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# SOLUTIONS TO KIRCHHOFF EQUATIONS WITH COMBINED NONLINEARITIES 

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$$
\begin{aligned}
& \text { AbSTRACT. We prove the existence of multiple positive solutions for the Kirch- } \\
& \text { hoff equation } \\
& \qquad-\left(a+b \int_{\Omega}|\nabla u|^{2} d x\right) \Delta u=h(x) u^{q}+f(x, u), \quad x \in \Omega, \\
& \qquad u=0, \quad x \in \partial \Omega, \\
& \text { Here } \Omega \text { is an open bounded domain in } R^{N}(N=1,2,3), h(x) \in L^{\infty}(\Omega), \\
& f(x, s) \text { is a continuous function which is asymptotically linear at zero and is } \\
& \text { asymptotically 3-linear at infinity. Our main tools are the Ekeland's variational } \\
& \text { principle and the mountain pass lemma. }
\end{aligned}
$$

## 1. Introduction and main results

In this article, we study the existence of positive solutions for the Kirchhoff equation

$$
\begin{gather*}
-\left(a+b \int_{\Omega}|\nabla u|^{2} d x\right) \Delta u=h(x) u^{q}+f(x, u), \quad x \in \Omega  \tag{1.1}\\
u=0, \quad x \in \partial \Omega
\end{gather*}
$$

where $\Omega$ is a bounded smooth domain in $R^{N}(N=1,2,3), a>0, b>0,0<q<1$.
To state the assumptions, we recall some results about the following two eigenvalue problems:

$$
\begin{equation*}
-\Delta u=\lambda u \text { in } \Omega, \quad u=0 \text { on } \Omega \tag{1.2}
\end{equation*}
$$

and

$$
\begin{equation*}
-\left(\int_{\Omega}|\nabla u|^{2} d x\right) \Delta u=\mu u^{3} \text { in } \Omega, \quad u=0 \text { on } \Omega . \tag{1.3}
\end{equation*}
$$

Let $\lambda_{1}$ be the principal eigenvalue of 1.2 and let $\phi_{1}>0$ be its associated eigenfunction. It is known that $\lambda_{1}$ can be characterized by

$$
\lambda_{1}=\inf \left\{\int_{\Omega}|\nabla u|^{2} d x: u \in H_{0}^{1}(\Omega), \int_{\Omega}|u|^{2} d x=1\right\}
$$

[^0]where $H_{0}^{1}(\Omega)$ is the usual Sobolev space defined as the completion of $C_{0}^{\infty}(\Omega)$ with respect to the norm $\|u\|=\left(\int_{\Omega}|\nabla u|^{2} d x\right)^{1 / 2}$. Moreover, define
$$
\mu_{1}=\inf \left\{\|u\|^{4}: u \in H_{0}^{1}(\Omega), \int_{\Omega}|u|^{4} d x=1\right\}
$$

As shown in [13], there exists $\mu_{1}>0$ which is the principle eigenvalue of 1.3 ) and there is a corresponding eigenfunction of $\varphi_{1}>0$ in $\Omega$.

In this article, we assume that $h, f$ satisfy the following conditions:
(H1) $h \in L^{\infty}(\Omega)$ and $h(x) \not \equiv 0$;
(F1) $f \in C(\Omega \times \mathbb{R}), f(x, 0)=0$ for all $x \in \Omega, f(x, s) \geq 0$ for all $x \in \Omega$ and $s \geq 0$;

$$
\begin{equation*}
\lim _{s \rightarrow 0^{+}} \frac{f(x, s)}{a \lambda_{1} s+b \mu_{1} s^{3}}=\alpha \in[0,1), \quad \lim _{s \rightarrow+\infty} \frac{f(x, s)}{a \lambda_{1} s+b \mu_{1} s^{3}}=\beta \in(1,+\infty) \tag{F2}
\end{equation*}
$$

uniformly for a.e. $x \in \Omega$.
It is obvious that the values of $f(x, s)$ for $s<0$ are irrelevant for us to seek for positive solutions of 1.1 , and we may define

$$
f(x, s)=0 \quad \text { for } x \in \Omega, s \leq 0
$$

The problem

$$
\begin{align*}
-\left(a+b \int_{\Omega}|\nabla u|^{2} d x\right) \Delta u & =g(x, u), \quad x \in \Omega  \tag{1.4}\\
u=0, \quad x & \in \partial \Omega
\end{align*}
$$

is related to the stationary analogue of the Kirchhoff equation which was proposed by Kirchhoff in 1883 [9] as an generalization of the well-known d'Alembert's equation

$$
\rho \frac{\partial^{2} u}{\partial t^{2}}-\left(\frac{P_{0}}{h}+\frac{E}{2 L} \int_{0}^{L}\left|\frac{\partial u}{\partial x}\right|^{2} d x\right) \frac{\partial^{2} u}{\partial x^{2}}=g(x, u)
$$

for free vibrations of elastic strings. Kirchhoffs model takes into account the changes in length of the string produced by transverse vibrations. Here, $L$ is the length of the string, $h$ is the area of the cross section, $E$ is the Young modulus of the material, $\rho$ is the mass density and $P_{0}$ is the initial tension. In [1], it was pointed out that the problem 1.4 models several physical systems, where $u$ describes a process which depends on the average of itself. Nonlocal effect also finds its applications in biological systems. After [2] and [14], there are abundant results about Kirchhoff's equations.

Some interesting studies by variational methods can be found in 4, 12, 13, 19, 18, 17, 3, 16, 15] references therein and for Kirchhoff-type problem (1.4), they consider it in a bounded domain $\Omega$. For example, Perera and Zhang 13 obtain nontrivial solutions of (1.4) with asymptotically 4 -linear terms by using Yang index. In [19], they revisit problem (1.4) and establish the existence of a positive, a negative and a sign-changing solution by means of invariant sets of descent flow. Similar results can also be found in Mao and Zhang [12] and in Yang and Zhang [18]. Yang and Zhang in [17] obtain the existence of nontrivial solutions for 1.4 by using the local linking theory. Sun and Tang [16] prove the existence of a mountain pass type positive solution for problem (1.4) with the nonlinearity which is asymptotically linear near zero and superlinear at infinity. Sun and Liu [15] obtain a nontrivial solution via Morse theory by computing the relevant critical groups for problem (1.4) with the nonlinearity which is superlinear near zero but asymptotically 4 -linear
at infinity and asymptotically near zero but 4-linear at infinity. In [11, the authors obtain the existence of positive solutions for (1.1) with $h \equiv 0$ and $f(x, t)=\nu h(x, t)$ by using the topological degree argument and variational method, where $h$ is a continuous function which is asymptotically linear at zero and is asymptotically 3 linear at infinity. Inspired by [11], we shall study the existence of positive solutions for problem (1.1) with $h \not \equiv 0$ and $f$ which is asymptotically linear at zero and asymptotically 3 -linear infinity by using the Ekeland's variational principle and Mountain Pass Lemma different from [11]. In [11], when $N=1,2,3$, the authors studied equation 1.1 with $h \equiv 0$ and obtain the existence results of positive solution for equation (1.1) under the conditions: $a, b>0$, and $f$ satisfies (F1) and (F2) with $\alpha>1$ and $\beta<1 ; a \geq 0, b>0$, and $f$ satisfies (F1) and (F2) with $\alpha<1$ and $\beta>1$, respectively. But equation 1.1 with $h \not \equiv 0$ has not been studied. We shall obtain the existence of two positive solution for equation (1.1) because of the nonlinearity term $h(x) t^{q}(0<q<1)$. By the way, recently, Cheng, Wu and Liu [5] apply variant mountain pass theorem and Ekeland variational principle to study the existence of multiple nontrivial solutions for a class of Kirchhoff type problems with concave nonlinearity similar to our problem. But in their article, the nonlinear term is superlinear at infinity.

In this article, we denote by $\|\cdot\|_{p}$ the $L^{p}(\Omega)$-normal $(1 \leq p \leq \infty)$. We say that $u \in H_{0}^{1}(\Omega)$ is a positive (nonnegative) weak solution to problem (1.1) if $u>0$ $(u \geq 0)$ a.e. $\Omega$ and satisfies

$$
\left(a+b \int_{\Omega}|\nabla u|^{2} d x\right) \int_{\Omega} \nabla u \cdot \nabla v d x=\int_{\Omega} h(x) u^{q} v d x+\int_{\Omega} f(x, u) v d x
$$

for all $v \in H_{0}^{1}(\Omega)$. By assumption (F1), we know that to seek a nonnegative weak solution of 1.1 is equivalent to finding a nonzero critical point of the following functional on $H_{0}^{1}(\Omega)$ :
$I(u)=\frac{1}{2} \int_{\Omega}|\nabla u|^{2} d x+\frac{b}{4}\left(\int_{\Omega}|\nabla u|^{2} d x\right)^{2}-\frac{1}{q+1} \int_{\Omega} h(x)\left(u^{+}\right)^{q+1} d x-\int_{\Omega} F\left(x, u^{+}\right) d x$,
where $u^{+}=\max \{0, u\}, F(x, s)=\int_{0}^{s} f(x, \sigma) d \sigma$. By (F1) and (F2), I is a $C^{1}$ functional. By the strong maximum principle, the nonzero critical points of $I$ are positive solutions to problem (1.1) if $h(x) \geq 0$.

Our results are as follows.
Theorem 1.1. Suppose that $N=1,2,3, a>0, b>0,0<q<1, h$ and $f$ satisfy (H1), (F1), (F2). Assume further that exists $v \in H_{0}^{1}(\Omega)$ such that
(H2) $\int_{\Omega} h(x)\left(v^{+}\right)^{q+1} d x>0$.
Then there exists a constant $m>0$ such that if $\|h\|_{\infty}<m$, problem (1.1) has a solution $u_{1} \in H_{0}^{1}(\Omega), u_{1} \geq 0$ and $I\left(u_{1}\right)<0$. Moreover, if $h(x) \geq 0$, then $u_{1}>0$. e. in $\Omega$.

Theorem 1.2. Suppose that $N=1,2,3, a>0, b>0,0<q<1, h$ and $f$ satisfy (H1), (F1), (F2). Assume further $\beta \mu_{1}$ is not an eigenvalue of 1.3). Then there exists a constant $m>0$ such that if $\|h\|_{\infty}<m$, problem 1.1) has a nonnegative solution $u_{2} \in H_{0}^{1}(\Omega)$ with $u_{2}>0$ and $I\left(u_{2}\right)>0$ if $h(x) \geq 0$.

Remark 1.3. Theorem 1.1 for problem (1.1) with $a, b>0$ generalizes [10, Theorem 1.1] where (1.1) with $a=1$ and $b=0$.

Corollary 1.4. Suppose that $N=1,2,3, a>0, b>0,0<q<1, h$ and $f$ satisfy (H1), (F1), (F2). Assume further that $\beta \mu_{1}$ is not an eigenvalue of (1.3) and $h(x) \geq(\not \equiv) 0$. Then there exists a constant $m>0$ such that for all $h \in L^{\infty}(\Omega)$ with $\|h\|_{\infty}<m$, problem (1.1) has at least two positive solutions $u_{1}, u_{2} \in H_{0}^{1}(\Omega)$ such that $I\left(u_{1}\right)<0<I\left(u_{2}\right)$.

Remark 1.5. If $h(x) \geq(\not \equiv) 0$, it is easy to see that (H2) is always satisfied. Therefore, Corollary 1.1 is a straightforward conclusion of Theorems 1.1 and 1.2 by applying the strong maximum principle 8].

This paper is organized as follows. In Section 1, we obtain the existence of a local minimum solution by the Ekeland's variational principle. In Section 2, by using the Mountain Pass Lemma, we obtain the existence of a mountain pass solution. In the following discussion, we denote various positive constants as $C$ or $C_{i}, i=1,2,3, \ldots$.

## 2. Existence of a local minimum

In this section, we prove Theorem 1.1 by Ekeland's variational principle. We need the following Lemmas.

Lemma 2.1. Suppose that $N=1,2,3, a>0, b>0,0<q<1, h$ and $f$ satisfy (H1), (F1), (F2). Then there exists a constant $m>0$ such that if $\|h\|_{\infty}<m$, we have
(a) There exist $\rho, \gamma>0$ such that $\left.I(u)\right|_{\|u\|=\rho} \geq \gamma>0$.
(b) There exists an $e \in \mathbb{R} \backslash B_{\rho}(0)$ such that $I(e)<0$.

Proof. (a) By (F2), $\beta \in(1,+\infty)$ and noticing that $f(x, s) / s^{p-1} \rightarrow 0$ as $s \rightarrow+\infty$ uniformly in $x \in \Omega$ for any fixed $p \in(4,6)$ if $N=3 ; p \in(4,+\infty)$ if $N=1,2$. Given $\varepsilon \in(0,1)$, there exist $\delta, M_{\varepsilon}>0$ satisfying $0<\delta<+\infty$ such that

$$
f(x, s)<(\alpha+\varepsilon)\left(a \lambda_{1} s+b \mu_{1} s^{3}\right), \quad 0<s<\delta
$$

and

$$
f(x, s)<M_{\varepsilon} s^{p-1}, \quad \delta<s
$$

where $p \in(4,6)$ if $N=3 ; p \in(4,+\infty)$ if $N=1,2$. Together with (F1) and $f(x, s)=0$ for $x \in \Omega, s \leq 0$, we obtain

$$
f(x, s)<a \lambda_{1}(\alpha+\varepsilon)|s|+b \mu_{1}(\alpha+\varepsilon)|s|^{3}+M_{\varepsilon} s^{p-1}, \quad s \in R .
$$

This yields

$$
\begin{equation*}
F(x, s) \leq \frac{a \lambda_{1}}{2}(\alpha+\varepsilon)|s|^{2}+\frac{b \mu_{1}}{4}(\alpha+\varepsilon)|s|^{4}+A|s|^{p}, \quad s \in R \tag{2.1}
\end{equation*}
$$

where $A=M_{\varepsilon} / p$. Furthermore, by (F2), for the above $\varepsilon$, we have

$$
f(x, s)>(\beta-\varepsilon)\left(a \lambda_{1} s+b \mu_{1} s^{3}\right), \quad s>\delta_{\infty}
$$

Thus, we obtain

$$
F(x, s)>(\beta-\varepsilon)\left(\frac{a \lambda_{1}}{2} s^{2}+\frac{b \mu_{1}}{4} s^{4}\right), \quad s>\delta_{\infty}
$$

Together with (F1) and $f(x, s)=0$ for $x \in \Omega, s \leq 0$, there exists a constant $B>0$ such that

$$
\begin{equation*}
F(x, s) \geq \frac{a}{2}(\beta-\varepsilon) \lambda_{1}|s|^{2}+\frac{b}{4}(\beta-\varepsilon) \mu_{1}|s|^{4}-B, \quad s \in R . \tag{2.2}
\end{equation*}
$$

Since $\alpha<1$, we can choose $\varepsilon>0$ such that $\varepsilon<1-\alpha$. By (H1), 2.1), $\lambda_{1}\|u\|_{2}^{2} \leq$ $\|u\|^{2}, \mu_{1}\|u\|_{4}^{4} \leq\|u\|^{2}$, the Sobolev's embedding theorem: $\|u\|_{q+1}^{q+1} \leq K\|u\|^{q+1}$, $\|u\|_{p+1}^{p+1} \leq M\|u\|^{p+1}$ and the Young inequality, we have

$$
\begin{align*}
& I(u) \\
&= \frac{a}{2} \int_{\Omega}|\nabla u|^{2} d x+\frac{b}{4}\left(\int_{\Omega}|\nabla u|^{2} d x\right)^{2}-\frac{1}{q+1} \int_{\Omega} h(x)\left(u^{+}\right)^{q+1} d x-\int_{\Omega} F\left(x, u^{+}\right) d x \\
& \geq \frac{a}{2}\|u\|^{2}+\frac{b}{4}\|u\|^{4}-\frac{\|h\|_{\infty}}{q+1}\left\|u^{+}\right\|_{q+1}^{q+1}-\frac{a}{2}(\alpha+\varepsilon) \lambda_{1}\left\|u^{+}\right\|_{2}^{2} \\
&-\frac{b}{4}(\alpha+\varepsilon) \mu_{1}\left\|u^{+}\right\|_{4}^{4}-A\left\|u^{+}\right\|_{p}^{p} \\
& \geq \frac{a}{2}\|u\|^{2}+\frac{b}{4}\|u\|^{4}-\frac{\|h\|_{\infty}}{q+1}\|u\|_{q+1}^{q+1}-\frac{a}{2}(\alpha+\varepsilon)\|u\|^{2}-\frac{b}{4}(\alpha+\varepsilon)\|u\|^{4}-A\|u\|_{p}^{p} \\
& \geq \frac{a[1-(\alpha+\varepsilon)]}{2}\|u\|^{2}+\frac{b[1-(\alpha+\varepsilon)]}{4}\|u\|^{4}-\frac{\|h\|_{\infty} K}{q+1}\|u\|^{q+1}-A M\|u\|^{p} \\
& \geq\|u\|^{2}\left(C_{1}-C_{2}\|h\|_{\infty}\|u\|^{q-1}-C_{3}\|u\|^{p-2}\right), \tag{2.3}
\end{align*}
$$

where $C_{1}=\frac{a[1-(\alpha+\varepsilon)]}{2}, C_{2}=\frac{K}{q+1}$ and $C_{3}=A M$. Let

$$
g(t)=C_{2}\|h\|_{\infty} t^{q-1}+C_{3} t^{p-2} \quad \text { for } t \geq 0
$$

Clearly,

$$
g^{\prime}(t)=C_{2}(q-1)\|h\|_{\infty} t^{q-2}+(p-2) C_{3} t^{p-3}
$$

From $g^{\prime}\left(t_{0}\right)=0$, we have

$$
t_{0}=\left(C_{4}\|h\|_{\infty}\right)^{\frac{1}{p-q-1}}, \quad 0<q<1<4<p
$$

where $C_{4}=\frac{C_{2}(1-q)}{(p-2) C_{3}}$. Then

$$
g\left(t_{0}\right)=C_{2}\|h\|_{\infty}\left(C_{4}\|h\|_{\infty}\right)^{\frac{q-1}{p-q-1}}+C_{3}\left(C_{4}\|h\|_{\infty}\right)^{\frac{p-2}{p-q-1}}=C_{5}\|h\|_{\infty}^{\frac{p-2}{p-q-1}}
$$

where $C_{5}=C_{2} C_{4}^{\frac{q-1}{p-q-1}}+C_{3} C_{4}^{\frac{p-2}{p-q-1}}$ and $\frac{p-2}{p-q-1}>0$ because $0<q<1<4<p$. Thus, for any $p>4$, there exists $m>0$ such that $g\left(t_{0}\right)<C_{1}$ if $\|h\|_{\infty}<m$. Then, if $\|h\|_{\infty}<m$ and taking $\rho=t_{0}$, from (2.3), (a) is proved.
(b) For $t>0$ large enough, by 2.2 and $0<q<1$, taking $\varepsilon>0$ such that $\varepsilon<\min \{\beta-1,1-\alpha\}$, we have

$$
\begin{aligned}
I\left(t \varphi_{1}\right)= & \frac{a t^{2}}{2} \int_{\Omega}\left|\nabla \varphi_{1}\right|^{2} d x+\frac{b t^{4}}{4}\left(\int_{\Omega}\left|\nabla \varphi_{1}\right|^{2} d x\right)^{2}-\frac{t^{q+1}}{q+1} \int_{\Omega} h(x) \varphi_{1}^{q+1} d x \\
& -\int_{\Omega} F\left(x, t \varphi_{1}\right) d x \\
\leq & \frac{a t^{2}}{2}\left\|\varphi_{1}\right\|^{2}+\frac{b t^{4}}{4}\left\|\varphi_{1}\right\|^{4}-\frac{t^{q+1}}{q+1} \int_{\Omega} h(x) \varphi_{1}^{q+1} d x-\frac{a t^{2}}{2}(\beta-\varepsilon) \lambda_{1}\left\|\varphi_{1}\right\|_{2}^{2} \\
& -\frac{b t^{4}}{4}(\beta-\varepsilon) \mu_{1}\left\|\varphi_{1}\right\|_{4}^{4}+B|\Omega| \\
\leq & \frac{a t^{2}}{2}\left\|\varphi_{1}\right\|^{2}+\frac{b t^{4}}{4}\left\|\varphi_{1}\right\|^{4}-\frac{t^{q+1}}{q+1} \int_{\Omega} h(x) \varphi_{1}^{q+1} d x-\frac{b t^{4}}{4}(\beta-\varepsilon)\left\|\varphi_{1}\right\|^{4}+B|\Omega|
\end{aligned}
$$

$$
\begin{aligned}
& =\frac{a t^{2}}{2}\left\|\varphi_{1}\right\|^{2}-\frac{b t^{4}}{4}(\beta-\varepsilon-1)\left\|\varphi_{1}\right\|^{4}-\frac{t^{q+1}}{q+1} \int_{\Omega} h(x) \varphi_{1}^{q+1} d x+B|\Omega| \\
& \rightarrow-\infty
\end{aligned}
$$

as $t \rightarrow \infty$. So we can choose $t^{0}>0$ large enough and $e=t \varphi_{1}$ so that $I(e)<0$ and $\|e\|>\rho$.
Proof of Theorem 1.1. Set $\rho$ as in Lemma 2.1(a), define

$$
\bar{B}_{\rho}=\left\{u \in H_{0}^{1}(\Omega):\|u\| \leq \rho\right\}, \quad \partial B_{\rho}=\left\{u \in H_{0}^{1}(\Omega):\|u\|=\rho\right\}
$$

and $\bar{B}_{\rho}$ is a complete metric space with the distance

$$
\operatorname{dist}(u, v)=\|u-v\| \text { for } u, v \in \bar{B}_{\rho}
$$

By Lemma 2.1,

$$
\begin{equation*}
\left.I(u)\right|_{\partial B_{\rho}} \geq \gamma>0 \tag{2.4}
\end{equation*}
$$

Clearly, $I \in C^{1}\left(\bar{B}_{\rho}, \mathbb{R}\right)$, hence $I$ is lower semicontinuous and bounded from below on $\bar{B}_{\rho}$. Let

$$
\begin{equation*}
c_{1}=\inf \left\{I(u): u \in \bar{B}_{\rho}\right\} . \tag{2.5}
\end{equation*}
$$

We claim that

$$
\begin{equation*}
c_{1}<0 \tag{2.6}
\end{equation*}
$$

Indeed, let $v \in H_{0}^{1}(\Omega)$ be given by (H2), that is, $\int_{\Omega} h(x)\left(v^{+}\right)^{q+1} d x>0$, then for $t>0$ small enough such that for any $\varepsilon>0$, we have $|t v|<\varepsilon$. Therefore, together (F2) and $\alpha>1$ imply

$$
\begin{aligned}
I(t v)= & \frac{a t^{2}}{2} \int_{\Omega}|\nabla v|^{2} d x+\frac{b t^{4}}{4}\left(\int_{\Omega}|\nabla v|^{2} d x\right)^{2}-\frac{t^{q+1}}{q+1} \int_{\Omega} h(x)\left(v^{+}\right)^{q+1} d x \\
& -\int_{\Omega} F\left(x, t v^{+}\right) d x \\
\leq & \frac{a t^{2}}{2}\|v\|^{2}+\frac{b t^{4}}{4}\|v\|^{4}-\frac{t^{q+1}}{q+1} \int_{\Omega} h(x)\left(v^{+}\right)^{q+1} d x \\
& -\frac{a t^{2}}{2}(\alpha+\varepsilon) \lambda_{1}\|v\|_{2}^{2}-\frac{b t^{4}}{4}(\alpha+\varepsilon) \mu_{1}\|v\|_{4}^{4}<0
\end{aligned}
$$

if $t>0$ small enough, because $0<q<1$. So 2.6 is proved.
By the Ekeland's variational principle [6, Theorem 1.1] in $\bar{B}_{\rho}$ and (2.5), there is a minimizing sequence $\left\{u_{n}\right\} \subset \bar{B}_{\rho}$ such that
(i) $c_{1}<I\left(u_{n}\right)<c_{1}+\frac{1}{n}$,
(ii) $I(w) \geq I\left(u_{n}\right)-\frac{1}{n}\left\|w-u_{n}\right\|$ for all $w \in \bar{B}_{\rho}$.

So, $I^{\prime}\left(u_{n}\right) \rightarrow 0$ in $H_{0}^{-1}(\Omega)$ as $n \rightarrow \infty$. Moreover, by (i) and (ii), we obtain $I\left(u_{n}\right) \rightarrow c_{1}<0$ as $n \rightarrow \infty$.

From the above discussion, we know that $\left\{u_{n}\right\}$ is a bounded $(P S)$ sequence, there exist a subsequence (still denoted by $\left\{u_{n}\right\}$ ) and $u_{1} \in H_{0}^{1}(\Omega)$ such that

$$
\begin{gather*}
u_{n} \rightharpoonup u_{1} \quad \text { weakly in } H_{0}^{1}(\Omega), \\
u_{n} \rightarrow u_{1} \quad \text { a.e. in } \Omega  \tag{2.7}\\
u_{n} \rightarrow u_{1} \quad \text { strongly in } L^{r}(\Omega)
\end{gather*}
$$

as $n \rightarrow \infty$, where $r \in[1,6]$ if $N=3$ and $r \in(1,+\infty)$ if $N=1,2$. Thus, we have $\lim _{n \rightarrow \infty}\left\langle I^{\prime}\left(u_{n}\right), v\right\rangle=\left\langle I^{\prime}\left(u_{1}\right), v\right\rangle=0$ for all $v \in H_{0}^{1}(\Omega)$ and $\lim _{n \rightarrow \infty} I\left(u_{n}\right)=c_{1}<0$. Moreover, it follows from $\left\langle I^{\prime}\left(u_{1}\right), u_{1}^{-}\right\rangle=\left(a+b\left\|u_{1}\right\|^{2}\right)\left\|u_{1}^{-}\right\|^{2}=0$ that $u_{1}=u_{1}^{+} \geq$

0 , where $u_{1}^{-}=\max \left\{-u_{1}, 0\right\}$. Therefore, $u_{1}$ is a nonnegative critical point of $I$. Furthermore, if $h(x) \geq 0$, the strong maximum principle [8] implies that $u_{1}$ is a positive solution of problem (1.1).

## 3. Existence of a mountain pass solution

In this section, we use a variant version of mountain pass theorem to obtain a nonzero critical point of functional I; this theorem is used also in [10] and its proof can be found in [7], let us recall first this theorem.

Lemma 3.1 (Mountain Pass Theorem). Let E be a real Banach space with its dual space $E^{*}$ and suppose that $I \in C^{1}(E, R)$ satisfy the condition

$$
\max \{I(0), I(e)\} \leq \kappa<\gamma \leq \inf _{\|u\|=\rho}\{I(u)\}
$$

for some $\kappa<\gamma, \rho>0$ and $e \in E$ with $\|e\|>\rho$. Let $c \geq \gamma$ be characterized by

$$
c=\inf _{h \in \Gamma} \max _{t \in[0,1]} I(h(t))
$$

where $\Gamma=\{h \in([0,1], E) \mid h(0)=0, h(1)=e\}$ is the set of continuous paths joining 0 and $e$. Then, there exists a sequence $\left\{u_{n}\right\} \subset E$ such that

$$
I\left(u_{n}\right) \rightarrow c \geq \gamma \quad \text { and } \quad\left(1+\left\|u_{n}\right\|\right)\left\|I^{\prime}\left(u_{n}\right)\right\|_{E^{-1}} \rightarrow 0
$$

as $n \rightarrow \infty$.
Proof of Theorem 1.2. Let $\rho, \gamma$ and $e$ be given in Lemma 2.1, applying Lemma 3.1 with $\kappa=0, E=H_{0}^{1}(\Omega)$, and for $c$ defined as in Lemma 3.1, then there exists a sequence $\left\{u_{n}\right\} \subset H_{0}^{1}(\Omega)$ such that

$$
I\left(u_{n}\right) \rightarrow c \geq \gamma \quad \text { and } \quad\left(1+\left\|u_{n}\right\|\right)\left\|I^{\prime}\left(u_{n}\right)\right\|_{E^{-1}} \rightarrow 0
$$

as $n \rightarrow \infty$. This implies that

$$
\begin{align*}
& \quad \frac{a}{2} \int_{\Omega}\left|\nabla u_{n}\right|^{2} d x+\frac{b}{4}\left(\int_{\Omega}\left|\nabla u_{n}\right|^{2} d x\right)^{2}-\frac{1}{q+1} \int_{\Omega} h(x)\left(u_{n}^{+}\right)^{q+1} d x  \tag{3.1}\\
& \quad-\int_{\Omega} F\left(x, u_{n}^{+}\right) d x=c+o(1) \\
& a \int_{\Omega} \nabla u_{n} \cdot \nabla \varphi d x+b \int_{\Omega}\left|\nabla u_{n}\right|^{2} d x \int_{\Omega} \nabla u_{n} \cdot \nabla \varphi d x-\frac{1}{q+1} \int_{\Omega} h(x)\left(u_{n}^{+}\right)^{q} \varphi \\
& -\int_{\Omega} f\left(x, u_{n}^{+}\right) \varphi d x=o(1), \quad \text { for } \varphi \in H_{0}^{1}(\Omega),  \tag{3.2}\\
& a \int_{\Omega}\left|\nabla u_{n}\right|^{2} d x+b\left(\int_{\Omega}\left|\nabla u_{n}\right|^{2} d x\right)^{2}-\int_{\Omega} h(x)\left(u_{n}^{+}\right)^{q+1} d x-\int_{\Omega} f\left(x, u_{n}^{+}\right) u_{n}^{+} d x=o(1) . \tag{3.3}
\end{align*}
$$

By the compactness of Sobolev embedding and the standard procedures, we know that, if $\left\{u_{n}\right\}$ is bounded in $H_{0}^{1}(\Omega)$, there exists $u_{2} \in H_{0}^{1}(\Omega)$ such that $I^{\prime}\left(u_{2}\right)=0$ and $I\left(u_{2}\right)=c>0$ and $u_{2}$ is a nonnegative weak solution of problem 1.1, which is positive if $h(x) \geq 0$ by the strong maximum principle. Moreover, $u_{2}$ is different from the solution $u_{1}$ obtained in Theorem 1.1 since $I\left(u_{1}\right)=c_{1}<0$. So, to prove Theorem 1.2, we only need to prove that $\left\{u_{n}\right\}$ given by $(3.1)-(3.3)$ is bounded in $H_{0}^{1}(\Omega)$.

Next, we shall show that $\left\{u_{n}\right\}$ is bounded in $H_{0}^{1}(\Omega)$. By contradiction, we suppose that $\left\|u_{n}\right\| \rightarrow \infty$ as $n \rightarrow \infty$, and set $w_{n}=\frac{u_{n}}{\left\|u_{n}\right\|}$. Clearly, $\left\{w_{n}\right\}$ is bounded
in $H_{0}^{1}(\Omega)$. Thus, there exist a subsequence, still denoted by $\left\{w_{n}\right\}$, and $w \in H_{0}^{1}(\Omega)$, such that

$$
\begin{gathered}
w_{n} \rightarrow w \text { weakly in } H_{0}^{1}(\Omega), \\
w_{n} \rightarrow w \text { a.e. in } \Omega, \\
w_{n} \rightarrow w \text { strongly in } L^{r}(\Omega)
\end{gathered}
$$

as $n \rightarrow \infty$, where $r \in[1,6]$ if $N=3$ and $r \in(1,+\infty)$ if $N=1,2$.
Similarly, $w_{n}^{+}=\frac{u_{n}^{+}}{\left\|u_{n}\right\|}$ also satisfies

$$
\begin{gathered}
w_{n}^{+} \rightharpoonup w^{+} \quad \text { weakly in } H_{0}^{1}(\Omega), \\
w_{n}^{+} \rightarrow w^{+} \quad \text { a.e. in } \Omega, \\
w_{n}^{+} \rightarrow w^{+} \quad \text { strongly in } L^{r}(\Omega)
\end{gathered}
$$

as $n \rightarrow \infty$. We first claim that $w \not \equiv 0$. Indeed, if $w \equiv 0$, then by (H1), we have

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \int_{\Omega} h(x)\left(w_{n}^{+}\right)^{q+1} d x=0 \tag{3.4}
\end{equation*}
$$

Moreover, by (F1)-(F2), for any $\varepsilon>0$, if $s>0$ large enough, we obtain

$$
(\beta-\varepsilon) a \lambda_{1} s+(\beta-\varepsilon) b \mu_{1} s^{3}<f(x, s)<(\beta+\varepsilon) a \lambda_{1} s+(\beta+\varepsilon) b \mu_{1} s^{3} .
$$

Therefore, we deduce

$$
(\beta-\varepsilon) a \lambda_{1} s-\varepsilon b \mu_{1} s^{3}<f(x, s)-\beta b \mu_{1} s^{3}<(\beta+\varepsilon) a \lambda_{1} s+\varepsilon b \mu_{1} s^{3} .
$$

It implies that

$$
\begin{aligned}
& \frac{(\beta-\varepsilon) \lambda_{1}}{\left\|u_{n}\right\|^{2}} \int_{\Omega} w_{n}^{+} \varphi d x-\varepsilon b \mu_{1} \int_{\Omega}\left(w_{n}^{+}\right)^{3} \varphi d x \\
& <\int_{\Omega} \frac{f\left(x, u_{n}^{+}\right)-b \beta \mu_{1}\left(u_{n}^{+}\right)^{3}}{\left\|u_{n}\right\|^{3}} \varphi d x \\
& <\frac{(\beta+\varepsilon) \lambda_{1}}{\left\|u_{n}\right\|^{2}} \int_{\Omega} w_{n}^{+} \varphi d x-\varepsilon b \mu_{1} \int_{\Omega}\left(w_{n}^{+}\right)^{3} \varphi d x
\end{aligned}
$$

for any $\varphi \in H_{0}^{1}(\Omega)$. By the arbitrariness of $\varepsilon$, we obtain

$$
\begin{equation*}
\lim _{n \rightarrow+\infty} \int_{\Omega} \frac{f\left(x, u_{n}^{+}\right)-b \beta \mu_{1}\left(u_{n}^{+}\right)^{3}}{\left\|u_{n}\right\|^{3}} \varphi d x=0 . \tag{3.5}
\end{equation*}
$$

Multiplying (3.2) by $\frac{1}{\left\|u_{n}\right\|^{3}}$, we have

$$
\begin{align*}
& \frac{a}{\left\|u_{n}\right\|^{2}} \int_{\Omega} \nabla w_{n} \cdot \nabla \varphi d x+b \int_{\Omega} \nabla w_{n} \cdot \nabla \varphi d x-\frac{1}{\left\|u_{n}\right\|^{3-q}} \int_{\Omega} h(x)\left(w_{n}^{+}\right)^{q} \varphi d x \\
& -b \beta \mu_{1} \int_{\Omega}\left(w_{n}^{+}\right)^{3} \varphi d x-\int_{\Omega} \frac{f\left(x, u_{n}^{+}\right)-b \beta \mu_{1}\left(u_{n}^{+}\right)^{3}}{\left\|u_{n}\right\|^{3}} \varphi d x=o(1) . \tag{3.6}
\end{align*}
$$

Letting $n \rightarrow \infty$ in (3.6), according to $\left\|u_{n}\right\| \rightarrow \infty$ as $n \rightarrow \infty$, (3.4, (3.5) and $b \neq 0$, we have

$$
\int_{\Omega} \nabla w \cdot \nabla \varphi d x=\beta \mu_{1} \int_{\Omega}\left(w^{+}\right)^{3} \varphi d x
$$

and $w \neq 0$. Hence, $\beta \mu_{1}$ is an eigenvalue of 1.3), which contradicts with the assumption. The proof is complete.

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