Electronic Journal of Differential Equations, Vol. 2008(2008), No. 06, pp. 1-14. ISSN: 1072-6691. URL: http://ejde.math.txstate.edu or http://ejde.math.unt.edu ftp ejde.math.txstate.edu (login: ftp)

# EXISTENCE OF GLOBAL SOLUTIONS FOR A PREDATOR-PREY MODEL WITH CROSS-DIFFUSION 

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#### Abstract

In this article, we prove the existence of global classical solutions for a prey-predator model when the space dimension $n<10$. Under certain conditions on the coefficients of the reaction functions, the convergence of solutions is established for the system with large diffusion by constructing a Lyapunov function.


## 1. Introduction

To investigate the spatial segregation under the self and cross population pressure, Shigesada, Kawasaki and Teramoto [1] proposed a competition model in 1979. Then there have been established many results in the literatures; see for example [2, 3, 4, 5, 6, 7, 8, 2]. For the cross-diffusion systems with prey-predator type reaction functions, there are a few results mainly on the steady-state problems with the elliptic systems, see [10, 11, 12, 13, 14.

In this paper, we study the following cross-diffusion system, with prey-predator type reactions,

$$
\begin{array}{cl}
u_{t}-\Delta\left[\left(d_{1}+\alpha_{11} u+\alpha_{12} v\right) u\right]=u\left(a_{1}-b_{1} u-c_{1} v\right) & \text { in } \Omega \times[0, \infty), \\
v_{t}-\Delta\left[\left(d_{2}+\alpha_{21} u+\alpha_{22} v\right) v\right]=v\left(a_{2}+b_{2} u-c_{2} v\right) & \text { in } \Omega \times[0, \infty),  \tag{1.1}\\
\partial_{\eta} u=\partial_{\eta} v=0 \quad \text { on } \partial \Omega \times[0, \infty), & \\
u(x, 0)=u_{0}(x) \geq 0, \quad v(x, 0)=v_{0}(x) \geq 0 & \text { in } \Omega,
\end{array}
$$

where $\Omega \subset \mathbb{R}^{n}(n \geq 1)$ is a bounded domain with smooth boundary $\partial \Omega, \eta$ is the outward unit normal vector of the boundary $\partial \Omega$, and $\partial_{\eta}=\partial / \partial_{\eta} . \alpha_{i j}$ are given nonnegative constants for $i, j=1,2$. And $d_{i}, b_{i}, c_{i}(i=1,2)$ and $a_{1}$ are positive constants only $a_{2}$ may be non-positive.

In system 1.1 , $u$ and $v$ are nonnegative functions which represent the population densities of the prey and predator species, respectively, $d_{1}$ and $d_{2}$ are the random diffusion rates of the two species, $\alpha_{11}$ and $\alpha_{22}$ are self-diffusion rates, and $\alpha_{12}$ and $\alpha_{21}$ are the so-called cross-diffusion rates. When $\alpha_{i j}=0(i, j=1,2)$, the system is the well-known Lotka-Volterra prey-predator model. For more details on the biological background, see references [1, 18].

[^0]Local existence (in time) of solutions to (1.1) was established by Amann in a series of important papers [15, 16, 17]. His result can be summarized as follows:

Theorem 1.1. Suppose that $u_{0}, v_{0}$ are in $W_{p}^{1}(\Omega)$ for some $p>n$. Then 1.1) has a unique non-negative smooth solution $u(x, t), v(x, t)$ in

$$
C\left([0, T), W_{p}^{1}(\Omega)\right) \bigcap C^{\infty}\left((0, T), C^{\infty}(\Omega)\right)
$$

with maximal existence time T. Moreover, if the solution $(u, v)$ satisfies the estimate

$$
\sup _{0 \leq t \leq T}\|u(., t)\|_{W_{p}^{1}(\Omega)}<\infty \quad \text { and } \quad \sup _{0 \leq t \leq T}\|v(., t)\|_{W_{p}^{1}(\Omega)}<\infty
$$

then $T=\infty$.
However, little is known about global existence of solutions to 1.1). In 2006, Shim [18] proved the existence of global solutions to 1.1) in two cases: Case(A) $n=1, d_{1}=d_{2}$ and $\alpha_{11}=\alpha_{22}=0 ; \operatorname{Case}(\mathrm{B}) n=1,0<\alpha_{21}<8 \alpha_{11}$ and $0<\alpha_{12}<8 \alpha_{22}$.

In 19 the author considered the case when $\alpha_{11}, \alpha_{12}, \alpha_{22}>0$ and $\alpha_{21}=0$ for the system (1.1), and established the existence of global solutions with $n=1$.

We shall prove the existence of global solutions to the following system (namely, the system (1.1) for $\alpha_{12}=0$ )

$$
\begin{gather*}
u_{t}-\Delta\left[\left(d_{1}+\alpha_{11} u\right) u\right]=u\left(a_{1}-b_{1} u-c_{1} v\right) \quad \text { in } \Omega \times[0, \infty), \\
v_{t}-\Delta\left[\left(d_{2}+\alpha_{21} u+\alpha_{22} v\right) v\right]=v\left(a_{2}+b_{2} u-c_{2} v\right) \quad \text { in } \Omega \times[0, \infty), \\
\partial_{\eta} u=\partial_{\eta} v=0 \quad \text { on } \partial \Omega \times[0, \infty),  \tag{1.2}\\
u(x, 0)=u_{0}(x) \geq 0, \quad v(x, 0)=v_{0}(x) \geq 0 \quad \text { in } \Omega .
\end{gather*}
$$

This paper draws on ideas from two papers [6] and [9] which deal with crossdiffusion system with competition type reactions. Duo to the difference in the reaction functions. Therefore, in order to obtain the $L^{p}$-estimate of $v$, we have to estimate the term $u v^{p}$. We also obtain result on the asymptotic stability of the global solution to (1.2) if the diffusion coefficients are large enough by an important Lemma 5.1 from [21]. We summarize our results in the following theorems:

Theorem 1.2. Let $\alpha_{22}>0$ and assume that $u_{0} \geq 0, v_{0} \geq 0$ satisfy zero Neumann boundary condition and belong to $C^{2+\lambda}(\bar{\Omega})$ for some $\lambda>0$. Then (1.2) possesses a unique non-negative solution $u, v \in C^{2+\lambda, \frac{2+\lambda}{2}}\left(\bar{\Omega} \times[0, \infty)\right.$ ) if either (i) $\alpha_{11}=0$ or (ii) $\alpha_{11}>0$ and $n<10$.

Theorem 1.3. Assume that all conditions in Theorem 1.2 are satisfied. Assume further that

$$
\begin{align*}
-\frac{a_{1} b_{2}}{b_{1} c_{2}} & <\frac{a_{2}}{c_{2}}<\frac{a_{1}}{c_{1}},  \tag{1.3}\\
4 \rho \overline{u v} d_{1} d_{2} & >m^{2}\left(\bar{v} \alpha_{21}\right)^{2} . \tag{1.4}
\end{align*}
$$

Then (1.2) has the unique positive equilibrium point $(\bar{u}, \bar{v})$ which is global asymptotic stable, where $m$ is the positive constant in Lemma 2.1 (independent of $d_{1}, d_{2}$ ), $\rho=\left(b_{2} c_{1}+2 b_{1} c_{2}\right) b_{2}^{-2}$ and

$$
(\bar{u}, \bar{v})=\left(\frac{a_{1} c_{2}-a_{2} c_{1}}{b_{1} c_{2}+b_{2} c_{1}}, \frac{a_{2} b_{1}+a_{1} b_{2}}{b_{1} c_{2}+b_{2} c_{1}}\right) .
$$

The paper is organized as follows. In section 2, we collect some well known results and prove three new lemmas that are needed in section 3 and section 4. In section 3, we establish $L^{r}$-estimates of the solution $v$ of 1.2 and in section 4 we give a proof of Theorem 1.2. In section 5, we give a proof of Theorem 1.3.

## 2. Preliminaries

We list here some notation.

$$
\begin{aligned}
& Q_{T}=\Omega \times[0, T) \\
& \|u\|_{L^{p, q}\left(Q_{T}\right)}=\left(\int_{0}^{T}\left(\int_{\Omega}|u(x, t)|^{p} d x\right)^{\frac{q}{p}} d t\right)^{1 / q}, L^{p}\left(Q_{T}\right):=L^{p, p}\left(Q_{T}\right) \\
& \|u\|_{W_{P}^{2,1}\left(Q_{T}\right)}:=\|u\|_{L^{p}\left(Q_{T}\right)}+\left\|u_{t}\right\|_{L^{p}\left(Q_{T}\right)}+\|\nabla u\|_{L^{p}\left(Q_{T}\right)}+\left\|\nabla^{2} u\right\|_{L^{p}\left(Q_{T}\right)} \\
& \|u\|_{V_{2}\left(Q_{T}\right)}:=\sup _{0 \leq t \leq T}\|u(., t)\|_{L^{2}(\Omega)}+\|\nabla u(x, t)\|_{L^{2}\left(Q_{T}\right)} .
\end{aligned}
$$

Firstly, we present some useful lemmas.
Lemma 2.1. Let $u$, $v$ be a solution of 1.2 in $[0, T)$. Then $0 \leq u \leq m$ and $v \geq 0$ in $Q_{T}$, where $m=\max \left\{\frac{a_{1}}{b_{1}},\left\|u_{0}\right\|_{L^{\infty}(\Omega)}\right\}$.
Proof. The first equation in 1.2 is expressed as

$$
\begin{equation*}
u_{t}=\left(d_{1}+2 \alpha_{11} u\right) \Delta u+2 \alpha_{11} \nabla u \cdot \nabla u+u\left(a_{1}-b_{1} u-c_{1} v\right) \tag{2.1}
\end{equation*}
$$

and the second equation is written as

$$
\begin{equation*}
v_{t}=\left(d_{2}+\alpha_{21} u+2 \alpha_{22} v\right) \Delta v+2\left(\alpha_{21} \nabla u+\alpha_{22} \nabla v\right) \nabla v+v\left(\alpha_{21} \Delta u+a_{2}+b_{2} u-c_{2} v\right) \tag{2.2}
\end{equation*}
$$

Then application of the maximum principle for 2.1 and 2.2 yields the nonnegative of $u$ and $v$. Applying the maximum principle to 2.1) again one can also show the boundedness of $u$.

Lemma 2.2. There exists a positive $C_{1}(T)$ such that

$$
\begin{gathered}
\sup _{0 \leq t \leq T}\|u(., t)\|_{L^{1}(\Omega)}<C_{1}(T), \quad \sup _{0 \leq t \leq T}\|v(., t)\|_{L^{1}(\Omega)}<C_{1}(T), \\
\|u\|_{L^{2}\left(Q_{T}\right)}<C_{1}(T), \quad\|v\|_{L^{2}\left(Q_{T}\right)}<C_{1}(T)
\end{gathered}
$$

Proof. Integrating the first equation in (1.2) over the domain $\Omega$, we have

$$
\begin{align*}
\frac{d}{d t} \int_{\Omega} u d x & =a_{1} \int_{\Omega} u d x-b_{1} \int_{\Omega} u^{2} d x-c_{1} \int_{\Omega} u v d x \\
& \leq a_{1} \int_{\Omega} u d x-b_{1} \int_{\Omega} u^{2} d x  \tag{2.3}\\
& \leq a_{1} \int_{\Omega} u d x-\frac{b_{1}}{|\Omega|}\left(\int_{\Omega} u d x\right)^{2}
\end{align*}
$$

where we used Hölder's inequality. Then we have $\|u(., t)\|_{L^{1}(\Omega)} \leq M_{1}^{\prime}$, where $M_{1}^{\prime}=$ $\max \left\{\left\|u_{0}\right\|_{L^{1}(\Omega)}, \frac{a_{1}}{b_{1}}|\Omega|\right\}$. Furthermore,

$$
\begin{equation*}
\sup _{0 \leq t \leq T}\|u(., t)\|_{L^{1}(\Omega)}<C_{1}(T) \tag{2.4}
\end{equation*}
$$

Since

$$
\begin{equation*}
\frac{d}{d t} \int_{\Omega} u d x \leq a_{1} \int_{\Omega} u d x-b_{1} \int_{\Omega} u^{2} d x \tag{2.5}
\end{equation*}
$$

Integrating 2.5 from 0 to $T$, we have

$$
\|u\|_{L^{2}\left(Q_{T}\right)}^{2} \leq \frac{a_{1}}{b_{1}} M_{1}^{\prime}\left|Q_{T}\right|+\left\|u_{0}\right\|_{L^{1}(\Omega)}
$$

Therefore,

$$
\begin{equation*}
\|u\|_{L^{2}\left(Q_{T}\right)} \leq C_{1}(T) \tag{2.6}
\end{equation*}
$$

Now, integrating the second equation in the system 1.2 over the domain $\Omega$ we have

$$
\begin{equation*}
\frac{d}{d t} \int_{\Omega} v d x=a_{2} \int_{\Omega} v d x+b_{2} \int_{\Omega} u v d x-c_{2} \int_{\Omega} v^{2} d x . \tag{2.7}
\end{equation*}
$$

Multiplying 2.3) by $\frac{b_{2}}{c_{1}}$ and adding it to (2.7), we have

$$
\begin{equation*}
\frac{d}{d t} \int_{\Omega}\left(\frac{b_{2}}{c_{1}} u+v\right) d x \leq \frac{a_{1} b_{2}}{c_{1}} \int_{\Omega} u d x+\left|a_{2}\right| \int_{\Omega} v d x-\frac{b_{1} b_{2}}{c_{1}} \int_{\Omega} u^{2} d x-c_{2} \int_{\Omega} v^{2} d x \tag{2.8}
\end{equation*}
$$

Then

$$
\begin{aligned}
& \min \left\{1, \frac{b_{2}}{c_{1}}\right\} \frac{d}{d t} \int_{\Omega}(u+v) d x \\
& \leq \max \left\{\frac{a_{1} b_{2}}{c_{1}},\left|a_{2}\right|\right\} \int_{\Omega}(u+v) d x-\min \left\{\frac{b_{1} b_{2}}{c_{1}}, c_{2}\right\} \int_{\Omega}\left(u^{2}+v^{2}\right) d x \\
& \leq \max \left\{\frac{a_{1} b_{2}}{c_{1}},\left|a_{2}\right|\right\} \int_{\Omega}(u+v) d x-\frac{1}{2} \min \left\{\frac{b_{1} b_{2}}{c_{1}}, c_{2}\right\}\left[\int_{\Omega}(u+v) d x\right]^{2} .
\end{aligned}
$$

Therefore, $\|v(., t)\|_{L^{1}(\Omega)} \leq M_{2}^{\prime}$, where $M_{2}^{\prime}=\max \left\{\frac{A_{1}}{A_{2}},\left\|u_{0}+v_{0}\right\|_{L^{1}(\Omega)}\right\}$,

$$
A_{1}=\frac{\max \left\{\frac{a_{1} b_{2}}{c_{1}},\left|a_{2}\right|\right\}}{\min \left\{1, \frac{b_{2}}{c_{1}}\right\}}, \quad A_{2}=\frac{\min \left\{\frac{b_{1} b_{2}}{c_{1}}, c_{2}\right\}}{2 \min \left\{1, \frac{b_{2}}{c_{1}}\right\}}
$$

Then

$$
\begin{equation*}
\sup _{0 \leq t \leq T}\|v(., t)\|_{L^{1}(\Omega)}<C_{1}(T) \tag{2.9}
\end{equation*}
$$

Integrating (2.8) from 0 to $T$, we have

$$
c_{2} \int_{Q_{T}} v^{2} d x d t \leq \frac{a_{1} b_{2}}{c_{1}} \int_{0}^{T} M_{1}^{\prime} d t+\left|a_{2}\right| \int_{0}^{T} M_{2}^{\prime} d t+\frac{b_{2}}{c_{1}}\left\|u_{0}\right\|_{L^{1}(\Omega)}+\left\|v_{0}\right\|_{L^{1}(\Omega)}
$$

which implies $\|v\|_{L^{2}\left(Q_{T}\right)} \leq C_{1}(T)$.
Lemma 2.3. Let $w_{1}=\left(d_{1}+\alpha_{11} u\right) u$. Then there exists a constant $C_{2}(T)$, depending on $\left\|u_{0}\right\|_{W_{2}^{1}(\Omega)}$ and $\left\|u_{0}\right\|_{L^{\infty}(\Omega)}$ such that

$$
\begin{equation*}
\left\|w_{1}\right\|_{W_{2}^{2,1}\left(Q_{T}\right)} \leq C_{2}(T) \tag{2.10}
\end{equation*}
$$

Furthermore, $\nabla w_{1} \in V_{2}\left(Q_{T}\right)$.
Proof. Note that $w_{1}$ satisfies the equation

$$
\begin{equation*}
w_{1 t}=\left(d_{1}+2 \alpha_{11} u\right) \Delta w_{1}+n_{1}+n_{2} v \tag{2.11}
\end{equation*}
$$

where $n_{1}=u\left(d_{1}+2 \alpha_{11} u\right)\left(a_{1}-b_{1} u\right), n_{2}=-c_{1}\left(d_{1}+2 \alpha_{11} u\right) u$ depend on $u$ and are bounded functions because of Lemma 2.1. Multiplying the above equation by $-\Delta w_{1}$ and integration by parts over $\Omega$, we have

$$
\begin{equation*}
\frac{1}{2} \frac{d}{d t} \int_{\Omega}\left|\nabla w_{1}\right|^{2} d x=-\int_{\Omega}\left(d_{1}+2 \alpha_{11} u\right)\left(\Delta w_{1}\right)^{2} d x-\int_{\Omega}\left(n_{1}+n_{2} v\right) \Delta w_{1} d x \tag{2.12}
\end{equation*}
$$

Integrating 2.12 from 0 to $t$, we obtain

$$
\begin{aligned}
& \frac{1}{2} \int_{\Omega}\left|\nabla w_{1}(x, t)\right|^{2} d x-\frac{1}{2} \int_{\Omega}\left|\nabla w_{1}(x, 0)\right|^{2} d x \\
& =-\int_{Q_{t}}\left(d_{1}+2 \alpha_{11} u\right)\left(\Delta w_{1}\right)^{2} d x d t-\int_{Q_{t}}\left(n_{1}+n_{2} v\right) \Delta w_{1} d x d t \\
& \leq-d_{1} \int_{Q_{t}}\left|\Delta w_{1}\right|^{2} d x d t+\int_{Q_{t}}\left(n_{1}+n_{2} v\right) \cdot\left|\Delta w_{1}\right| d x d t
\end{aligned}
$$

By Young's inequality and Hölder's inequality, we have

$$
\begin{aligned}
& \frac{1}{2} \int_{\Omega}\left|\nabla w_{1}(x, t)\right|^{2} d x+d_{1} \int_{Q_{T}}\left|\Delta w_{1}\right|^{2} d x d t \\
& \leq\left(\left\|n_{1}\right\|_{L^{2}\left(Q_{T}\right)}+\left\|n_{2} v\right\|_{L^{2}\left(Q_{T}\right)}\right) \cdot\left\|\Delta w_{1}\right\|_{L^{2}\left(Q_{T}\right)}+\frac{1}{2} \int_{\Omega}\left|\nabla w_{1}(x, 0)\right|^{2} d x \\
& \leq m_{1}\left(1+\|v\|_{L^{2}\left(Q_{T}\right)}\right) \cdot\left\|\Delta w_{1}\right\|_{L^{2}\left(Q_{T}\right)}+\frac{1}{2} \int_{\Omega}\left|\nabla w_{1}(x, 0)\right|^{2} d x \\
& \leq m_{1}\left(1+C_{1}(T)\right) \cdot\left\|\Delta w_{1}\right\|_{L^{2}\left(Q_{T}\right)}+\frac{1}{2} \int_{\Omega}\left|\nabla w_{1}(x, 0)\right|^{2} d x \\
& \leq \frac{d_{1}}{2}\left\|\Delta w_{1}\right\|_{L^{2}\left(Q_{T}\right)}^{2}+\frac{m_{1}^{2}\left(1+C_{1}(T)\right)^{2}}{2 d_{1}}+\frac{1}{2} \int_{\Omega}\left|\nabla w_{1}(x, 0)\right|^{2} d x .
\end{aligned}
$$

Therefore,

$$
\sup _{0 \leq t \leq T} \int_{\Omega}\left|\nabla w_{1}\right|^{2}(x, t) d x+d_{1} \int_{Q_{T}}\left|\Delta w_{1}\right|^{2} d x d t \leq m_{2}
$$

where $m_{2}$ depends on $\left\|u_{0}\right\|_{W_{2}^{1}(\Omega)}$ and $\left\|u_{0}\right\|_{L^{\infty}(\Omega)}$. This implies $\nabla w_{1} \in V_{2}\left(Q_{T}\right)$. Since $w_{1} \in L^{2}\left(Q_{T}\right)$ we have from the elliptic regularity estimate [2, Lemma 2.3]

$$
\int_{Q_{T}}\left\|\left(w_{1}\right)_{x_{i} x_{j}}\right\|^{2} d x d t \leq m_{3} \quad \text { for } i, j=1, \ldots, n
$$

From 2.11, since $n_{1}, n_{2}$ and $u$ are bounded and $v \in L^{2}\left(Q_{T}\right)$, we have $w_{1 t} \in$ $L^{2}\left(Q_{T}\right)$. Hence, $w_{1} \in W_{2}^{2,1}\left(Q_{T}\right)$.

Let $a(x, t, \xi)$ be continuous and $(x, \xi)$-differentiable for $(x, t, \xi) \in \mathrm{Q}_{\mathrm{T}} \times \mathrm{R}$. Assume also that $a(x, t, \xi)$ satisfies the following conditions
(i) There is $d>0$ such that $a(x, t, \xi) \geq d$ and $a_{\xi}(x, t, \xi) \geq 0$ for all $(x, t) \in Q_{T}$ and $\xi$ in R .
(ii) There is a continuous function $M$ on R such that $a(x, t, \xi) \leq M(\xi)$ for all $(x, t) \in Q_{T}$.
(iii) For any bounded measurable function $g$ on $Q_{T},\left|\nabla_{x} a(., ., g(.,)).\right|$ is in the space $L^{2 p}\left(Q_{T}\right)$.

Lemma 2.4. Assume that $w \in W_{p}^{2,1}\left(Q_{T}\right) \bigcap C^{2,1}(\bar{\Omega} \times[0, T))$ is a bounded function satisfying

$$
w_{t} \leq a(x, t, w) \Delta w+f(x, t) \quad \text { in } \quad Q_{T}
$$

with boundary condition $\frac{\partial w}{\partial \nu} \leq 0$ on $\partial_{Q_{T}}$, where $f \in L^{p}\left(Q_{T}\right)$. Then, $\nabla w$ is in $L^{2 p}\left(Q_{T}\right)$.

The proof of the above lemma can be found in [9, Proposition 2.1].

Lemma 2.5. Let $q>1, \widetilde{q}=2+\frac{4 q}{n(q+1)}, \widetilde{\beta}$ in $(0,1)$ and let $C_{T}>0$ be any number which may depend on $T$. Then there is a constant $M_{1}$ depending on $q, n, \Omega, \widetilde{\beta}$ and $C_{T}$ such that for any $g$ in $C\left([0, T), W_{2}^{1}(\Omega)\right)$ with $\left(\int_{\Omega}|g(., t)|^{\widetilde{\beta}} d x\right)^{1 / \widetilde{\beta}} \leq C_{T}$ for all $t \in[0, T]$, we have the inequality

$$
\|g\|_{L^{\widetilde{q}}\left(Q_{T}\right)} \leq M_{1}\left\{1+\left(\sup _{0 \leq t \leq T}\|g(., t)\|_{L^{2 q / q+1}(\Omega)}\right)^{4 q / n(q+1) \widetilde{q}}\|\nabla g\|_{L^{2}\left(Q_{T}\right)}^{2 / \widetilde{q}}\right\}
$$

The proof of the above lemma can be found in [6, Lemmas 2.3, 2.4].

## 3. $L^{r}$-ESTIMATES FOR $v$

Lemma 3.1. There exists a constant $C_{3}(T)$ such that $\|\nabla u\|_{L^{4}\left(Q_{T}\right)} \leq C_{3}(T)$.
Proof. Let $\delta=\alpha_{11} / d_{1}, w_{1}=(1+\delta u) u$. By Lemma 2.1. $u$ is bounded. Therefore, $w_{1}$ is also bounded. By Lemma 2.3 , we have $w_{1} \in W_{2}^{2,1}\left(Q_{T}\right)$. Moreover, $w_{1}$ satisfies

$$
\begin{aligned}
w_{1 t} & \leq d_{1}(1+2 \delta u) \Delta w_{1}+a_{1} u(1+2 \delta u) \\
& =\sqrt{d_{1}^{2}+4 \delta d_{1} w_{1}} \Delta w_{1}+a_{1} u(1+2 \delta u)
\end{aligned}
$$

By Lemma 2.4 with $p=2, a(x, t, \xi)=\sqrt{d_{1}^{2}+4 \delta d_{1} \xi}, f(x, t)=a_{1} u(x, t)(1+$ $2 \delta u(x, t)$ ), we obtain the desired result.

Lemma 3.2. Let $r>2$ and $p_{r}=\frac{2 r}{r-2}$ be two positive numbers. Assume that $\alpha_{22}>0$ and assume also that there is a constant $M_{r, T}>0$ depending only on $r, T, \Omega$ and the coefficients of 1.2 such that

$$
\|\nabla u\|_{L^{r}\left(Q_{T}\right)} \leq M_{r, T} .
$$

Then for any $q>1$, there exists a constant $C(r, q, T)>0$ such that

$$
\begin{align*}
& \|v(., t)\|_{L^{q}(\Omega)}^{q}+\left\|\nabla\left(v^{q / 2}\right)\right\|_{L^{2}\left(Q_{t}\right)}^{2}+\left\|\nabla\left(v^{(q+1) / 2}\right)\right\|_{L^{2}\left(Q_{t}\right)}^{2} \\
& \leq C(r, q, T)\left(1+\|v\|_{L^{\frac{p_{r(q-1)}^{2}}{2}}\left(Q_{t}\right)}^{q-1}\right) . \tag{3.1}
\end{align*}
$$

Proof. For any constant $q>1$, multiplying the second equation of 1.2 by $q v^{q-1}$ and using the integration by parts, we obtain

$$
\begin{aligned}
& \frac{\partial}{\partial t} \int_{\Omega} v^{q} d x \\
& =q \int_{\Omega} v^{q-1} \nabla \cdot\left[\left(d_{2}+\alpha_{21} u+2 \alpha_{22} v\right) \nabla v+\alpha_{21} v \nabla u\right] d x+q \int_{\Omega} v^{q}\left(a_{2}+b_{2} u-c_{2} v\right) d x \\
& =-q(q-1) \int_{\Omega} v^{q-2}\left(d_{2}+\alpha_{21} u+2 \alpha_{22} v\right)|\nabla v|^{2} d x-\alpha_{21}(q-1) \int_{\Omega} \nabla\left(v^{q}\right) \cdot \nabla u d x \\
& \quad+q \int_{\Omega} v^{q}\left(a_{2}+b_{2} u-c_{2} v\right) d x \\
& \leq-q(q-1) d_{2} \int_{\Omega} v^{q-2}|\nabla v|^{2} d x-2 \alpha_{22} q(q-1) \int_{\Omega} v^{q-1}|\nabla v|^{2} d x \\
& \quad-\alpha_{21}(q-1) \int_{\Omega} \nabla\left(v^{q}\right) \cdot \nabla u d x+q \int_{\Omega} v^{q}\left(a_{2}+b_{2} u-c_{2} v\right) d x \\
& =-\frac{4(q-1) d_{2}}{q} \int_{\Omega}\left|\nabla\left(v^{\frac{q}{2}}\right)\right|^{2} d x-\frac{8 \alpha_{22} q(q-1)}{(q+1)^{2}} \int_{\Omega}\left|\nabla\left(v^{\frac{q+1}{2}}\right)\right|^{2} d x
\end{aligned}
$$

$$
-\alpha_{21}(q-1) \int_{\Omega} \nabla\left(v^{q}\right) \cdot \nabla u d x+q \int_{\Omega} v^{q}\left(a_{2}+b_{2} u-c_{2} v\right) d x
$$

Integrating the above inequality from 0 to $t$, we have

$$
\begin{align*}
& \int_{\Omega} v^{q}(x, t) d x+\frac{4(q-1) d_{2}}{q} \int_{Q_{t}}\left|\nabla\left(v^{\frac{q}{2}}\right)\right|^{2} d x d t+\frac{8 \alpha_{22} q(q-1)}{(q+1)^{2}} \int_{Q_{t}}\left|\nabla\left(v^{\frac{q+1}{2}}\right)\right|^{2} d x d t \\
& \leq \int_{\Omega} v^{q}(x, 0) d x-\alpha_{21}(q-1) \int_{Q_{t}} \nabla\left(v^{q}\right) \cdot \nabla u d x d t+q \int_{Q_{t}} v^{q}\left(a_{2}+b_{2} u-c_{2} v\right) d x d t \tag{3.2}
\end{align*}
$$

By Hölder's inequality, we have

$$
\begin{align*}
& q \int_{Q_{t}} v^{q}\left(a_{2}+b_{2} u-c_{2} v\right) d x d t \\
& =a_{2} q \int_{Q_{t}} v^{q} d x d t-c_{2} q \int_{Q_{t}} v^{q+1} d x d t+b_{2} q \int_{Q_{t}} u v^{q} d x d t \\
& \leq-c_{2} q\|v\|_{L^{q+1}\left(Q_{t}\right)}^{q+1}+\left|a_{2}\right| q\left|Q_{T}\right|^{\frac{1}{q+1}}\|v\|_{L^{q+1}\left(Q_{t}\right)}^{q}+b_{2} q \int_{Q_{t}} u v^{q} d x d t  \tag{3.3}\\
& \leq-c_{2} q\|v\|_{L^{q+1}\left(Q_{t}\right)}^{q+1}+\left|a_{2}\right| q\left[\varepsilon\|v\|_{L^{q+1}\left(Q_{t}\right)}^{q+1}+\varepsilon^{-q}\left|Q_{T}\right|^{\frac{q}{q+1}}\right]+b_{2} q \int_{Q_{t}} u v^{q} d x d t \\
& \leq B_{1}+b_{2} q \int_{Q_{t}} u v^{q} d x d t
\end{align*}
$$

where $\varepsilon=\frac{c_{2}}{\left|a_{2}\right|}, B_{1}$ depends on $T, q,|\Omega|$ and the coefficients of $1.2 \mid$.
On the other hand, since that $\frac{1}{r}+\frac{1}{2}+\frac{1}{p_{r}}=1$, using the Hölder's inequality and Poincaré inequality, we have

$$
\begin{align*}
\int_{Q_{t}} u v^{q} d x d t & =\int_{Q_{t}} u \cdot v^{\frac{q-1}{2}} \cdot v^{\frac{q+1}{2}} d x d t \\
& \leq\left\|v^{\frac{q-1}{2}}\right\|_{L^{p_{r}}\left(Q_{t}\right)} \cdot\left\|v^{\frac{q+1}{2}}\right\|_{L^{2}\left(Q_{t}\right)} \cdot\|u\|_{L^{r}\left(Q_{t}\right)}  \tag{3.4}\\
& \leq C_{4} m\|v\|_{L^{\frac{p_{r}(q-1)}{2}}\left(Q_{t}\right)}^{(q-1) / 2} \cdot\left\|\nabla\left(v^{\frac{q+1}{2}}\right)\right\|_{L^{2}\left(Q_{t}\right)} .
\end{align*}
$$

The substitution (3.4) into (3.3) leads to

$$
\begin{equation*}
q \int_{Q_{t}} v^{q}\left(a_{2}+b_{2} u-c_{2} v\right) d x d t \leq B_{1}+C_{5}\|v\|_{L^{\frac{p_{r}(q-1)}{2}}\left(Q_{t}\right)}^{(q-1) / 2} \cdot\left\|\nabla\left(v^{\frac{q+1}{2}}\right)\right\|_{L^{2}\left(Q_{t}\right)} . \tag{3.5}
\end{equation*}
$$

Since that $\frac{1}{r}+\frac{1}{2}+\frac{1}{p_{r}}=1$ and $\nabla u$ is in $L^{r}\left(Q_{T}\right)$, using the Hölder's inequality, we have

$$
\begin{aligned}
\left|-\int_{Q_{t}} \nabla\left(v^{q}\right) \cdot \nabla u d x d t\right| & =\frac{2 q}{q+1}\left|\int_{Q_{t}} v^{\frac{q-1}{2}} \cdot \nabla\left(v^{\frac{(q+1)}{2}}\right) \cdot \nabla u d x d t\right| \\
& \leq \frac{2 q}{q+1}\left\|v^{\frac{q-1}{2}}\right\|_{L^{p_{r}}\left(Q_{t}\right)} \cdot\left\|\nabla\left(v^{\frac{q+1}{2}}\right)\right\|_{L^{2}\left(Q_{t}\right)} \cdot\|\nabla u\|_{L^{r}\left(Q_{t}\right)} \\
& \leq \frac{2 q}{q+1}\|v\|_{L^{\frac{q-1}{2}}}^{\frac{p_{r}(q-1)}{2}}\left(Q_{t}\right) \\
& \left.\leq \frac{2 q}{q+1} M_{r, T}\|v\|^{\frac{q-1}{2}} v_{L^{\frac{p_{r}(q-1)}{2}}\left(Q_{t}\right)}^{\frac{q+1}{2}}\right)\left\|_{L^{2}\left(Q_{t}\right)} \cdot\right\| \nabla u\left\|_{L^{r}\left(Q_{t}\right)} \cdot\right\| \nabla\left(v^{\frac{q+1}{2}}\right) \|_{L^{2}\left(Q_{t}\right)} .
\end{aligned}
$$

The substitution (3.5) and the above inequality into (3.2) leads to

$$
\begin{align*}
& \int_{\Omega} v^{q}(x, t) d x+\frac{4(q-1) d_{2}}{q} \int_{Q_{t}}\left|\nabla\left(v^{\frac{q}{2}}\right)\right|^{2} d x d t+\frac{8 \alpha_{22} q(q-1)}{(q+1)^{2}} \int_{Q_{t}}\left|\nabla\left(v^{\frac{q+1}{2}}\right)\right|^{2} d x d t \\
& \leq B_{2}+C_{6}\|v\|_{L^{\frac{q-1}{2}}}^{\frac{p_{r}(q-1)}{2}}\left(Q_{t}\right) \\
& \left.\leq B_{2}+\frac{C_{6}}{4 \varepsilon}\|v\|_{L^{\frac{q+1}{2}}}\right) \|_{L^{2}\left(Q_{t}\right)}^{q-1}  \tag{3.6}\\
& { }_{L}^{\frac{p_{r}(q-1)}{2}}\left(Q_{t}\right)
\end{align*}+C_{6} \varepsilon\left\|\nabla\left(v^{\frac{q+1}{2}}\right)\right\|_{L^{2}\left(Q_{t}\right)}^{2}, ~ l
$$

where $B_{2}>0$ depending on $q, T, \Omega$ coefficients of 1.2 and initial datal $v_{0}$. For any $\varepsilon>0$, from (3.6) and by choosing a sufficiently small $\varepsilon$, such that $C_{6} \varepsilon<\frac{8 \alpha_{22} q(q-1)}{(q+1)^{2}}$, we get (3.1). This completes the proof of the lemma.

For any number $a$, we denote $a_{+}=\max \{a, 0\}$.
Proposition 3.3. Let $\alpha_{22}>0$.
(i) If $\alpha_{11}>0$, then there is a constant $C_{7}(T)>0$ such that

$$
\|v\|_{V_{2}\left(Q_{T}\right)} \leq C_{7}(T)
$$

Moreover, for any constant $r<\frac{4(n+1)}{(n-2)_{+}}$, there exists a positive constant $C_{T}$ such that

$$
\|v\|_{L^{r}\left(Q_{T}\right)} \leq C_{T}
$$

(ii) If $\alpha_{11}=0$, then

$$
\|v\|_{L^{r}\left(Q_{T}\right)} \leq C_{T} \quad \text { for any } \quad r>1
$$

Proof. (i) Set $w=v^{(q+1) / 2}$ so that $v^{q}=w^{2 q /(q+1)}$ and $v^{q+1}=w^{2}$. Then

$$
\begin{aligned}
E & \equiv \sup _{0 \leq t \leq T} \int_{\Omega} v^{q}(x, t) d x+\int_{Q_{T}}\left|\nabla\left(v^{(q+1) / 2}\right)\right|^{2} d x d t \\
& =\sup _{0 \leq t \leq T} \int_{\Omega} w^{2 q / q+1} d x+\int_{Q_{T}}|\nabla w|^{2} d x d t
\end{aligned}
$$

Let $r_{0}=4, p_{0}=\frac{2 r_{0}}{r_{0}-2}$. By Lemma 3.1. we see that $\nabla u$ is in $L^{r_{0}}\left(Q_{T}\right)$. So, from Lemma 3.2. we have

$$
\begin{equation*}
E+\left\|\nabla\left(v^{\frac{q}{2}}\right)\right\|_{L^{2}\left(Q_{T}\right)}^{2} \leq C\left(r_{0}, q, T\right)\left(1+\|w\|_{L^{\frac{2(q-1)}{q+1}}}^{\frac{p_{0}(q-1)}{q+1}}\left(Q_{T}\right),\right. \tag{3.7}
\end{equation*}
$$

where $C\left(r_{0}, q, T\right)>0$ depending only $T, \Omega$, initial data $u_{0}, v_{0}$ and the coefficients of 1.2 . Since $q>1$, if we restrict our $q$ so that

$$
\begin{equation*}
\left(n p_{0}-2 n-4\right) q \leq 2 n+n p_{0} \tag{3.8}
\end{equation*}
$$

Then, $\frac{p_{0}(q-1)}{q+1} \leq \widetilde{q}$, where $\widetilde{q}=2+\frac{4 q}{n(q+1)}$. Therefore, by Hölder's inequality

$$
\begin{equation*}
\|w\|_{L^{\frac{p_{0}(q-1)}{q+1}}\left(Q_{T}\right)} \leq C_{8}(q, T)\|w\|_{L^{\tilde{q}}\left(Q_{T}\right)} \tag{3.9}
\end{equation*}
$$

where $C_{8}(q, T)=\left|Q_{T}\right|^{\frac{q+1}{p_{0}(q-1)}-\frac{1}{q}}$. Setting $\widetilde{\beta}=2 /(q+1) \in(0,1)$, by Lemma 2.2 we have

$$
\begin{equation*}
\|w(., t)\|_{L^{\tilde{\beta}}(\Omega)}=\|v(., t)\|_{L^{1}(\Omega)}^{\frac{1}{\tilde{\beta}}} \leq\left(C_{1}(T)\right)^{\frac{1}{\tilde{\beta}}}, \quad \forall t \in[0, T) \tag{3.10}
\end{equation*}
$$

Hence, by Lemma 2.5 and the definition of $E$, 3.10 yields

$$
\begin{equation*}
\|w\|_{L^{p_{0}(q-1) / q+1}\left(Q_{T}\right)} \leq C_{8}(q, T)\|w\|_{L^{\widetilde{q}}\left(Q_{T}\right)} \leq C_{8}(q, T) M_{1}\left\{1+E^{2 / n \widetilde{q}} E^{\frac{1}{q}}\right\} \tag{3.11}
\end{equation*}
$$

Then (3.7) together with the above inequality, we can find a constant $C_{9}(q, T)>0$ such that

$$
\begin{equation*}
E \leq C_{9}(q, T)\left(1+E^{\mu} E^{\nu}\right) \tag{3.12}
\end{equation*}
$$

with

$$
\mu=\frac{4(q-1)}{n \widetilde{q}(q+1)}, \quad \nu=\frac{2(q-1)}{\widetilde{q}(q+1)} .
$$

Since

$$
\mu+\nu=\frac{2(q-1)}{\widetilde{q}(q+1)}\left[\frac{2}{n}+1\right]<\frac{1}{\widetilde{q}}\left[\frac{4 q}{n(q+2)}+2\right]=1
$$

it is easy to see from 3.12 that $E$ is bounded. Therefore, from 3.11) and 3.12 we get $w \in L^{\widetilde{q}}\left(Q_{T}\right)$ which in turn implies that $v \in L^{r}\left(Q_{T}\right)$ with $r=\frac{\tilde{q}(q+1)}{2}$ for any $q$ satisfying 3.8. Now, looking at (3.8), if $n \leq 2$, we have

$$
\begin{equation*}
n p_{0}-2 n-4=2(n-2) \leq 0, \tag{3.13}
\end{equation*}
$$

then (3.8) holds for all $q$. so for $n \leq 2, v \in L^{r}\left(Q_{T}\right)$ for all $r>1$. Now, suppose that $n>2$, we see 3.8 is equivalent to

$$
1<q<q_{0}:=\frac{2 n+n p_{0}}{\left(n p_{0}-2 n-4\right)}=\frac{3 n}{n-2}
$$

Then, we have

$$
\frac{\widetilde{q}(q+1)}{2}=q+1+\frac{2 q}{n} \leq \bar{r}_{1}:=q_{0}+1+\frac{2 q_{0}}{n}=\frac{4(n+1)}{n-2}
$$

So, we see that $v$ is in $L^{r}\left(Q_{T}\right)$ for all $1<r \leq \bar{r}_{1}$. Since 3.8 holds true for $q=2$. So when $q=2$, we have $E$ is finite. Therefore, from (3.7) and (3.11), we see that $\|v\|_{V_{2}\left(Q_{T}\right)}$ is bounded for any $n$, this completes the proof of Proposition 3.3 when $\alpha_{11}>0$ and $r<\frac{4(n+2)}{(n-2)_{+}}$.

Next, we consider the case $\alpha_{11}=0$. By Hölder's inequality, we have

$$
\begin{align*}
& q \int_{Q_{t}} v^{q}\left(a_{2}+b_{2} u-c_{2} v\right) d x d t \\
& =a_{2} q \int_{Q_{t}} v^{q} d x d t-c_{2} q \int_{Q_{t}} v^{q+1} d x d t+b_{2} q \int_{Q_{t}} u v^{q} d x d t \\
& \leq-c_{2} q\|v\|_{L^{q+1}\left(Q_{t}\right)}^{q+1}+\left.\left.\left|a_{2}\right| q\right|_{T}\right|^{\frac{1}{q+1}}\|v\|_{L^{q+1}\left(Q_{t}\right)}^{q}  \tag{3.14}\\
& \quad+b_{2} q\|v\|_{L^{q+1}\left(Q_{t}\right)}^{q} \cdot\|u\|_{L^{q+1}\left(Q_{t}\right)} \\
& \leq-c_{2} q\|v\|_{L^{q+1}\left(Q_{t}\right)}^{q+1}+\left.\left.\left|a_{2}\right| q\right|_{T}\right|^{\frac{1}{q+1}}\|v\|_{L^{q+1}\left(Q_{t}\right)}^{q}+b_{2} q m\|v\|_{L^{q+1}\left(Q_{t}\right)}^{q} \\
& \leq-c_{2} q\|v\|_{L^{q+1}\left(Q_{t}\right)}^{q+1}+q \varepsilon\|v\|_{L^{q+1}\left(Q_{t}\right)}^{q+1}+B_{3} \\
& \leq B_{3}
\end{align*}
$$

where $\varepsilon=c_{2}$ and $B_{3}>0$ which depends only on $T, q,|\Omega|,\left\|u_{0}\right\|_{L^{\infty}(\Omega)}$ and the coefficients of 1.2 .

We can integrate by parts once to obtain from Lemma 2.1 and analogue of 20 , Theorem 9.1, p. 341-342] for Neumann boundary condition [20, p.351]

$$
\begin{align*}
& \left|-\int_{Q_{t}} \nabla\left(v^{q}\right) \cdot \nabla u d x d t\right| \\
& =\left|-\int_{Q_{t}} v^{q} \Delta u d x d t\right| \\
& \leq\|v\|_{L^{q+1}\left(Q_{T}\right)}^{q} \cdot\|\Delta u\|_{L^{q+1}\left(Q_{T}\right)}  \tag{3.15}\\
& \leq C_{10}\|v\|_{L^{q+1}\left(Q_{T}\right)}^{q}\left(\left\|u\left(a_{1}-b_{1} u-c_{1} v\right)\right\|_{L^{q+1}\left(Q_{T}\right)}+\left\|u_{0}\right\|_{W_{q+1}^{2-\frac{2}{q+1}}(\Omega)}\right) \\
& \leq C_{11}\left(1+\|v\|_{L^{q+1}\left(Q_{T}\right)}^{q+1}\right)
\end{align*}
$$

The substitution of $(3.14)$ and $(3.15)$ into $(3.2)$ leads to

$$
\begin{equation*}
\sup _{0 \leq t \leq T}\left\|v^{q}(t)\right\|_{L^{q}(\Omega)}^{q}+\left\|\nabla\left(v^{(q+1) / 2}\right)\right\|_{L^{2}\left(Q_{T}\right)}^{2} \leq C_{12}\left(1+\|v\|_{L^{q+1}\left(Q_{T}\right)}^{q+1}\right) \tag{3.16}
\end{equation*}
$$

We introduce $w=v^{\frac{q+1}{2}}$, then (3.16 leads to

$$
\begin{equation*}
E \equiv \sup _{0 \leq t \leq T}\|w(t)\|_{L^{\frac{2 q}{q+1}}(\Omega)}^{\frac{2 q}{q+1}}+\|\nabla w\|_{L^{2}\left(Q_{T}\right)}^{2} \leq C_{12}\left(1+\|w\|_{L^{2}\left(Q_{T}\right)}^{2}\right) \tag{3.17}
\end{equation*}
$$

Recall that Lemma 2.2 implies $v \in L^{2}\left(Q_{T}\right)$, so $\|w\|_{L^{\frac{4}{q+1}}\left(Q_{T}\right)} \leq C_{13}$. Since $\frac{4}{q+1}<$ $2 \leq \widetilde{q}$. Then we see from Hölder's inequality

$$
\begin{equation*}
\|w\|_{L^{2}\left(Q_{T}\right)}^{2} \leq\|w\|_{L^{\tilde{q}}\left(Q_{T}\right)}^{2(1-\lambda)}\|w\|_{L^{\frac{4}{q+1}}\left(Q_{T}\right)}^{2 \lambda} \leq C_{13}^{2 \lambda}\|w\|_{L^{\tilde{q}}\left(Q_{T}\right)}^{2(1-\lambda)}, \tag{3.18}
\end{equation*}
$$

where $\lambda=\left(\frac{1}{2}-\frac{1}{\widetilde{q}}\right) /\left(\frac{q+1}{4}-\frac{1}{\widetilde{q}}\right)$. Setting $\widetilde{\beta}=2 /(q+1) \in(0,1)$, we have $\|w(., t)\|_{L^{\tilde{\beta}}(\Omega)}=$ $\|v(., t)\|_{L^{1}(\Omega)}^{\frac{1}{\beta}} \leq C_{1}(T)^{\frac{1}{\beta}}$ for all $t \in[0, T)$ by Lemma 2.2. Then it follow from (3.17), (3.18) and Lemma 2.5 that

$$
\begin{equation*}
E \leq C_{14}\left(1+E^{\alpha}\right) \tag{3.19}
\end{equation*}
$$

with

$$
\alpha=\frac{2(1-\lambda)}{\widetilde{q}}\left(\frac{2}{n}+1\right)<1 .
$$

Thus (3.19) implies

$$
\sup _{0 \leq t \leq T}\|w(t)\|_{L^{\frac{2 q}{q+1}}(\Omega)}^{\frac{2 q}{q+1}} \leq E \leq C_{15}
$$

with some $C_{15}>0$, let $r=q>1$, so that $\sup _{0 \leq t \leq T}\|v(t)\|_{L^{r}(\Omega)} \leq C_{T}$ and the proof is complete.

## 4. Proof of Theorem 1.2

The first step of the proof is to show $v$ is in $L^{r}\left(Q_{T}\right)$ for any $r>1$.
Lemma 4.1. Let $\alpha_{11}>0$ and suppose that there are $r_{1}>\max \left\{\frac{n+2}{2}, 3\right\}$ and a positive constant $C_{r_{1}, T}$ such that

$$
\|v\|_{L^{r_{1}}\left(Q_{T}\right)} \leq C_{r_{1}, T}
$$

Then, $v$ is in $L^{r}\left(Q_{T}\right)$ for any $r>1$.

Proof. The proof is almost identical to [9, Lemma 4.1], but for completeness we repeat it here. First, the equation for $u$ can be written in the divergence form as

$$
\begin{equation*}
u_{t}=\nabla \cdot\left[\left(d_{1}+2 \alpha_{11} u\right) \nabla u\right]+u\left(a_{1}-b_{1} u-c_{1} v\right) \tag{4.1}
\end{equation*}
$$

where $d_{1}+2 \alpha_{11} u$ is bounded in $\bar{Q}_{T}$ by Lemma 2.1 and $u\left(a_{1}-b_{1} u-c_{1} v\right)$ is in $L^{r_{1}}$ with $r_{1}>\frac{n+2}{2}$. Application of the Hölder continuity result in [20, Theorem 10.1, p. 204] to (4.1) yields

$$
\begin{equation*}
u \in C^{\beta, \frac{\beta}{2}}\left(\bar{Q}_{T}\right) \quad \text { with some } \beta>0 . \tag{4.2}
\end{equation*}
$$

Moreover, we have

$$
\begin{equation*}
w_{1 t}=\left(d_{1}+2 \alpha_{11} u\right) \Delta w_{1}+f_{1} \tag{4.3}
\end{equation*}
$$

where $w_{1}=\left(d_{1}+\alpha_{11} u\right) u$ is as in the proof of Lemma 2.3, $f_{1}=\left(d_{1}+2 \alpha_{11} u\right) u\left(a_{1}-\right.$ $b_{1} u-c_{1} v$ ). Since $u$ is bounded and by the assumption of this Lemma, we see that $f_{1}$ is in $L^{r_{1}}\left(Q_{T}\right)$. From 4.2, Lemma 2.1 and Proposition 3.3, applying [20, Theorem 9.1, pp. 341-342] and its remark [20, P. 351], we have

$$
\begin{equation*}
w_{1} \in W_{r_{1}}^{2,1}\left(Q_{T}\right) \tag{4.4}
\end{equation*}
$$

This implies $\nabla u=\frac{1}{d_{1}+2 \alpha_{11} u} \nabla w_{1}$ in $L^{r_{1}}\left(Q_{T}\right)$. Now, following the proof of Proposition 3.3 with $r_{1}$ instead of $r_{0}$ and $p_{1}=\frac{2 r_{1}}{r_{1}-2}$ instead of $p_{0}$, we see that either $v$ is in $L^{r}\left(Q_{T}\right)$ for any $r>1$ or else $v$ is in $L^{r_{2}}\left(Q_{T}\right)$ with

$$
r_{2}:=\frac{(n+1) r_{1}}{n+2-r_{1}} .
$$

The later case happens if and only if $n+2-r_{1}>0$.
If $v$ is in $L^{r_{2}}\left(Q_{T}\right)$, we see that $f_{1}$ is in $L^{r_{2}}\left(Q_{T}\right)$. Therefore, applying [20, Theorem 9.1, p. 341-342] and its remark [20, p. 351] again, we have $\nabla u$ in $L^{r_{2}}\left(Q_{T}\right)$. Then we go back and do the same argument again. Keep doing likes this we will get a sequence of numbers

$$
\begin{equation*}
r_{k+1}:=\frac{(n+1) r_{k}}{n+2-r_{k}} . \tag{4.5}
\end{equation*}
$$

We stop and get the conclusion that $v$ is in $L^{r}\left(Q_{T}\right)$ for any $r>1$ when

$$
\begin{equation*}
n+2-r_{k} \leq 0 \tag{4.6}
\end{equation*}
$$

Since $r_{1}>3$, from 4.5 we can prove by induction that $r_{k}>3, k=1,2, \ldots$ Then, we have

$$
\begin{equation*}
\frac{r_{k+1}}{r_{k}}=\frac{n+1}{n+2-r_{k}} \geq \frac{n+1}{n-1}>1 \tag{4.7}
\end{equation*}
$$

Thus, the sequence $r_{k}$ is strictly increasing. Therefore, there must be some $k$ such that (4.6) holds. we stop at this $k$ and conclude that $v$ is in $L^{r}\left(Q_{T}\right)$ for any $r>1$, namely, there is a positive constant $C_{16}$ such that $\|v\|_{L^{r}\left(Q_{T}\right)} \leq C_{16}$, where $C_{16}>0$ depending on $q, T, \Omega$ and the coefficients of the system (1.2) but not on $r$.

So, from Proposition 3.3 and Lemma 4.1, we have the following lemma.
Lemma 4.2. Let $\alpha_{22}>0$ and suppose (i) $\alpha_{11}=0$ or (ii) $\alpha_{11}>0$ and $n<10$. Then there exists $M_{2}$ such that

$$
\|v\|_{L^{r}\left(Q_{T}\right)} \leq M_{2} \quad \text { for any } r>1
$$

Moreover, for any $r>1, v$ is in $V_{2}\left(Q_{T}\right)$.

Proof of Theorem 1.2. We give the proof only in case $\alpha_{11}>0$ because the proof for $\alpha_{11}=0$ is essentially the same. By Lemma 4.2, $v$ is bounded in $\bar{Q}_{T}$. From 4.3), we have

$$
w_{1 t}=\left(d_{1}+2 \alpha_{11} u\right) \Delta w_{1}+f_{1}
$$

where $f_{1}=\left(d_{1}+2 \alpha_{11} u\right) u\left(a_{1}-b_{1} u-c_{1} v\right)$ is bounded in $\bar{Q}_{T}$ by Lemma 2.1 and Lemma 4.2. $\left(d_{1}+2 \alpha_{11} u\right) \in C^{\beta, \frac{\beta}{2}}\left(Q_{T}\right)$ by 4.2). By [20, Theorem 9.1, p.341-342], we have

$$
\left\|w_{1}\right\|_{W_{r}^{2,1}}\left(Q_{T}\right)<M_{3} \quad \text { for } \frac{n+2}{2}<r<\frac{4(n+1)}{(n-2)_{+}}
$$

Hence it follows from [20, Lemma 3.3, p.80] that

$$
\begin{equation*}
w_{1} \in C^{1+\beta^{*}}, \frac{\left(1+\beta^{*}\right)}{2}\left(\bar{Q}_{T}\right), \quad \forall 0<\beta^{*}<1 \tag{4.8}
\end{equation*}
$$

Since $u=\frac{-d_{1}+\sqrt{d_{1}^{2}+4 w_{1} \alpha_{11}}}{2 \alpha_{11}}$, it follow from 4.8 that

$$
\begin{equation*}
u \in C^{1+\beta^{*}}, \frac{\left(1+\beta^{*}\right)}{2}\left(\bar{Q}_{T}\right), \quad \forall 0<\beta^{*}<1 \tag{4.9}
\end{equation*}
$$

Next, we rewrite the equation for $v$ in divergence form as

$$
v_{t}=\nabla \cdot\left[\left(d_{2}+\alpha_{21} u+2 \alpha_{22} v\right) \nabla v+\alpha_{21} v \nabla u\right]+f_{2}(x, t),
$$

where $f_{2}(x, t)=v\left(a_{2}+b_{2} u-c_{2} v\right), u, v$ and $\nabla u$ are all bounded functions because of Lemma 2.1. Lemma 4.2 and 4.9. By [20, Theorem 10.1, p.204], we have

$$
\begin{equation*}
v \in C^{\sigma, \frac{\sigma}{2}}\left(\bar{Q}_{T}\right) \text { with some } 0<\sigma<1 \tag{4.10}
\end{equation*}
$$

Now, we then return to the equation for $u$ and write it as

$$
\begin{equation*}
u_{t}=\left(d_{1}+2 \alpha_{11} u\right) \Delta u+f_{3}(x, t) \tag{4.11}
\end{equation*}
$$

where $f_{3}(x, t)=2 \alpha_{11}|\nabla u|^{2}+u\left(a_{1}-b_{1} u-c_{1} v\right) \in C^{\sigma, \frac{\sigma}{2}}\left(\bar{Q}_{T}\right)$ by 4.9) and 4.10). Then the Schuader estimate in [20, Theorem 5.3, p.320-321] applied to 4.11) yields

$$
\begin{equation*}
u \in C^{2+\sigma^{*}, \frac{2+\sigma^{*}}{2}}\left(\bar{Q}_{T}\right) \quad \text { with } \sigma^{*}=\min \{\lambda, \sigma\} \tag{4.12}
\end{equation*}
$$

Let $w_{2}=\left(d_{2}+\alpha_{21} u+\alpha_{22} v\right) v$. Then $w_{2}$ satisfies

$$
\begin{equation*}
w_{2 t}=\left(d_{2}+\alpha_{21} u+2 \alpha_{22} v\right) \Delta w_{2}+f_{4}(x, t) \tag{4.13}
\end{equation*}
$$

where $f_{4}(x, t)=\left(d_{2}+\alpha_{21} u+2 \alpha_{22} v\right) v\left(a_{2}+b_{2} u-c_{2} v\right)+\alpha_{21} v u_{t} \in C^{\sigma^{*}, \frac{\sigma^{*}}{2}}\left(\bar{Q}_{T}\right)$ by (4.11) and 4.12), $d_{2}+\alpha_{21} u+2 \alpha_{22} v \in C^{\sigma, \frac{\sigma}{2}}\left(\bar{Q}_{T}\right)$ by 4.9) and 4.10, by applying the Schuader estimate to the equation 4.13), we have

$$
\begin{equation*}
w_{2} \in C^{2+\sigma^{*}, \frac{2+\sigma^{*}}{2}}\left(\bar{Q}_{T}\right) \tag{4.14}
\end{equation*}
$$

Then

$$
\begin{equation*}
v=\frac{-\left(d_{2}+\alpha_{21} u\right)+\sqrt{\left(d_{2}+\alpha_{21} u\right)^{2}+4 w_{2} \alpha_{22}}}{2 \alpha_{22}} \in C^{2+\sigma^{*}, \frac{2+\sigma^{*}}{2}}\left(\bar{Q}_{T}\right) \tag{4.15}
\end{equation*}
$$

Now repeat the procedure by making use of 4.12 and 4.15 in place of 4.9 and (4.10), we have

$$
\begin{equation*}
u, v \in C^{2+\lambda, \frac{2+\lambda}{2}}\left(\bar{Q}_{T}\right) \tag{4.16}
\end{equation*}
$$

Finally, the estimates (4.12) and 4.15 imply that the hypotheses of Theorem 1.1 are satisfied. So that $(u, v)$ exists globally in time. The proof of Theorem 1.2 is now complete.

## 5. Stability

In this section, we discuss global asymptotic stability of positive equilibrium point $(\bar{u}, \bar{v})$ for $\sqrt{1.2}$, namely to prove Theorem 1.3 .

Proof of Theorem 1.3. Define the Lyapunov function:

$$
H(u, v)=\int_{\Omega}\left[\left(u-\bar{u}-\bar{u} \ln \frac{u}{\bar{u}}\right)+\rho\left(v-\bar{v}-\bar{v} \ln \frac{v}{\bar{v}}\right)\right] d x
$$

where $\rho=\left(b_{2} c_{1}+2 b_{1} c_{2}\right) b_{2}^{-2}$. Obviously, $H(u, v)$ is nonnegative and $H(u, v)=0$ if and only if $(u, v)=(\bar{u}, \bar{v})$. By Theorem 1.2, $H(u, v)$ is well-posed for $t \geq 0$ if $(u, v)$ is positive solution to system 1.2 . The time derivative of $H(u, v)$ for system 1.2 satisfies

$$
\begin{aligned}
& \frac{d H(u, v)}{d t} \\
& =\int_{\Omega}\left(\frac{u-\bar{u}}{u} u_{t}+\rho \frac{v-\bar{v}}{v} v_{t}\right) d x \\
& =\int_{\Omega}\left\{\frac{u-\bar{u}}{u} \nabla \cdot\left[\left(d_{1}+2 \alpha_{11} u\right) \nabla u\right]+(u-\bar{u})\left(a_{1}-b_{1} u-c_{1} v\right)\right. \\
& \left.\quad+\rho \frac{v-\bar{v}}{v} \nabla \cdot\left[\left(d_{2}+\alpha_{21} u+2 \alpha_{22} v\right) \nabla v+\alpha_{21} v \nabla u\right]+\rho(v-\bar{v})\left(a_{2}+b_{2} u-c_{2} v\right)\right\} d x \\
& =-\int_{\Omega}\left[\frac{\left(d_{1}+2 \alpha_{11} u\right) \bar{u}}{u^{2}}|\nabla u|^{2}+\frac{\rho \alpha_{21} \bar{v}}{v} \nabla u \cdot \nabla v+\frac{\rho\left(d_{2}+\alpha_{21} u+2 \alpha_{22} v\right) \bar{v}}{v^{2}}|\nabla v|^{2}\right] d x \\
& \quad-\int_{\Omega}\left[b_{1}(u-\bar{u})^{2}+\left(c_{1}-\rho b_{2}\right)(u-\bar{u})(v-\bar{v})+c_{2} \rho(v-\bar{v})^{2}\right] d x .
\end{aligned}
$$

The second integrand in the above equality is positive definite by the choice of $\rho$. Meanwhile the first integrand is positive semi-definite if

$$
\begin{equation*}
4 \rho \overline{u v}\left(d_{1}+2 \alpha_{22} u\right)\left(d_{2}+\alpha_{21} u+2 \alpha_{22} v\right)>u^{2}\left(\alpha_{21} \bar{v}\right)^{2} . \tag{5.1}
\end{equation*}
$$

By the Lemma 2.1 and Theorem 1.2 , the condition $\sqrt{1.4}$ implies (5.1). Therefore, when all conditions in Theorem 1.3 hold, there exists positive constant $\delta$ depending on $b_{1}, b_{2}, c_{1}$ and $c_{2}$ such that

$$
\begin{equation*}
\frac{d H(u, v)}{d t} \leq-\delta \int_{\Omega}\left[(u-\bar{u})^{2}+(v-\bar{v})^{2}\right] d x \tag{5.2}
\end{equation*}
$$

To obtain the uniform convergence of the solution to 1.2 , we recall the following result which can be find in 21 .

Lemma 5.1. Let $a$ and $b$ positive constant. Assume that $\varphi, \psi \in C^{1}[a,+\infty)$, $\psi(t) \geq 0, \varphi$ is bounded. If $\varphi^{\prime}(t) \leq-b \psi(t)$ and $\psi^{\prime}(t)$ is bounded in $[a,+\infty)$, then $\lim _{t \rightarrow \infty} \psi(t)=0$.

Using integration by parts, Hölder's inequality, Lemma 2.1. and Lemma 4.2, one can easily verify that $\frac{d}{d t} \int_{\Omega}\left[(u-\bar{u})^{2}+(v-\bar{v})^{2}\right] d x$ is bounded from above. Then from Lemma 5.1 and 5.2 , we have

$$
\|u(\cdot, t)-\bar{u}\|_{L^{\infty}(\Omega)} \rightarrow 0, \quad\|v(\cdot, t)-\bar{v}\|_{L^{\infty}(\Omega)} \rightarrow 0 \quad(t \rightarrow \infty)
$$

Namely, $(u, v)$ converges uniformly to $(\bar{u}, \bar{v})$. By the fact that $H(u, v)$ is decreasing for $t \geq 0$, it is obvious that $(\bar{u}, \bar{v})$ is global asymptotic stable, and the proof of Theorem 1.3 is complete.

Acknowledgements. The author would like to thank professor Sheng-mao Fu for the encouragement and useful discussions, also the anonymous referee for the very careful reading of the original manuscript and the helpful suggestions.

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[^0]:    2000 Mathematics Subject Classification. 35K57, 35B35, 92D25.
    Key words and phrases. Cross-diffusion; global solution; gradient estimates; stability.
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    Submitted October 13, 2007. Published January 12, 2008.

