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# EXISTENCE AND UNIQUENESS FOR ONE-PHASE STEFAN PROBLEMS OF NON-CLASSICAL HEAT EQUATIONS WITH TEMPERATURE BOUNDARY CONDITION AT A FIXED FACE 

ADRIANA C. BRIOZZO, DOMINGO A. TARZIA


#### Abstract

We prove the existence and uniqueness, local in time, of a solution for a one-phase Stefan problem of a non-classical heat equation for a semiinfinite material with temperature boundary condition at the fixed face. We use the Friedman-Rubinstein integral representation method and the Banach contraction theorem in order to solve an equivalent system of two Volterra integral equations.


## 1. Introduction

The one-phase Stefan problem for a semi-infinite material for the classical heat equation requires the determination of the temperature distribution $u$ of the liquid phase (melting problem) or of the solid phase (solidification problem), and the evolution of the free boundary $x=s(t)$. Phase-change problems appear frequently in industrial processes and other problems of technological interest [2, 3, 6, 8, 5, 10 , 11, 12, 18, 29]. A large bibliography on the subject was given in [25].

Non-classical heat conduction problem for a semi-infinite material was studied in 4, 7, 17, 27, 28, e.g. problems of the type

$$
\begin{gather*}
u_{t}-u_{x x}=-F\left(u_{x}(0, t)\right), \quad x>0, t>0, \\
u(0, t)=0, \quad t>0  \tag{1.1}\\
u(x, 0)=h(x), \quad x>0
\end{gather*}
$$

where $h(x), x>0$, and $F(V), V \in \mathbb{R}$, are continuous functions. The function $F$, henceforth referred as control function, is assumed to satisfy the condition
$(\mathrm{H} 1) \quad F(0)=0$.
As observed in [27, 28, the heat flux $w(x, t)=u_{x}(x, t)$ for problem (1.1) satisfies a classical heat conduction problem with a nonlinear convective condition at $x=0$,

[^0]which can be written in the form
\[

$$
\begin{align*}
& w_{t}-w_{x x}=0, \quad x>0, \quad t>0 \\
& w_{x}(0, t)=F(w(0, t)), \quad t>0  \tag{1.2}\\
& w(x, 0)=h^{\prime}(x) \geq 0, \quad x>0
\end{align*}
$$
\]

The literature concerning problem $\sqrt{1.2}$ has increased rapidly since the publication of the papers [19, 21, 22]. Related problems have been also studied; see for example [1, 14, 16]. In [26], a one-phase Stefan problem for a non-classical heat equation for a semi-infinite material was presented. There the free boundary problem consists in determining the temperature $u=u(x, t)$ and the free boundary $x=s(t)$ with a control function $F$ which depends on the evolution of the heat flux at the extremum $x=0$ is given by the conditions

$$
\begin{gather*}
u_{t}-u_{x x}=-F\left(u_{x}(0, t)\right), \quad 0<x<s(t), 0<t<T \\
u(0, t)=f(t) \geqslant 0, \quad 0<t<T \\
u(s(t), t)=0, \quad u_{x}(s(t), t)=-\dot{s}(t), \quad 0<t<T  \tag{1.3}\\
u(x, 0)=h(x) \geq 0, \quad 0 \leqslant x \leqslant b=s(0) \quad(b>0)
\end{gather*}
$$

The goal in this paper is to prove the existence and uniqueness, local in time, of a solution to the one-phase Stefan problem (1.3) for a non-classical heat equation with temperature boundary condition at the fixed face $x=0$. First, we prove that problem (1.3) is equivalent to a system of two Volterra integral equations (2.4)- (2.5) following the Friedman-Rubinstein's method given in [13, 23]. Then, we prove that the problem (2.4)-2.5) has a unique local solution by using the Banach contraction theorem.

## 2. Existence and Uniqueness of Solutions

We have the following equivalence for the existence of solutions to the nonclassical free boundary problem (1.3).

Theorem 2.1. The solution of the free-boundary problem (1.3) is

$$
\begin{align*}
u(x, t)= & \int_{0}^{b} G(x, t ; \xi, 0) h(\xi) d \xi+\int_{0}^{t} G_{\xi}(x, t ; 0, \tau) f(\tau) d \tau  \tag{2.1}\\
& +\int_{0}^{t} G(x, t ; s(\tau), \tau) v(\tau) d \tau-\iint_{D(t)} G(x, t ; \xi, \tau) F(V(\tau)) d \xi d \tau \\
& s(t)=b-\int_{0}^{t} v(\tau) d \tau \tag{2.2}
\end{align*}
$$

where $D(t)=\{(x, \tau): 0<x<s(\tau), 0<\tau<t\}$, with $f \in C^{1}[0, T), h \in C^{1}[0, b]$, $h(b)=0, h(0)=f(0), F$ is a Lipschitz function over $C^{0}[0, T]$, and the functions $v \in C^{0}[0, T], V \in C^{0}[0, T]$ defined by

$$
\begin{equation*}
v(t)=u_{x}(s(t), t), \quad V(t)=u_{x}(0, t) \tag{2.3}
\end{equation*}
$$

must satisfy the following system of Volterra integral equations

$$
\begin{align*}
v(t)= & 2 \int_{0}^{b} N(s(t), t ; \xi, 0) h^{\prime}(\xi) d \xi-2 \int_{0}^{t} N(s(t), t ; 0, \tau) \dot{f}(\tau) d \tau \\
& +2 \int_{0}^{t} G_{x}(s(t), t ; s(\tau), \tau) v(\tau) d \tau  \tag{2.4}\\
& +2 \int_{0}^{t}[N(s(t), t ; s(\tau), \tau)-N(s(t), t ; 0, \tau)] F(V(\tau)) d \tau \\
V(t)= & \int_{0}^{b} N(0, t ; \xi, 0) h^{\prime}(\xi) d \xi \\
& -\int_{0}^{t} N(0, t ; 0, \tau) \dot{f}(\tau) d \tau+\int_{0}^{t} G_{x}(0, t ; s(\tau), \tau) v(\tau) d \tau  \tag{2.5}\\
& +\int_{0}^{t}[N(0, t ; s(\tau), \tau)-N(0, t ; 0, \tau)] F(V(\tau)) d \tau
\end{align*}
$$

where $G, N$ are the Green and Neumann functions and $K$ is the fundamental solution of the heat equation, defined respectively by

$$
\begin{gathered}
G(x, t, \xi, \tau)=K(x, t, \xi, \tau)-K(-x, t, \xi, \tau), \\
N(x, t, \xi, \tau)=K(x, t, \xi, \tau)+K(-x, t, \xi, \tau), \\
K(x, t, \xi, \tau)= \begin{cases}\frac{1}{2 \sqrt{\pi(t-\tau)}} \exp \left(-\frac{(x-\xi)^{2}}{4(t-\tau)}\right) & t>\tau \\
0 & t \leq \tau\end{cases}
\end{gathered}
$$

where $s(t)$ is given by 2.2 ,
Proof. Let $u(x, t)$ be the solution to $\sqrt{1.3}$ ). We integrate, on the domain $D_{t, \varepsilon}=$ $\{(\xi, \tau): 0<\xi<s(\tau), \varepsilon<\tau<t-\varepsilon\}$, the Green identity

$$
\begin{equation*}
\left(G u_{\xi}-u G_{\xi}\right)_{\xi}-(G u)_{\tau}=G F\left(u_{\xi}(0, \tau)\right) \tag{2.6}
\end{equation*}
$$

Now we let $\varepsilon \rightarrow 0$, to obtain the following integral representation for $u(x, t)$,

$$
\begin{aligned}
u(x, t)= & \int_{0}^{b} G(x, t ; \xi, 0) h(\xi) d \xi+\int_{0}^{t} G_{\xi}(x, t ; 0, \tau) f(\tau) d \tau \\
& +\int_{0}^{t} G(x, t ; s(\tau), \tau) u_{\xi}(s(\tau), \tau) d \tau-\iint_{D(t)} G(x, t ; \xi, \tau) F\left(u_{\xi}(0, \tau)\right) d \xi d \tau
\end{aligned}
$$

From the definition of $v(t)$ and $V(t)$ by 2.3 , we obtain 2.1 and 2.2 . If we differentiate $u(x, t)$ in variable $x$ and we let $x \rightarrow 0^{+}$and $x \rightarrow s(t)$, by using the jump relations, we obtain the integral equations for $v$ and $V$ given by 2.4 and (2.5).

Conversely, the function $u(x, t)$ defined by 2.1 where $v$ and $V$ are the solutions of (2.4) and $\sqrt{2.5}$, satisfy the conditions 1.3 ) (i), (ii), (iv) and (v). In order to prove condition (1.3) (iii) we define $\psi(t)=u(s(t), t)$. Taking into account that $u$ satisfy the conditions 1.3 ) (i),(ii),(iv) and (v), if we integrate the Green identity 2.6 over
the domain $D_{t, \varepsilon},(\varepsilon>0)$ and we let $\varepsilon \rightarrow 0$ we obtain that

$$
\begin{aligned}
u(x, t)= & \int_{0}^{b} G(x, t ; \xi, 0) h(\xi) d \xi+\int_{0}^{t} G(x, t ; s(\tau), \tau) v(\tau) d \tau \\
& +\int_{0}^{t} \psi(\tau)\left[G_{x}(x, t ; s(\tau), \tau)-G(x, t ; s(\tau), \tau) v(\tau)\right] d \tau \\
& +\int_{0}^{t} G_{\xi}(x, t ; 0, \tau) f(\tau) d \tau-\iint_{D(t)} G(x, t ; \xi, \tau) F(V(\tau)) d \xi d \tau
\end{aligned}
$$

Then, if we compare this last expression with 2.1), we deduce that

$$
\begin{equation*}
M(x, t)=\int_{0}^{t} \psi(\tau)\left[G_{x}(x, t ; s(\tau), \tau)-G(x, t ; s(\tau), \tau) v(\tau)\right] d \tau \equiv 0 \tag{2.7}
\end{equation*}
$$

for $0<x<s(t), 0<t<\sigma$. We let $x \rightarrow s(t)$ in 2.7 and by using the jump relations we have that $\psi$ satisfy the integral equation

$$
\frac{1}{2} \psi(t)+\int_{0}^{t} \psi(\tau)\left[G_{x}(s(t), t ; s(\tau), \tau)-G(s(t), t ; s(\tau), \tau) v(\tau)\right] d \tau=0
$$

Then we deduce that

$$
\begin{aligned}
|\psi(t)| & \leq C \int_{0}^{t} \frac{|\psi(\tau)|}{\sqrt{t-\tau}} d \tau \\
& \leq C^{2} \int_{0}^{t} \frac{d \tau}{\sqrt{t-\tau}} \int_{0}^{\tau} \frac{|\psi(\eta)|}{\sqrt{\tau-\eta}} d \eta \\
& =C^{2} \int_{0}^{t}|\psi(\eta)| d \eta \int_{\eta}^{t} \frac{d \tau}{[(t-\tau)(\tau-\eta)]^{1 / 2}} \\
& =\pi C^{2} \int_{0}^{t}|\psi(\eta)| d \eta
\end{aligned}
$$

where $C=C(t)$; therefore by using the Gronwall inequality we have that $\psi(t)=0$ over $[0, \sigma]$.

Next, we use the Banach fixed point theorem in order to prove the local existence and uniqueness of solution $v, V \in C^{0}[0, \sigma]$ to the system of two Volterra integral equations $(2.4)-(2.5)$ where $\sigma$ is a positive small number. Consider the Banach space

$$
C_{M, \sigma}=\left\{\vec{w}=\binom{v}{V}: v, V:[0, \sigma] \rightarrow \mathbb{R}, \text { continuous, with }\|\vec{w}\|_{\sigma} \leq M\right\}
$$

with

$$
\|\vec{w}\|_{\sigma}:=\|v\|_{\sigma}+\|V\|_{\sigma}:=\max _{t \in[0, \sigma]}|v(t)|+\max _{t \in[0, \sigma]}|V(t)|
$$

We define $A: C_{M, \sigma} \longrightarrow C_{M, \sigma}$, such that

$$
\overrightarrow{\widetilde{w}}(t)=A(\vec{w}(t))=\binom{A_{1}(v(t), V(t))}{A_{2}(v(t), V(t))}
$$

where

$$
\begin{equation*}
A_{1}(v(t), V(t))=F_{0}(v(t))+2 \int_{0}^{t}[N(s(t), t, s(\tau), \tau)-N(s(t), t, 0, \tau)] F(V(\tau)) d \tau \tag{2.8}
\end{equation*}
$$

with

$$
\begin{aligned}
F_{0}(v(t))= & 2 \int_{0}^{b} N(s(t), t, \xi, 0) h^{\prime}(\xi) d \xi-2 \int_{0}^{t} N(s(t), t, 0, \tau) \dot{f}(\tau) d \tau \\
& +2 \int_{0}^{t} G_{x}(s(t), t, s(\tau), \tau) v(\tau) d \tau
\end{aligned}
$$

and

$$
\begin{align*}
A_{2}(v(t), V(t))= & \int_{0}^{b} N(0, t, \xi, 0) h^{\prime}(\xi) d \xi-\int_{0}^{t} N(0, t, 0, \tau) \dot{f}(\tau) d \tau \\
& +\int_{0}^{t} G_{x}(0, t, s(\tau), \tau) v(\tau) d \tau  \tag{2.9}\\
& +\int_{0}^{t}[N(0, t, s(\tau), \tau)-N(0, t, 0, \tau)] F(V(\tau)) d \tau
\end{align*}
$$

Lemma 2.2. If $v \in C^{0}[0, \sigma]$, $\max _{t \in[0, \sigma]}|v(t)| \leq M$ and $2 M \sigma \leq b$ then $s(t)$ defined by (2.2) satisfies

$$
|s(t)-s(\tau)| \leq M|t-\tau| \quad|s(t)-b| \leq \frac{b}{2}, \quad \forall t, \tau \in[0, \sigma]
$$

To prove the following Lemmas we need the inequality

$$
\begin{equation*}
\exp \left(\frac{-x^{2}}{\alpha(t-\tau)}\right) /(t-\tau)^{n / 2} \leq\left(\frac{n \alpha}{2 e x^{2}}\right)^{n / 2}, \quad \alpha, x>0, t>\tau, n \in \mathbb{N} \tag{2.10}
\end{equation*}
$$

Lemma 2.3. Let $\sigma \leq 1, M \geq 1, f \in C^{1}[0, T), h \in C^{1}[0, b], F$ a Lipschitz function over $C^{0}[0, T]$. Under the hypothesis of Lemma 2.2, we have the following properties:

$$
\begin{gather*}
\int_{0}^{t}\left|N(s(t), t, 0, \tau)\|\dot{f}(\tau) \mid d \tau \leq\| \dot{f} \|_{t} C_{1}(b) t\right.  \tag{2.11}\\
\int_{0}^{t}\left|G_{x}(s(t), t, s(\tau), \tau) \| v(\tau)\right| d \tau \leq M^{2} C_{2}(b) \sqrt{t}  \tag{2.12}\\
\int_{0}^{b}\left|N(s(t), t, \xi, 0)\left\|h^{\prime}(\xi) \mid d \xi \leq\right\| h^{\prime} \|\right.  \tag{2.13}\\
\int_{0}^{t}|N(s(t), t, s(\tau), \tau)-N(s(t), t, 0, \tau) \| F(V(\tau))| d \tau \leq C_{4}(L) M \sqrt{t}  \tag{2.14}\\
\int_{0}^{b}\left|N(0, t, \xi, 0)\left\|h^{\prime}(\xi) \mid d \xi \leq\right\| h^{\prime} \|\right.  \tag{2.15}\\
\int_{0}^{t}|N(0, t, 0, \tau) \| \dot{f}(\tau)| d \tau \leq \frac{2\|\dot{f}\|_{\sigma}}{\sqrt{\pi}} \sqrt{t}  \tag{2.16}\\
\int_{0}^{t}\left|G_{x}(0, t, s(\tau), \tau) \| v(\tau)\right| d \tau \leq C_{3}(b) M t  \tag{2.17}\\
\int_{0}^{t}|N(0, t, s(\tau), \tau)-N(0, t, 0, \tau) \| F(V(\tau))| d \tau \leq C_{4}(L) M \sqrt{t} \tag{2.18}
\end{gather*}
$$

where $L$ is the Lipschitz constant of $F$ and

$$
\begin{gather*}
C_{1}(b)=\left(\frac{8}{e b^{2}}\right)^{1 / 2} \frac{1}{\sqrt{\pi}}, \quad C_{2}(b)=\frac{1}{2 \sqrt{\pi}}+\frac{3 b}{4 \sqrt{\pi}}\left(\frac{2}{3 e b^{2}}\right)^{3 / 2}  \tag{2.19}\\
C_{3}(b)=\frac{3 b}{8 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{3 / 2}, \quad C_{4}(L)=\frac{4 L}{\sqrt{\pi}} .
\end{gather*}
$$

Proof. To prove 2.11, we have

$$
\begin{aligned}
|N(s(t), t, 0, \tau)| & =|K(s(t), t, 0, \tau)+K(-s(t), t, 0, \tau)|=2 K(s(t), t, 0, \tau) \\
& =\exp \left(\frac{-s^{2}(t)}{4(t-\tau)}\right) \frac{(t-\tau)^{-1 / 2}}{\sqrt{\pi}} \\
& \leq \exp \left(\frac{-b^{2}}{16(t-\tau)}\right) \frac{(t-\tau)^{-1 / 2}}{\sqrt{\pi}} \\
& \leq\left(\frac{8}{e b^{2}}\right)^{1 / 2} \frac{1}{\sqrt{\pi}}=C_{1}(b)
\end{aligned}
$$

then 2.11 holds. To prove 2.12 , we have

$$
\begin{aligned}
\left|G_{x}(s(t), t, s(\tau), \tau)\right|= & \left|K_{x}(s(t), t, s(\tau), \tau)+K_{x}(-s(t), t, s(\tau), \tau)\right| \\
= & \frac{(t-\tau)^{-3 / 2}}{4 \sqrt{\pi}} \left\lvert\,(s(t)-s(\tau)) \exp \left(\frac{-(s(t)-s(\tau))^{2}}{4(t-\tau)}\right)\right. \\
& \left.-(s(t)+s(\tau)) \exp \left(\frac{-(s(t)+s(\tau))^{2}}{4(t-\tau)}\right) \right\rvert\, \\
\leq & \frac{(t-\tau)^{-3 / 2}}{4 \sqrt{\pi}}\left(M(t-\tau)+3 b \exp \left(\frac{-9 b^{2}}{4(t-\tau)}\right)\right) \\
\leq & \frac{1}{4 \sqrt{\pi}}\left(M(t-\tau)^{-1 / 2}+3 b\left(\frac{2}{3 e b^{2}}\right)^{3 / 2}\right)
\end{aligned}
$$

Then

$$
\begin{aligned}
\int_{0}^{t}\left|G_{x}(s(t), t, s(\tau), \tau)\right||v(\tau)| d \tau & \leq \frac{M}{4 \sqrt{\pi}}\left(2 M \sqrt{t}+3 b\left(\frac{2}{3 e b^{2}}\right)^{3 / 2} t\right) \\
& \leq M^{2} \sqrt{t}\left(\frac{1}{2 \sqrt{\pi}}+\frac{3 b}{M 4 \sqrt{\pi}}\left(\frac{2}{3 e b^{2}}\right)^{3 / 2}\right) \\
& \leq M^{2} C_{2}(b) \sqrt{t}
\end{aligned}
$$

which implies 2.12 . To prove 2.13 , we have

$$
\int_{0}^{b}\left|N(s(t), t, \xi, 0)\left\|h^{\prime}(\xi)\left|d \xi \leq\left\|h^{\prime}\right\| \int_{0}^{\infty}\right| N(s(t), t, \xi, 0) \mid d \xi \leq\right\| h^{\prime} \|\right.
$$

because

$$
\int_{0}^{\infty}|N(s(t), t, \xi, 0)| d \xi \leq 1
$$

To prove 2.14 , by taking into account that

$$
|N(s(t), t, s(\tau), \tau)-N(s(t), t, 0, \tau)| \leq \frac{2}{\sqrt{\pi(t-\tau)}}
$$

we obtain

$$
\begin{aligned}
\int_{0}^{t}|N(s(t), t, s(\tau), \tau)-N(s(t), t, 0, \tau)||F(V(\tau))| d \tau & \leq L M \int_{0}^{t} \frac{2}{\sqrt{\pi(t-\tau)}} d \tau \\
& =C_{4}(L) M \sqrt{t}
\end{aligned}
$$

The inequality 2.15 is prove in the same way as 2.13 . To prove 2.16, we have

$$
\begin{aligned}
\int_{0}^{t}|N(0, t, 0, \tau) \| \dot{f}(\tau)| d \tau & \leq\|\dot{f}\|_{\sigma} \int_{0}^{t}|N(0, t, 0, \tau)| d \tau \\
& =\|\dot{f}\|_{\sigma} \int_{0}^{t} \frac{1}{\sqrt{\pi(t-\tau)}} d \tau \\
& =\frac{\|\dot{f}\|_{\sigma}}{\sqrt{\pi}} 2 \sqrt{t}
\end{aligned}
$$

To prove 2.17), we have

$$
\begin{aligned}
\left|G_{x}(0, t, s(\tau), \tau)\right| & =\frac{(t-\tau)^{-3 / 2}}{4 \sqrt{\pi}} s(\tau) \exp \left(\frac{-(s(\tau))^{2}}{4(t-\tau)}\right) \\
& \leq \frac{3 b}{8 \sqrt{\pi}}(t-\tau)^{-3 / 2} \exp \left(\frac{-b^{2}}{16(t-\tau)}\right) \\
& \leq \frac{3 b}{8 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{3 / 2}
\end{aligned}
$$

To prove 2.18), as in 2.14, we prove that

$$
|N(0, t, s(\tau), \tau)-N(0, t, 0, \tau)| \leq \frac{2}{\sqrt{\pi(t-\tau)}}
$$

and therefore 2.18 holds.

Lemma 2.4. Let $s_{1}$, $s_{2}$ be the functions corresponding to $v_{1}$, $v_{2}$ in $C^{0}[0, \sigma]$, respectively, with $\max _{t \in[0, \sigma]}\left|v_{i}(t)\right| \leq M, i=1,2$, Then we have

$$
\begin{gather*}
\left|s_{2}(t)-s_{1}(t)\right| \leq t\left\|v_{2}-v_{1}\right\|_{t} \\
\left|s_{i}(t)-s_{i}(\tau)\right| \leq M|t-\tau|, \quad i=1,2  \tag{2.20}\\
\frac{b}{2} \leq s_{i}(t) \leq \frac{3 b}{2}, \quad \forall t \in[0, \sigma], \quad i=1,2
\end{gather*}
$$

Lemma 2.5. Let $f \in C^{1}[0, T), h \in C^{1}[0, b], F$ a Lipschitz function in $C^{0}[0, T]$. We have

$$
\begin{gather*}
\left|F_{0}\left(v_{2}(t)\right)-F_{0}\left(v_{1}(t)\right)\right| \leq E(b, h, f) \sqrt{t}\left\|v_{2}-v_{1}\right\|_{t}  \tag{2.21}\\
\int_{0}^{t}\left|N\left(s_{2}(t), t, s_{2}(\tau), \tau\right)-N\left(s_{2}(t), t, 0, \tau\right)\right|\left|F\left(V_{2}(\tau)\right)-F\left(V_{1}(\tau)\right)\right| d \tau  \tag{2.22}\\
\leq C_{4}(L) \sqrt{t}\left\|V_{2}-V_{1}\right\|_{t} ; \\
\int_{0}^{t}\left|N\left(s_{2}(t), t, 0, \tau\right)-N\left(s_{1}(t), t, 0, \tau\right)\right|\left|F\left(V_{1}(\tau)\right)\right| d \tau  \tag{2.23}\\
\leq C_{5}(b, L, M) t\left\|v_{2}-v_{1}\right\|_{t} ; \\
\int_{0}^{t}\left|N\left(s_{2}(t), t, s_{2}(\tau), \tau\right)-N\left(s_{1}(t), t, s_{1}(\tau), \tau\right)\right|\left|F\left(V_{1}(\tau)\right)\right| d \tau  \tag{2.24}\\
\leq\left[C_{6}(L, M) \sqrt{t}+C_{7}(b, L, M) t\right]\left\|v_{2}-v_{1}\right\|_{t} \\
\int_{0}^{t}\left|G_{x}\left(0, t, s_{2}(\tau), \tau\right)\right|\left|v_{2}(\tau)-v_{1}(\tau)\right| d \tau \leq C_{8}(b) t\left\|v_{2}-v_{1}\right\|_{t}  \tag{2.25}\\
\int_{0}^{t}\left|G_{x}\left(0, t, s_{2}(\tau), \tau\right) v_{2}(\tau)-G_{x}\left(0, t, s_{1}(\tau), \tau\right) v_{1}(\tau)\right| d \tau  \tag{2.26}\\
\leq\left(C_{8}(b) t+C_{9}(b, M) t^{2}\right)\left\|v_{2}-v_{1}\right\|_{t} ; \\
\int_{0}^{t} \mid\left[N\left(0, t, s_{2}(\tau), \tau\right)-N(0, t, 0, \tau)\right] F\left(V_{2}(\tau)\right) \\
-\left[N\left(0, t, s_{1}(\tau), \tau\right)-N(0, t, 0, \tau)\right] F\left(V_{1}(\tau)\right) \mid d \tau  \tag{2.27}\\
\leq C_{4}(L) \sqrt{t}\left\|V_{2}-V_{1}\right\|_{t}+C_{5}(b, L, M) t^{2}\left\|v_{2}-v_{1}\right\|_{t}
\end{gather*}
$$

where the constants are defined by

$$
\begin{gather*}
C_{4}(L)=\frac{4 L}{\sqrt{\pi}}, \quad C_{5}(b, L, M)=L M \frac{3 b}{8 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{3 / 2} \\
C_{6}(L, M)=\frac{L M^{3}}{\sqrt{\pi}}, \quad C_{7}(b, L, M)=\left(\frac{6}{e b^{2}}\right)^{3 / 2} \frac{3 b L M^{2}}{2 \sqrt{\pi}},  \tag{2.28}\\
C_{8}(b)=\frac{3}{4 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{3 / 2}, \quad C_{9}(b, M)=\left[\left(\frac{40}{e b^{2}}\right)^{\frac{5}{2}} \frac{9 b^{2}}{16 \sqrt{\pi}}+\frac{1}{2 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{3 / 2}\right] \frac{M}{2} .
\end{gather*}
$$

Proof. The proof of 2.21 can be found in [13]. To prove 2.22, we have

$$
\left|N\left(s_{2}(t), t, s_{2}(\tau), \tau\right)-N\left(s_{2}(t), t, 0, \tau\right)\right| \leq \frac{2}{\sqrt{\pi(t-\tau)}}
$$

Then

$$
\begin{aligned}
& \int_{0}^{t}\left|N\left(s_{2}(t), t, s_{2}(\tau), \tau\right)-N\left(s_{2}(t), t, 0, \tau\right)\right|\left|F\left(V_{2}(\tau)\right)-F\left(V_{1}(\tau)\right)\right| d \tau \\
& \leq \frac{4 L}{\sqrt{\pi}} \sqrt{t}\left\|V_{2}-V_{1}\right\|_{t}
\end{aligned}
$$

To prove 2.23, we use the mean value theorem: There exists $c=c(t, \tau)$ between $s_{1}(t)$ and $s_{2}(t)$ such that

$$
\begin{aligned}
& \left|N\left(s_{2}(t), t, 0, \tau\right)-N\left(s_{1}(t), t, 0, \tau\right)\right|\left|F\left(V_{1}(\tau)\right)\right| \\
& =\left|N_{x}(c, t, 0, \tau)\right|\left|s_{2}(\tau)-s_{1}(\tau)\right|\left|F\left(V_{1}(\tau)\right)\right| \\
& \leq|c| \exp \left(-\frac{c^{2}}{4(t-\tau)}\right) \frac{(t-\tau)^{-3 / 2}}{2 \sqrt{\pi}} L M \tau\left|v_{2}(\tau)-v_{1}(\tau)\right| \\
& \leq \frac{3 b}{4 \sqrt{\pi}} \exp \left(-\frac{b^{2}}{16(t-\tau)}\right)(t-\tau)^{-3 / 2} L M \tau\left|v_{2}(\tau)-v_{1}(\tau)\right| \\
& \leq \frac{3 b}{4 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{3 / 2} L M \tau\left|v_{2}(\tau)-v_{1}(\tau)\right|
\end{aligned}
$$

Then

$$
\begin{aligned}
& \int_{0}^{t}\left|N\left(s_{2}(t), t, 0, \tau\right)-N\left(s_{1}(t), t, 0, \tau\right) \| F\left(V_{1}(\tau)\right)\right| d \tau \\
& \leq \frac{3 b}{8 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{\frac{3}{2}} L M t\left\|v_{2}-v_{1}\right\|_{t}=C_{5}(b, L, M) t\left\|v_{2}-v_{1}\right\|_{t}
\end{aligned}
$$

To prove 2.24 , we have

$$
\begin{aligned}
& N\left(s_{2}(t), t, s_{2}(\tau), \tau\right)-N\left(s_{1}(t), t, s_{1}(\tau), \tau\right) \\
& =K\left(s_{2}(t), t, s_{2}(\tau), \tau\right)-K\left(s_{1}(t), t, s_{1}(\tau), \tau\right) \\
& \quad+K\left(-s_{2}(t), t, s_{2}(\tau), \tau\right)-K\left(-s_{1}(t), t, s_{1}(\tau), \tau\right)
\end{aligned}
$$

As in [24], for each $(t, \tau), 0<\tau<t$, we define

$$
f_{t, \tau}(x)=\exp \left(\frac{-x^{2}}{4(t-\tau)}\right)
$$

Then we have

$$
\begin{aligned}
& K\left(s_{2}(t), t, s_{2}(\tau), \tau\right)-K\left(s_{1}(t), t, s_{1}(\tau), \tau\right) \\
& =\frac{(t-\tau)^{-1 / 2}}{2 \sqrt{\pi}}\left[\exp \left(-\frac{\left(s_{2}(t)-s_{2}(\tau)\right)^{2}}{4(t-\tau)}\right)-\exp \left(-\frac{\left(s_{1}(t)-s_{1}(\tau)\right)^{2}}{4(t-\tau)}\right)\right] \\
& =\frac{(t-\tau)^{-1 / 2}}{2 \sqrt{\pi}}\left[f_{t, \tau}\left(s_{2}(t)-s_{2}(\tau)\right)-f_{t, \tau}\left(s_{1}(t)-s_{1}(\tau)\right)\right]
\end{aligned}
$$

and

$$
\begin{aligned}
& K\left(-s_{2}(t), t, s_{2}(\tau), \tau\right)-K\left(-s_{1}(t), t, s_{1}(\tau), \tau\right) \\
& =\frac{(t-\tau)^{-1 / 2}}{2 \sqrt{\pi}}\left[\exp \left(-\frac{\left(s_{2}(t)+s_{2}(\tau)\right)^{2}}{4(t-\tau)}\right)-\exp \left(-\frac{\left(s_{1}(t)+s_{1}(\tau)\right)^{2}}{4(t-\tau)}\right)\right] \\
& =\frac{(t-\tau)^{-1 / 2}}{2 \sqrt{\pi}}\left[f_{t, \tau}\left(s_{2}(t)+s_{2}(\tau)\right)-f_{t, \tau}\left(s_{1}(t)+s_{1}(\tau)\right)\right]
\end{aligned}
$$

By the mean value theorem there exists $c=c(t, \tau)$ between $s_{2}(t)-s_{2}(\tau)$ and $s_{1}(t)-s_{1}(\tau)$ such that

$$
\begin{aligned}
& f_{t, \tau}\left(s_{2}(t)-s_{2}(\tau)\right)-f_{t, \tau}\left(s_{1}(t)-s_{1}(\tau)\right) \\
& =f_{t, \tau}^{\prime}(c)\left(s_{2}(t)-s_{2}(\tau)-s_{1}(t)+s_{1}(\tau)\right) \\
& =\frac{-c}{2(t-\tau)} \exp \left(-\frac{c^{2}}{4(t-\tau)}\right)\left(s_{2}(t)-s_{2}(\tau)-s_{1}(t)+s_{1}(\tau)\right)
\end{aligned}
$$

Taking into account that

$$
|c| \leq \max \left\{\left|s_{i}(t)-s_{i}(\tau)\right|, i=1,2\right\} \leq M(t-\tau)
$$

it results

$$
\begin{aligned}
\left|f_{t, \tau}\left(s_{2}(t)-s_{2}(\tau)\right)-f_{t, \tau}\left(s_{1}(t)-s_{1}(\tau)\right)\right| & \leq \frac{M}{2}\left[\mid s_{2}(t)\right)-s_{1}(t)\left|+\left|s_{2}(\tau)-s_{1}(\tau)\right|\right] \\
& \leq M^{2}\left\|v_{2}-v_{1}\right\|_{t}
\end{aligned}
$$

Then we have

$$
\left|K\left(s_{2}(t), t, s_{2}(\tau), \tau\right)-K\left(s_{1}(t), t, s_{1}(\tau), \tau\right)\right| \leq \frac{M^{2}}{2 \sqrt{\pi(t-\tau)}}\left\|v_{2}-v_{1}\right\|_{t}
$$

In the same way we have

$$
\begin{aligned}
& f_{t, \tau}\left(s_{2}(t)+s_{2}(\tau)\right)-f_{t, \tau}\left(s_{1}(t)+s_{1}(\tau)\right) \\
& =f_{t, \tau}^{\prime}\left(c^{*}\right)\left(s_{2}(t)+s_{2}(\tau)-s_{1}(t)-s_{1}(\tau)\right) \\
& =\frac{-c^{*}}{2(t-\tau)} \exp \left(-\frac{c^{* 2}}{4(t-\tau)}\right)\left(s_{2}(t)+s_{2}(\tau)-s_{1}(t)-s_{1}(\tau)\right)
\end{aligned}
$$

where $c^{*}=c^{*}(t, \tau)$ is between $s_{2}(t)+s_{2}(\tau)$ and $s_{1}(t)+s_{1}(\tau)$. Since $s_{1}(t)+s_{1}(\tau) \leq$ $c^{*} \leq s_{2}(t)+s_{2}(\tau)$, (or viceversa), we deduce that $b \leq c^{*} \leq 3 b$, that is $\exp \left(-c^{* 2} / 4(t-\right.$ $\tau)) \leq \exp \left(-b^{2} / 4(t-\tau)\right)$. Then we obtain

$$
\begin{aligned}
& \left|K\left(-s_{2}(t), t, s_{2}(\tau), \tau\right)-K\left(-s_{1}(t), t, s_{1}(\tau), \tau\right)\right| \\
& =\frac{(t-\tau)^{-1 / 2}}{2 \sqrt{\pi}}\left|f_{t, \tau}\left(s_{2}(t)+s_{2}(\tau)\right)-f_{t, \tau}\left(s_{1}(t)+s_{1}(\tau)\right)\right| \\
& \leq \frac{3 b}{4 \sqrt{\pi}(t-\tau)^{3 / 2}} \exp \left(-\frac{b^{2}}{4(t-\tau)}\right) 2 M\left\|v_{2}-v_{1}\right\|_{t} \\
& \leq\left(\frac{6}{e b^{2}}\right)^{3 / 2} \frac{3 b M}{2 \sqrt{\pi}}\left\|v_{2}-v_{1}\right\|_{t}
\end{aligned}
$$

and

$$
\begin{aligned}
& \left|N\left(s_{2}(t), t, s_{2}(\tau), \tau\right)-N\left(s_{1}(t), t, s_{1}(\tau), \tau\right)\right| \\
& \leq\left(\frac{M^{2}}{2 \sqrt{\pi(t-\tau)}}+\left(\frac{6}{e b^{2}}\right)^{3 / 2} \frac{3 b M}{2 \sqrt{\pi}}\right)\left\|v_{2}-v_{1}\right\|_{t}
\end{aligned}
$$

Therefore,

$$
\begin{aligned}
& \int_{0}^{t}\left|N\left(s_{2}(t), t, s_{2}(\tau), \tau\right)-N\left(s_{1}(t), t, s_{1}(\tau), \tau\right)\right|\left|F\left(V_{1}(\tau)\right)\right| d \tau \\
& \leq \int_{0}^{t}\left(\frac{M^{2}}{2 \sqrt{\pi(t-\tau)}}+\left(\frac{6}{e b^{2}}\right)^{3 / 2} \frac{3 b M}{2 \sqrt{\pi}}\right)\left\|v_{2}-v_{1}\right\|_{t}\left|F\left(V_{1}(\tau)\right)\right| d \tau \\
& \leq L M\left(\frac{M^{2} \sqrt{t}}{\sqrt{\pi}}+\left(\frac{6}{e b^{2}}\right)^{3 / 2} \frac{3 b M}{2 \sqrt{\pi}} t\right)\left\|v_{2}-v_{1}\right\|_{t} \\
& =\left(C_{6}(L, M) \sqrt{t}+C_{7}(L, M, b) t\right)\left\|v_{2}-v_{1}\right\|_{t}
\end{aligned}
$$

To prove 2.25, we take into account 2.10 :

$$
\begin{aligned}
G_{x}\left(0, t, s_{2}(\tau), \tau\right) & =K\left(0, t, s_{2}(\tau), \tau\right) \frac{s_{2}(\tau)}{t-\tau} \\
& =\exp \left(-\frac{s_{2}^{2}(\tau)}{4(t-\tau)}\right) \frac{(t-\tau)^{-3 / 2}}{2 \sqrt{\pi}} s_{2}(\tau) \\
& \leq \frac{1}{2 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{3 / 2} s_{2}(\tau) \leq \frac{3 b}{4 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{\frac{3}{2}}=C_{8}(b)
\end{aligned}
$$

To prove (2.26), we have

$$
\begin{aligned}
& \left|G_{x}\left(0, t, s_{2}(\tau), \tau\right) v_{2}(\tau)-G_{x}\left(0, t, s_{1}(\tau), \tau\right) v_{1}(\tau)\right| \\
& \quad \leq\left|G_{x}\left(0, t, s_{2}(\tau), \tau\right)\right|\left|v_{2}(\tau)-v_{1}(\tau)\right| \\
& \quad+\left|G_{x}\left(0, t, s_{2}(\tau), \tau\right)-G_{x}\left(0, t, s_{1}(\tau), \tau\right)\right|\left|v_{1}(\tau)\right| .
\end{aligned}
$$

Using the mean value theorem there exists $c=c(\tau)$ between $s_{2}(\tau)$ and $s_{1}(\tau)$ such that $G_{x}\left(0, t, s_{2}(\tau), \tau\right)-G_{x}\left(0, t, s_{1}(\tau), \tau\right)=G_{x \xi}(0, t, c, \tau)\left(s_{2}(\tau)-s_{1}(\tau)\right)$. Taking into account the following properties

$$
\begin{gathered}
G_{x \xi}(0, t, c, \tau)=\frac{K(0, t, c, \tau)}{t-\tau}\left(\frac{c^{2}}{2(t-\tau)}+1\right), \\
\frac{K(0, t, c, \tau)}{t-\tau}=\frac{1}{2 \sqrt{\pi}} \exp \left(-\frac{c^{2}}{4(t-\tau)}\right)(t-\tau)^{-3 / 2} \leq \frac{1}{2 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{3 / 2}, \\
K(0, t, c, \tau) \frac{c^{2}}{2(t-\tau)^{2}}=\frac{1}{4 \sqrt{\pi}} \exp \left(-\frac{c^{2}}{4(t-\tau)}\right)(t-\tau)^{-\frac{5}{2}} c^{2} \leq \frac{9 b^{2}}{16 \sqrt{\pi}}\left(\frac{40}{e b^{2}}\right)^{5 / 2}
\end{gathered}
$$

we have

$$
\begin{aligned}
& \left|G_{x}\left(0, t, s_{2}(\tau), \tau\right)-G_{x}\left(0, t, s_{1}(\tau), \tau\right)\right|\left|v_{1}(\tau)\right| \\
& \leq\left(\frac{1}{2 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{\frac{3}{2}}+\frac{9 b^{2}}{16 \sqrt{\pi}}\left(\frac{40}{e b^{2}}\right)^{\frac{5}{2}}\right)\left|s_{2}(\tau)-s_{1}(\tau)\right|\left|v_{1}(\tau)\right| \\
& \leq M\left(\frac{1}{2 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{\frac{3}{2}}+\frac{9 b^{2}}{16 \sqrt{\pi}}\left(\frac{40}{e b^{2}}\right)^{\frac{5}{2}}\right) \tau\left|v_{2}(\tau)-v_{1}(\tau)\right|
\end{aligned}
$$

Then

$$
\begin{aligned}
& \int_{0}^{t}\left|G_{x}\left(0, t, s_{2}(\tau), \tau\right)-G_{x}\left(0, t, s_{1}(\tau), \tau\right)\right|\left|v_{1}(\tau)\right| d \tau \\
& \leq\left(\frac{1}{2 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{\frac{3}{2}}+\frac{9 b^{2}}{16 \sqrt{\pi}}\left(\frac{40}{e b^{2}}\right)^{\frac{5}{2}}\right) \frac{M t^{2}}{2}\left\|v_{2}-v_{1}\right\|_{t}
\end{aligned}
$$

Then (2.26) holds by using (2.25). To prove 2.27), we have

$$
\begin{align*}
& {\left[N\left(0, t, s_{2}(\tau), \tau\right)-N(0, t, 0, \tau)\right] F\left(V_{2}(\tau)\right)} \\
& -\left[N\left(0, t, s_{1}(\tau), \tau\right)-N(0, t, 0, \tau)\right] F\left(V_{1}(\tau)\right) \\
& =\left[N\left(0, t, s_{2}(\tau), \tau\right)-N(0, t, 0, \tau)\right]\left[F\left(V_{2}(\tau)\right)-F\left(V_{1}(\tau)\right)\right]  \tag{2.29}\\
& \quad+\left[N\left(0, t, s_{2}(\tau), \tau\right)-N\left(0, t, s_{1}(\tau), \tau\right)\right] F\left(V_{1}(\tau)\right)
\end{align*}
$$

Using $\left|N\left(0, t, s_{2}(\tau), \tau\right)-N(0, t, 0, \tau)\right| \leq \frac{2}{\sqrt{\pi(t-\tau)}}$ we get
$\left|N\left(0, t, s_{2}(\tau), \tau\right)-N(0, t, 0, \tau)\right|\left|F\left(V_{2}(\tau)\right)-F\left(V_{1}(\tau)\right)\right| \leq \frac{2}{\sqrt{\pi(t-\tau)}} L\left|V_{2}(\tau)-V_{1}(\tau)\right|$,
and

$$
\begin{align*}
& \int_{0}^{t}\left|N\left(0, t, s_{2}(\tau), \tau\right)-N(0, t, 0, \tau) \| F\left(V_{2}(\tau)\right)-F\left(V_{1}(\tau)\right)\right| d \tau \\
& \leq \frac{4 \sqrt{t}}{\sqrt{\pi}} L\left\|V_{2}-V_{1}\right\|_{t}=C_{4}(L) \sqrt{t}\left\|V_{2}-V_{1}\right\|_{t} \tag{2.30}
\end{align*}
$$

Furthermore,

$$
\left|N\left(0, t, s_{2}(\tau), \tau\right)-N\left(0, t, s_{1}(\tau), \tau\right)\right|=\left|N_{\xi}(0, t, c, \tau)\right|\left|s_{2}(\tau)-s_{1}(\tau)\right|
$$

where $c=c(\tau)$ is between $s_{2}(\tau)$ and $s_{1}(\tau)$ and

$$
\begin{aligned}
\left|N_{\xi}(0, t, c, \tau)\right|\left|s_{2}(\tau)-s_{1}(\tau)\right| & =\left|-G_{x}(0, t, c, \tau)\right|\left|s_{2}(\tau)-s_{1}(\tau)\right| \\
& \leq \frac{|c|}{2 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{3 / 2} \tau\left|v_{2}(\tau)-v_{1}(\tau)\right| \\
& \leq \frac{3 b}{4 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{3 / 2} \tau\left|v_{2}(\tau)-v_{1}(\tau)\right|
\end{aligned}
$$

Then

$$
\begin{align*}
& \int_{0}^{t}\left|N\left(0, t, s_{2}(\tau), \tau\right)-N\left(0, t, s_{1}(\tau), \tau\right) \| F\left(V_{1}(\tau)\right)\right| d \tau \\
& \leq L M \frac{3 b}{4 \sqrt{\pi}}\left(\frac{24}{e b^{2}}\right)^{3 / 2} \frac{t^{2}}{2}\left\|v_{2}-v_{1}\right\|_{t}=C_{5}(L, M, b) t^{2}\left\|v_{2}-v_{1}\right\|_{t} \tag{2.31}
\end{align*}
$$

Therefore, by 2.29, 2.30, and 2.31), the inequality 2.27 holds.
Theorem 2.6. The map $A: C_{M, \sigma} \rightarrow C_{M, \sigma}$ is well defined and is a contraction map if $\sigma$ satisfies the following inequalities:

$$
\begin{gather*}
\sigma \leq 1,2 M \sigma \leq b  \tag{2.32}\\
\left(2\|\dot{f}\|_{\sigma} C_{1}(b)+M C_{3}(b)\right) \sigma+\left(2 M^{2} C_{2}(b)+\frac{2\|\dot{f}\|_{\sigma}}{\sqrt{\pi}}+3 M C_{4}(L)\right) \sqrt{\sigma} \leq 1  \tag{2.33}\\
D(b, f, h, L, M) \sqrt{\sigma}<1 \tag{2.34}
\end{gather*}
$$

where

$$
\begin{equation*}
M=1+3\left\|h^{\prime}\right\| \tag{2.35}
\end{equation*}
$$

and

$$
\begin{gathered}
D_{1}(b, f, h, L, M)=E(b, f, h)+2 C_{6}(L, M)+3 C_{4}(L) \\
D_{2}(b, L, M)=2\left[C_{5}(b, L, M)+2 C_{7}(b, L, M)+C_{8}(b)\right] \\
D_{3}(b, L, M)=C_{9}(b, M)+C_{5}(b, L, M) \\
D(b, f, h, L, M)=D_{1}(b, f, h, L, M)+D_{2}(b, L, M)+D_{3}(b, L, M)
\end{gathered}
$$

Then there exists a unique solution on $C_{M, \sigma}$ to the system of integral equations (2.4), 2.5).

Proof. Firstly we demonstrate that $A$ maps $C_{\sigma, M}$ into itself, that is

$$
\begin{equation*}
\|A(\vec{w})\|_{\sigma}=\max _{t \in[0, \sigma]}\left|A_{1}(v(t), V(t))\right|+\max _{t \in[0, \sigma]}\left|A_{2}(v(t), V(t))\right| \leq M \tag{2.36}
\end{equation*}
$$

Using the Lemmas $2.3,2.4$ and the definitions $(2.8)-(2.9)$, we have

$$
\begin{gathered}
\left|A_{1}(v(t), V(t))\right| \leq 2\|\dot{f}\|_{\sigma} C_{1}(b) t+2 M^{2} C_{2}(b) \sqrt{t}+2\left\|h^{\prime}\right\|+2 C_{4}(L) M \sqrt{t} \\
\left|A_{2}(v(t), V(t))\right| \leq\left\|h^{\prime}\right\|+\left(\frac{2\|\dot{f}\|_{\sigma}}{\sqrt{\pi}}+C_{4}(L) M\right) \sqrt{t}+C_{3}(b) M t
\end{gathered}
$$

Then

$$
\begin{aligned}
\|A(\vec{w})\|_{\sigma}= & \max _{t \in[0, \sigma]}\left|A_{1}(v(t), V(t))\right|+\max _{t \in[0, \sigma]}\left|A_{2}(v(t), V(t))\right| \\
\leq & 3\left\|h^{\prime}\right\|+\left(2\|\dot{f}\|_{\sigma} C_{1}(b)+C_{3}(b) M\right) \sigma \\
& +\left(2 M^{2} C_{2}(b)+\frac{2\|\dot{f}\|_{\sigma}}{\sqrt{\pi}}+3 M C_{4}(L)\right) \sqrt{\sigma} .
\end{aligned}
$$

Selecting $M$ by 2.35 and $\sigma$ such that 2.32 and 2.33 hold, we obtain 2.36).
Now, we prove that

$$
\left\|A\left(\overrightarrow{w_{2}}\right)-A\left(\overrightarrow{w_{1}}\right)\right\|_{\sigma} \leq D(b, h, f, L, M) \sqrt{\sigma}\left\|\overrightarrow{w_{2}}-\overrightarrow{w_{1}}\right\|_{\sigma}
$$

where $\overrightarrow{w_{1}}=\binom{v_{1}}{V_{1}}, \overrightarrow{w_{2}}=\binom{v_{2}}{V_{2}}$. By selecting $\sigma$ such that 2.34 holds, $A$ becomes a contraction mapping on $C_{\sigma, M}$ and therefore it has a unique fixed point. To prove this assertion we consider

$$
A\left(\vec{w}_{1}\right)(t)-A\left(\overrightarrow{w_{2}}\right)(t)=\binom{A_{1}\left(v_{2}(t), V_{2}(t)\right)-A_{1}\left(v_{1}(t), V_{1}(t)\right)}{A_{2}\left(v_{2}(t), V_{2}(t)\right)-A_{2}\left(v_{1}(t), V_{1}(t)\right)}
$$

where

$$
\begin{aligned}
& A_{1}\left(v_{2}(t), V_{2}(t)\right)-A_{1}\left(v_{1}(t), V_{1}(t)\right) \\
& =F_{0}\left(v_{2}(t)\right)-F_{0}\left(v_{1}(t)\right)+2 \int_{0}^{t}\left[N\left(s_{2}(t), t, s_{2}(\tau), \tau\right)-N\left(s_{2}(t), t, 0, \tau\right)\right] F\left(V_{2}(\tau)\right) d \tau \\
& \quad-2 \int_{0}^{t}\left[N\left(s_{1}(t), t, s_{1}(\tau), \tau\right)-N\left(s_{1}(t), t, 0, \tau\right)\right] F\left(V_{1}(\tau)\right) d \tau
\end{aligned}
$$

and

$$
\begin{aligned}
A_{2} & \left(v_{2}(t), V_{2}(t)\right)-A_{2}\left(v_{1}(t), V_{1}(t)\right) \\
= & \int_{0}^{t}\left[G_{x}\left(0, t, v_{2}(\tau), \tau\right) v_{2}(\tau)-G_{x}\left(0, t, v_{1}(\tau), \tau\right) v_{1}(\tau)\right] d \tau \\
& +\int_{0}^{t}\left\{\left[N\left(0, t, s_{2}(\tau), \tau\right)-N(0, t, 0, \tau)\right] F\left(V_{2}(\tau)\right)\right. \\
& \left.-\left[N\left(0, t, s_{1}(\tau), \tau\right)-N(0, t, 0, \tau)\right] F\left(V_{1}(\tau)\right)\right\} d \tau
\end{aligned}
$$

Taking into account the Lemmas 2.4 and 2.5 it results

$$
\begin{aligned}
& \left|A_{1}\left(v_{2}(t), V_{2}(t)\right)-A_{1}\left(v_{1}(t), V_{1}(t)\right)\right| \\
& \quad \leq E(b, h, f) \sqrt{t}\left\|v_{2}-v_{1}\right\|_{t}+2 C_{4}(L) \sqrt{t}\left\|V_{2}-V_{1}\right\|_{t} \\
& \quad+2 C_{5}(b, L, M) t\left\|v_{2}-v_{1}\right\|_{t}+2\left[C_{6}(L, M) \sqrt{t}+C_{7}(b, L, M) t\right]\left\|v_{2}-v_{1}\right\|_{t}
\end{aligned}
$$

and

$$
\begin{aligned}
& \left|A_{2}\left(v_{2}(t), V_{2}(t)\right)-A_{2}\left(v_{1}(t), V_{1}(t)\right)\right| \\
& \leq\left(C_{8}(b) t+C_{9}(b, M) t^{2}\right)\left\|v_{2}-v_{1}\right\|_{t} \\
& \quad+C_{4}(L) \sqrt{t}\left\|V_{2}-V_{1}\right\|_{t}+C_{5}(b, L, M) t^{2}\left\|v_{2}-v_{1}\right\|_{t}
\end{aligned}
$$

Therefore,

$$
\begin{aligned}
& \left\|A\left(\overrightarrow{w_{2}}\right)-A\left(\overrightarrow{w_{1}}\right)\right\|_{\sigma} \\
& \leq \max _{t \in[0, \sigma]}\left|A_{1}\left(v_{2}(t), V_{2}(t)\right)-A_{1}\left(v_{1}(t), V_{1}(t)\right)\right| \\
& \quad+\max _{t \in[0, \sigma]}\left|A_{2}\left(v_{2}(t), V_{2}(t)\right)-A_{2}\left(v_{1}(t), V_{1}(t)\right)\right| \\
& \leq\left\{D_{1}(b, f, h, L, M) \sqrt{\sigma}+D_{2}(b, L, M) \sigma+D_{3}(b, L, M) \sigma^{2}\right\}\left\|\overrightarrow{w_{2}}-\overrightarrow{w_{1}}\right\|_{\sigma} \\
& \leq \\
& D(b, f, h, L, M) \sqrt{\sigma}\left\|\overrightarrow{w_{2}}-\overrightarrow{w_{1}}\right\|_{\sigma}
\end{aligned}
$$

By hypothesis 2.34 we have that $A$ is a contraction.
Remark. If $F$ satisfies the conditions
(H2) $F(V)>0$, for all $V \neq 0$ and $F(0)=0$,
then by the maximum principle [5], $u$ is a sub-solution for the same problem with $F \equiv 0$, that is

$$
u(x, t) \leq u_{0}(x, t), \quad s(t) \leq s_{0}(t)
$$

where $u_{0}(x, t)$ and $s_{0}(t)$ solve the classical Stefan problem

$$
\begin{gathered}
u_{0 t}-u_{0 x x}=0, \quad 0<x<s_{0}(t), 0<t<T \\
u_{0}(0, t)=f(t) \geqslant 0, \quad 0<t<T \\
u_{0}\left(s_{0}(t), t\right)=0, \quad u_{0} x\left(s_{0}(t), t\right)=-\dot{s_{0}}(t) .0<t<T, \\
u_{0}(x, 0)=h(x) \quad 0 \leqslant x \leqslant b=s_{0}(0)
\end{gathered}
$$

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## References

[1] E. K. Afenya and C. P. Calderon, A remark on a nonlinear integral equation, Rev. Un. Mat. Argentina, 39 (1995), 223-227.
[2] V. Alexiades and A. D. Solomon, Mathematical modeling of melting and freezing processes, Hemisphere - Taylor \& Francis, Washington (1983).
[3] I. Athanasopoulos, G. Makrakis, and J. F. Rodrigues (Eds.), Free Boundary Problems: Theory and Applications, CRC Press, Boca Raton (1999).
[4] L. R. Berrone, D. A. Tarzia, and L. T. Villa, Asymptotic behavior of a non-classical heat conduction problem for a semi-infinite material, Math. Meth. Appl. Sci., 23 (2000), 11611177.
[5] M. Bertsch, A class of degenerate diffusion equations with a singular nonlinear term, Nonlinear Anal. Th. Meth. Appl., 7(1983), 117-127.
[6] J. R. Cannon, The one-dimensional heat equation, Addison-Wesley, Menlo Park (1984).
[7] J. R. Cannon and H. M. Yin, A class of non-linear non-classical parabolic equations, J. Diff. Eq., 79 (1989), 266-288.
[8] H. S. Carslaw and J. C. Jaeger, Conduction of heat in solids, Oxford University Press, London (1959).
[9] J. M. Chadam and H. Rasmussen H (Eds.), Free boundary problems involving solids, Pitman Research Notes in Mathematics Series \#281, Longman Essex (1993).
[10] J. Crank, Free and moving boundary problems, Clarendon Press Oxford (1984).
[11] J. I. Diaz, M. A. Herrero, A. Liñan and J. L. Vazquez (Eds.), "Free boundary problems: theory and applications", Pitman Research Notes in Mathematics Series \#323, Longman, Essex (1995).
[12] A. Fasano and M. Primicerio (Eds.), "Nonlinear diffusion problems", Lecture Notes in Math N. 1224, Springer Verlag, Berlin (1986).
[13] A. Friedman, Free boundary problems for parabolic equations I. Melting of solids, J. Math. Mech., 8 (1959), 499-517.
[14] K. Glashoff and J. Sprekels, The regulation of temperature by thermostats and set-valued integral equations, J. Integral Eq., 4 (1982), 95-112.
[15] G. Gripenberg, S.O. Londen, and O. Staffans, Volterra integral and functional equations, Cambridge Univ. Press, Cambridge (1990).
[16] R. A. Handelsman and W. E. Olmstead, Asymptotic solution to a class of nonlinear Volterra integral equations, SIAM J. Appl. Math., 22 (1972), 373-384.
[17] N. Kenmochi and M. Primicerio, One-dimensional heat conduction with a class of automatic heat source controls, IMA J. Appl. Math. 40 (1988), 205-216.
[18] V. J. Lunardini, Heat transfer with freezing and thawing, Elsevier, Amsterdam (1991).
[19] W. R. Mann and F. Wolf, Heat transfer between solid and gases under nonlinear boundary conditions, Quart. Appl. Math., 9 (1951), 163-184.
[20] R. K. Miller, Nonlinear Volterra integral equations, W. A. Benjamin, Menlo Park (1971).
[21] W. E. Olmstead and R. A. Handelsman, Diffusion in a semi-infinite region with nonlinear surface dissipation, SIAM Rev., 18 (1976), 275-291.
[22] J. H. Roberts and W. R. Mann, A certain nonlinear integral equation of the Volterra type, Pacific J. Math., 1 (1951), 431-445.
[23] L. I. Rubinstein, The Stefan problem, Trans. Math. Monographs \# 27, Amer. Math. Soc., Providence (1971).
[24] B. Sherman, A free boundary problem for the heat equation with prescribed flux at both fixed face and melting interface, Quart. Appl. Math., 25 (1967), 53-63,
[25] D. A. Tarzia, A bibliography on moving-free boundary problems for the heat-diffusion equation. The Stefan and related problems, MAT-Serie A, Rosario, \# 2 (2000), (with 5869 titles on the subject, 300 pages). See www.austral.edu.ar/MAT-SerieA/2(2000)/
[26] D. A. Tarzia, A Stefan problem for a non-classical heat equation, MAT-Serie A, Rosario, \# 3 (2001), 21-26.
[27] D. A. Tarzia and L. T. Villa, Some nonlinear heat conduction problems for a semi-infinite strip with a non-uniform heat source, Rev. Un. Mat. Argentina, 41 (2000), 99-114.
[28] L. T. Villa, Problemas de control para una ecuación unidimensional del calor, Rev. Un. Mat. Argentina, 32 (1986), 163-169.
[29] L. C. Wrobel and C. A. Brebbia (Eds.), Computational methods for free and moving boundary problems in heat and fluid flow, Computational Mechanics Publications, Southampton (1993).

Adriana C. Briozzo
Departamento de Matemática, FCE, Universidad Austral, Paraguay 1950, S2000FZF
Rosario, Argentina
E-mail address: Adriana.Briozzo@fce.austral.edu.ar
Domingo Alberto Tarzia
Departamento de Matemática - CONICET, FCE, Universidad Austral, Paraguay 1950, S2000FZF Rosario, Argentina

E-mail address: Domingo.Tarzia@fce.austral.edu.ar


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