1-\lambda} L(s) ds}{t^{2-\lambda} L(t)} \frac{t^{2-\lambda} k(t)}{\int_t^\infty s^{1-\lambda} k(s) ds} - 1 \right].$$

This together with Lemma 2.2 and the fact that k and L are in \mathcal{K} , implies

$$\lim_{t \rightarrow \infty} \frac{k(t)}{L(t)} (\tau_i t^{2-\lambda} k(t) g'(\omega_i(t)) + (\alpha - 1)\gamma) - \frac{\gamma}{\gamma + 1} = 0.$$

Now since $\frac{t^\nu B'(t)}{B(t)}$ is bounded for t large and by Lemma 2.1, we have $\frac{k}{L} \in \mathcal{K}$ and $\lim_{t \rightarrow \infty} \frac{t^{1-\nu} k(t)}{L(t)} = 0$, we deduce that

$$\lim_{t \rightarrow \infty} \frac{k(t)}{L(t)} \left(\frac{t B'(t)}{B(t)} \right) = \lim_{t \rightarrow \infty} \frac{t^{1-\nu} k(t)}{L(t)} \left(\frac{t^\nu B'(t)}{B(t)} \right) = 0.$$

So, for the fixed $\varepsilon > 0$, there exists $M_\varepsilon > 1$ such that for $t \geq M_\varepsilon$, we have

$$\begin{aligned} L_A \omega_2(t) + q(t)g(\omega_2(t)) &\leq g(\omega_2(t))t^{-\lambda} L(t) \left[\frac{\varepsilon}{3} + \frac{\varepsilon}{3} - \frac{\tau_2}{\gamma + 1} + a_2 + \frac{\varepsilon}{3} \right] = 0, \\ L_A \omega_1(t) + q(t)g(\omega_1(t)) &\geq g(\omega_1(t))t^{-\mu} L(t) \left[-\frac{\varepsilon}{3} - \frac{\varepsilon}{3} - \frac{\tau_1}{\gamma + 1} + a_1 - \frac{\varepsilon}{3} \right] = 0. \end{aligned}$$

Let $u \in C([0, \infty)) \cap C^2((0, \infty))$ be the unique solution of (1.7). As in the proof of Theorem 1.2, we choose $C > 0$ such that

$$\omega_1(t) - C v_0(t) \leq u(t) \leq \omega_2(t) + C v_0(t) \quad \text{for } t \geq M_\varepsilon.$$

Moreover, thanks to (H6), we have $\lim_{t \rightarrow \infty} k(t) = \infty$. So, using Lemma 2.2, we obtain

$$\begin{aligned} \lim_{t \rightarrow \infty} \frac{1}{t^{\alpha-1} \left((1 + \gamma)\tau_1 \int_t^\infty s^{1-\lambda} k(s) ds \right)^{\frac{1}{1+\gamma}}} &= \lim_{t \rightarrow \infty} \frac{1}{t^{\alpha-1} \left(t^{(2-\lambda)} \frac{\tau_1 k(t)}{\alpha-1} \right)^{\frac{1}{1+\gamma}}} \\ &= \lim_{t \rightarrow \infty} \left(\frac{\alpha - 1}{\tau_1 k(t)} \right)^{\frac{1}{1+\gamma}} = 0. \end{aligned}$$

This and (3.1) gives $\lim_{t \rightarrow \infty} \frac{v_0(t)}{\omega_1(t)} = 0$. Similarly, we obtain $\lim_{t \rightarrow \infty} \frac{v_0(t)}{\omega_2(t)} = 0$. So we have

$$\limsup_{t \rightarrow \infty} \frac{u(t)}{\omega_2(t)} \leq 1 \leq \liminf_{t \rightarrow \infty} \frac{u(t)}{\omega_1(t)}.$$

This implies that

$$\liminf_{t \rightarrow \infty} \frac{u(t)}{\left((1 + \gamma) \int_t^\infty s^{1-\lambda} k(s) ds \right)^{\frac{1}{1+\gamma}}} \geq \tau_1^{\frac{1}{1+\gamma}}.$$

Now, as ε tends to zero, we obtain

$$\liminf_{t \rightarrow \infty} \frac{u(t)}{\left((1 + \gamma) \int_t^\infty s^{1-\lambda} k(s) ds \right)^{\frac{1}{1+\gamma}}} \geq ((\gamma + 1)a_1)^{\frac{1}{1+\gamma}}.$$

Similarly, we obtain

$$\limsup_{t \rightarrow \infty} \frac{u(t)}{\left((1 + \gamma) \int_t^\infty s^{1-\lambda} k(s) ds \right)^{\frac{1}{1+\gamma}}} \leq ((\gamma + 1)a_2)^{\frac{1}{1+\gamma}}.$$

Now, since $(\gamma+1) \int_t^\infty s^{1-\lambda} k(s) ds \sim_{t \rightarrow \infty} \frac{t^{2-\lambda} k(t)}{\alpha-1} = \frac{t^{(1-\alpha)(\gamma+1)}}{\alpha} \int_1^t \frac{L(s)}{s} ds$, we deduce that

$$\begin{aligned} \left(\frac{(\gamma+1)a_1}{\alpha-1} \right)^{\frac{1}{1+\gamma}} &\leq \liminf_{t \rightarrow \infty} \frac{u(t)}{t^{1-\alpha} \left(\int_1^t \frac{L(s)}{s} ds \right)^{\frac{1}{1+\gamma}}} \\ &\leq \limsup_{t \rightarrow \infty} \frac{u(t)}{t^{1-\alpha} \left(\int_1^t \frac{L(s)}{s} ds \right)^{\frac{1}{1+\gamma}}} \\ &\leq \left(\frac{(\gamma+1)a_2}{\alpha} \right)^{\frac{1}{1+\gamma}}. \end{aligned}$$

In particular, if $a_1 = a_2$, we obtain

$$\lim_{t \rightarrow \infty} \frac{u(t)}{t^{1-\alpha} \left(\int_1^t \frac{L(s)}{s} ds \right)^{\frac{1}{1+\gamma}}} = \left(\frac{(\gamma+1)a_1}{\alpha-1} \right)^{\frac{1}{1+\gamma}}.$$

□

4. APPLICATIONS

Application 1. We consider the Dirichlet problem

$$\begin{aligned} -\frac{1}{A}(Au')' + \frac{\beta}{u}(u')^2 &= q(t)g(u), \quad t \in (0, \infty), \\ u &> 0, \quad \text{in } (0, \infty), \\ \lim_{t \rightarrow 0^+} A(t)u'(t) &= 0 \quad \lim_{t \rightarrow \infty} u(t) = 0, \end{aligned} \tag{4.1}$$

where $\beta < 1$ and $\lim_{t \rightarrow \infty} \frac{q(t)}{t^{-\lambda} L(t)} = a_0 > 0$ with $\lambda \geq 2$ and $L \in \mathcal{K}$ with $\int_1^\infty s^{1-\lambda} L(s) ds < \infty$.

We assume that A satisfies (H1) and g satisfies the following hypotheses:

(A1) The function $t \rightarrow t^{-\beta} g(t)$ is non-increasing from $(0, \infty)$ into $(0, \infty)$.

(A2) $\lim_{t \rightarrow 0} g'(t) \int_0^t \frac{1}{g(s)} ds = -C_g$ with $\max(0, \frac{\beta}{\beta-1}) \leq C_g \leq 1$.

(A3) $(\alpha-1) - (\lambda-2)(1-\beta)(1-C_g) > 0$

Note that for $\gamma > 0$ and $-\gamma < \beta < 1$, the function $g(t) = t^{-\gamma}$ satisfies (A1) and (A2). Put $u = v^{\frac{1}{1-\beta}}$. Then v satisfies

$$\begin{aligned} -\frac{1}{A}(Av')' &= (1-\beta)q(t)g(v^{\frac{1}{1-\beta}})v^{\frac{-\beta}{1-\beta}}, \quad t \in (0, \infty), \\ v &> 0, \quad \text{in } (0, \infty), \\ \lim_{t \rightarrow 0^+} A(t)v'(t) &= 0, \quad \lim_{t \rightarrow \infty} v(t) = 0, \end{aligned} \tag{4.2}$$

The function $f(r) = (1-\beta)g(r^{\frac{1}{1-\beta}})r^{-\frac{\beta}{1-\beta}}$ is non-increasing on $(0, \infty)$ and a simple computation shows that $\psi_g = (\psi_f)^{\frac{1}{1-\beta}}$ and

$$\lim_{r \rightarrow 0} f'(r) \int_0^r \frac{1}{f(s)} ds = (1-\beta)(1-C_g) - 1 =: -C_f, \quad \text{with } 0 \leq C_f \leq 1.$$

Applying Corollary 1.3 to problem (4.2), we deduce that there exists a unique solution v to (4.2) such that

(a) When $C_f = 1$, we have

$$(i) \lim_{t \rightarrow \infty} \frac{v(t)}{\psi_f(t^{2-\lambda} L(t))} = 1 \text{ if } \lambda > 2;$$

$$(ii) \lim_{t \rightarrow \infty} \frac{v(t)}{\psi_f(\int_t^\infty \frac{L(s)}{s} ds)} = 1 \text{ if } \lambda = 2;$$

(b) When $C_f < 1$, we have:

$$(i) \lim_{t \rightarrow \infty} \frac{v(t)}{\psi_f(t^{2-\lambda} L(t))} = \left[\frac{a_0}{\alpha+1-\lambda+(\lambda-2)C_f} \right]^{1-C_f} \text{ if } 2 < \lambda < 2 + \frac{\alpha-1}{(1-\beta)(1-c_g)};$$

$$(ii) \lim_{t \rightarrow \infty} \frac{v(t)}{\psi_f(\int_t^\infty \frac{L(s)}{s} ds)} = \left[\frac{a_0}{\alpha-1} \right]^{1-C_f} \text{ if } \lambda = 2.$$

This implies that problem (4.1) has a solution $u \in C([0, \infty)) \cap C^2((0, \infty))$ satisfying the following exact behavior

(a) When $C_g = 1$, we have:

(i) if $\lambda > 2$, then

$$\lim_{t \rightarrow \infty} \frac{u(t)}{\psi_g(t^{2-\lambda} L(t))} = 1;$$

(ii) if $\lambda = 2$, then

$$\lim_{t \rightarrow \infty} \frac{u(t)}{\psi_g(\int_t^\infty \frac{L(s)}{s} ds)} = 1;$$

(b) If $\max(0, \frac{\beta}{\beta-1}) \leq C_g < 1$, then:

(i) if $2 < \lambda < 2 + \frac{\alpha-1}{(1-\beta)(1-C_g)}$, then

$$\lim_{t \rightarrow \infty} \frac{u(t)}{\psi_g(t^{2-\lambda} L(t))} = \left[\frac{a_0}{\alpha-1-(\lambda-2)(1-\beta)(1-C_g)} \right]^{1-C_g}$$

(ii) if $\lambda = 2$, then

$$\lim_{|x| \rightarrow \infty} \frac{u(t)}{\psi_g(\int_t^\infty \frac{L(s)}{s} ds)} = \left[\frac{a_0}{\alpha-1} \right]^{1-C_g}.$$

Application 2. In this subsection, we assume that the function A satisfy the following hypothesis

(A4) A is a continuous function on $[0, \infty)$, positive and differentiable on $(0, \infty)$ such that $\frac{1}{A}$ is integrable near 0 and $\lim_{t \rightarrow \infty} \frac{tA'(t)}{A(t)} = \sigma \in \mathbb{R} - \{1\}$.

We are interested in the exact behavior at infinity of the unique positive solution of the problem

$$\begin{aligned} \frac{1}{A(t)}(A(t)u'(t))' &= -p(t)u^{-\gamma}, \quad t \in (0, \infty), \\ u &> 0, \quad \text{in } (0, \infty), \\ u(0) &= 0, \quad \lim_{t \rightarrow \infty} \frac{u(t)}{\rho(t)} = 0, \end{aligned} \tag{4.3}$$

where $\gamma > 0$ and $\rho(t) = \int_0^t \frac{ds}{A(s)}$. Let $u(t) = \rho(t)v(t)$ and $B(t) = A(t)\rho^2(t)$ for $t \in [0, \infty)$. Then u is a positive solution of (4.3) if and only if v is a positive solution of the problem

$$\begin{aligned} \frac{1}{B(t)}(B(t)v'(t))' &= -\frac{p(t)}{(\rho(t))^{\gamma+1}} v^{-\gamma}, \quad t \in (0, \infty), \\ v &> 0, \quad \text{in } (0, \infty), \\ \lim_{t \rightarrow 0^+} B(t)v'(t) &= 0, \quad \lim_{t \rightarrow \infty} v(t) = 0. \end{aligned} \tag{4.4}$$

First, we claim that if A satisfies (A4), then

$$\lim_{t \rightarrow \infty} \frac{tB'(t)}{B(t)} = 1 + |\sigma - 1| > 1. \quad (4.5)$$

Since $\frac{tB'(t)}{B(t)} = \frac{tA'(t)}{A(t)} + \frac{2t}{A(t)\rho(t)}$ and by Definition 2.5 and assertion (ii) of Lemma 2.6, we have $A(t) = t^\sigma L_0(t)$ for $t \geq a > 1$ with $L_0 \in \mathcal{K}$, then we deduce from Lemma 2.2 that

- For $\sigma < 1$, we have $\rho(\infty) = \infty$ and so

$$\rho(t) = \int_0^a \frac{ds}{A(s)} + \int_a^t \frac{1}{s^\sigma L_0(s)} ds \sim \frac{1}{1-\sigma} \frac{t^{1-\sigma}}{L_0(t)} \quad \text{as } t \rightarrow \infty.$$

So

$$\frac{2t}{A(t)\rho(t)} \sim (1-\sigma) \frac{L_0(t)}{t^{1-\sigma}} \frac{2t}{t^\sigma L_0(t)} = 2(1-\sigma) \quad \text{as } t \rightarrow \infty.$$

Consequently in this case we have

$$\lim_{t \rightarrow \infty} \frac{tB'(t)}{B(t)} = \sigma + 2(1-\sigma) = 2 - \sigma = 1 + |\sigma - 1|.$$

- For $\sigma > 1$, we have $\rho(\infty) = \int_0^\infty \frac{ds}{A(s)} ds < \infty$. So

$$\frac{2t}{A(t)\rho(t)} \sim \frac{2t}{t^\sigma L_0(t)\rho(\infty)} \rightarrow 0 \quad \text{as } t \rightarrow \infty.$$

In this case we have

$$\lim_{t \rightarrow \infty} \frac{tB'(t)}{B(t)} = \sigma = 1 + |\sigma - 1|.$$

This proves (4.5). Taking into account this fact, we assume that the function p satisfies the following hypotheses

- (A5) p is a nonnegative continuous function $(0, \infty)$ satisfying

$$0 < a_0 = \lim_{t \rightarrow \infty} \frac{t^\lambda p(t)}{L(t)(\rho(t))^{\gamma+1}} < \infty,$$

where $\lambda \geq 2$ and $L \in \mathcal{K}$ such that $\int_1^\infty s^{1-\lambda} L(s) ds < \infty$.

- (A6) $2 + |\sigma - 1| - \lambda + (\lambda - 2) \frac{\gamma}{\gamma+1} > 0$.

Assume that A and p satisfy (A4)–(A6) and let v be the unique positive solution of problem (4.4). Then v has the following exact behavior at infinity

- (i) if $\lambda > 2$, then

$$\lim_{t \rightarrow \infty} \frac{v(t)}{[(\gamma+1)t^{2-\lambda}L(t)]^{\frac{1}{1+\gamma}}} = \left[\frac{a_0}{(\lambda-2)(2+|\sigma-1|-\lambda+(\lambda-2)\frac{\gamma}{\gamma+1})} \right]^{\frac{1}{1+\gamma}}.$$

- (ii) if $\lambda = 2$, then

$$\lim_{t \rightarrow \infty} \frac{v(t)}{[(\gamma+1) \int_t^\infty \frac{L(s)}{s} ds]^{\frac{1}{1+\gamma}}} = \left[\frac{a_0}{|\sigma-1|} \right]^{\frac{1}{1+\gamma}}$$

Consequently, the unique positive solution u of problem (4.3) has the following exact behavior at infinity

- (i) if $\lambda > 2$, then

$$\lim_{t \rightarrow \infty} \frac{u(t)}{\rho(t)[(\gamma+1)t^{2-\lambda}L(t)]^{\frac{1}{1+\gamma}}} = \left[\frac{a_0}{(\lambda-2)(2+|\sigma-1|-\lambda+(\lambda-2)\frac{\gamma}{\gamma+1})} \right]^{\frac{1}{1+\gamma}}.$$

ii) if $\lambda = 2$, then

$$\lim_{t \rightarrow \infty} \frac{u(t)}{\rho(t)[(\gamma + 1) \int_t^\infty \frac{L(s)}{s} ds]^{\frac{1}{1+\gamma}}} = \left[\frac{a_0}{|\sigma - 1|} \right]^{\frac{1}{1+\gamma}}$$

Acknowledgements. This project was funded by the Deanship of Scientific Research (DSR), King Abdulaziz University, Jeddah, under grant No. (253-662-1436-G). The authors, therefore, acknowledge with thanks DSR technical and financial support.

The three authors have contributed equally for this article.

REFERENCES

- [1] R. P. Agarwal, D. O'Regan; Singular problems modelling phenomena in the theory of pseudoplastic fluids, *ANZIAM J.* **45** (2003), 167-179.
- [2] R. Alsaedi, H. Mâagli, N. Zeddini; Exact behavior of the unique positive solution to some singular elliptic problem in exterior domains, *Nonlinear Anal.* **119** (2015), 186-198.
- [3] V. Avakumovic; Sur l'équation différentielle de Thomas-Fermi, *Acad. Serbe Sci., Publ. Inst. Math.* **1** (1947), 101-113.
- [4] I. Bachar, H. Mâagli, N. Zeddini; Estimates on the Green function and existence of positive solutions of nonlinear singular elliptic equations, *Comm. Contemporary Math.* **5** (2003), 401-434.
- [5] S. Barile, A. Salvatore; Existence and multiplicity results for some Lane-Emden elliptic systems: subquadratic case, *Adv. Nonlinear Anal.* **4** (2015), 25-35.
- [6] A. Callegari, A. Nachman; Some singular, nonlinear differential equations arising in boundary layer theory, *J. Math. Anal. Appl.* **64** (1978), 96-105.
- [7] M. Cencelj, D. Repovš, Z. Virk; Multiple perturbations of a singular eigenvalue problem, *Nonlinear Anal.* **119** (2015), 37-45.
- [8] R. Chemmam, A. Dhifli, H. Mâagli; Asymptotic behavior of ground state solutions for sublinear and singular nonlinear Dirichlet problems, *Electronic Journal Differential Equations*, Vol. 2011 (2011), No. 88, pp. 1-12.
- [9] J. S. Fulks, J. S. Maybee; A singular nonlinear elliptic equation, *Osaka J. Math.* **12** (1960), 1-19.
- [10] J. A. Gatica, G. E. Hernandez, P. Waltman; Radially symmetric positive solutions for a class of singular elliptic equations, *Proceedings of the Edinburgh Mathematical Society* **33** (1990), 169-180.
- [11] J. A. Gatica, V. Olikar, P. Waltman; Singular nonlinear boundary value problems for second order ordinary differential equations, *J. Differential Equations* **79** (1989), 62-78.
- [12] P. G. de Gennes; Wetting: statics and dynamics, *Review of Modern Physics* **57** (1985), 827-863.
- [13] M. Ghergu, V. D. Rădulescu; Bifurcation and asymptotics for the Lane-Fowler equations, *C. R. Acad. Sci. Paris Ser. I* **337** (2003), 259-264.
- [14] M. Ghergu, V. D. Rădulescu; Multi-parameter bifurcation and asymptotics for the singular Lane-Emden-Fowler equation with a convection term, *Proc. Roy. Soc. Edinburgh Sect. A* **135** (2005), 61-83.
- [15] M. Ghergu, V.D. Rădulescu; *Singular Elliptic Problems: Bifurcation and Asymptotic Analysis*, Oxford University Press, 2008.
- [16] M. Ghergu, V.D. Rădulescu; *Nonlinear PDEs: Mathematical Models in Biology, Chemistry and Population Genetics*, Springer Monographs in Mathematics, Springer-Verlag, Heidelberg, 2012.
- [17] H. Mâagli, S. Masmoudi; Sur les solutions d'un operateur différentiel singulier semi-linéaire, *Potential Analysis* **10** (1999), 289-304.
- [18] H. Mâagli, N. Mhadhebi, N. Zeddini; Existence and exact asymptotic behavior of positive solutions for a fractional boundary value problem, *Abstract and Applied Analysis*, Volume 2013, Article ID 420514, 6 p.
- [19] V. Maric; *Regular Variation and Differential Equations*, Lecture Notes in Math., Vol. 1726, Springer-Verlag, Berlin, 2000.

- [20] V. Maric, M. Tomic; Asymptotic properties of solutions of the equation $y'' = f(x)\phi(y)$, *Math. Z.* **149** (1976), 261-266.
- [21] V. Maric, M. Tomic; Regular variation and asymptotic properties of solutions of nonlinear differential equations, *Publ. Inst. Math. (Belgr.)* **21** (1977), 119-129.
- [22] A. Nachman, A. Callegari; A nonlinear singular boundary value problem in the theory of pseudoplastic fluids, *SIAM J. Appl. Math.* **38** (1980), 275-281.
- [23] D. O'Regan; Singular nonlinear differential equations on the half line, *Topological Methods in Nonlinear Analysis* **8** (1996), 137-159.
- [24] V. Rădulescu, D. Repovš; Perturbation effects in nonlinear eigenvalue problems, *Nonlinear Anal.* **70** (2009), 3030-3038.
- [25] V. Rădulescu, D. Repovš; *Partial Differential Equations with Variable Exponents: Variational Methods and Qualitative Analysis*, Chapman and Hall/CRC, Taylor & Francis Group, Boca Raton, FL, 2015.
- [26] D. Repovš; Singular solutions of perturbed logistic-type equations, *Appl. Math. Comp.* **218** (2011), 4414-4422.
- [27] D. Repovš; Asymptotics for singular solutions of quasilinear elliptic equations with an absorption term, *J. Math. Anal. Appl.* **395** (2012), 78-85.
- [28] R. Seneta; *Regular Varying Functions*, Lectures Notes in Math., Vol. 508, Springer-Verlag, Berlin, 1976.
- [29] S. Taliaferro; A nonlinear singular boundary value problem, *Nonlinear Anal.* **3** (1979), 897-904.
- [30] H. Usami; On a singular elliptic boundary value problem in a ball, *Nonlinear Anal.* **13** (1989), 1163-1170.
- [31] Z. Zhao; Positive solutions of nonlinear second order ordinary differential equations, *Proc. Amer. Math. Soc.* **121** (1994), 465-469.

HABIB MÁAGLI

DEPARTMENT OF MATHEMATICS, COLLEGE OF SCIENCES AND ARTS, KING ABDULAZIZ UNIVERSITY,
RABIGH CAMPUS, P.O. BOX 344, RABIGH 21911, SAUDI ARABIA

E-mail address: `habib.maagli@fst.rnu.tn`

RAMZI ALSAEDI

DEPARTMENT OF MATHEMATICS, FACULTY OF SCIENCES, KING ABDULAZIZ UNIVERSITY, P.O. BOX
80203, JEDDAH 21589, SAUDI ARABIA

E-mail address: `ramzialsaedi@yahoo.co.uk`

NOUREDDINE ZEDDINI

DEPARTMENT OF MATHEMATICS, COLLEGE OF SCIENCES AND ARTS, KING ABDULAZIZ UNIVERSITY,
RABIGH CAMPUS, P.O. BOX 344, RABIGH 21911, SAUDI ARABIA

E-mail address: `noureddine.zeddini@ipein.rnu.tn`