# SIGN-CHANGING SOLUTIONS FOR NON-LOCAL ELLIPTIC EQUATIONS 

HUXIAO LUO<br>Communicated by Raffaella Servadei


#### Abstract

This article concerns the existence of sign-changing solutions for equations driven by a non-local integrodifferential operator with homogeneous Dirichlet boundary conditions, $$
\begin{gathered} -\mathcal{L}_{K} u=f(x, u), \quad x \in \Omega \\ u=0, \quad x \in \mathbb{R}^{n} \backslash \Omega \end{gathered}
$$ where $\Omega \subset \mathbb{R}^{n}(n \geq 2)$ is a bounded, smooth domain and the nonlinear term $f$ satisfies suitable growth assumptions. By using Brouwer's degree theory and Deformation Lemma and arguing as in 2 , we prove that there exists a least energy sign-changing solution. Our results generalize and improve some results obtained in 27 .


## 1. Introduction

In recent years, fractional and non-local operators of elliptic type have been widely investigated. This type of operators arises in several areas such as anomalous diffusion, the thin obstacle problem, optimization, finance, phase transitions, stratified materials, crystal dislocation, soft thin films, semipermeable membranes, flame propagation, conservation laws, quasi-geostrophic flows, multiple scattering and materials science. One can see [8, 9, 12, 14, 20, 21, 22, 23, 24, 25, 26, 28, and their references. Many publications [3, 4, 5, [6, (7, 10, 11, 13, [15, 17, 18, 19] are devoted to the study of the existence of sign-changing solutions of classical elliptic boundary value problems such as

$$
\begin{align*}
-\Delta u & =f(x, u), \quad x \in \Omega, \\
u & =0, \quad x \in \partial \Omega, \tag{1.1}
\end{align*}
$$

where $f \in C(\bar{\Omega} \times \mathbb{R}, \mathbb{R}), \Omega \subset \mathbb{R}^{n}(n \geq 2)$ is a bounded domain with smooth boundary $\partial \Omega$. There have been several methods developed in studying sign-changing solutions of nonlinear elliptic equations, such as the invariant sets of descending flow method developed by Liu and Sun [4, 15, 19, and the minimax method which is established by Berestycki and Lions in the classical paper [7]. Teng, Wang and Wang [27] established the existence of a least energy sign-changing solution for nonlinear

[^0]problems involving the fractional Laplacian by using a constrained minimization method.

Ambrosio and Isernia [2] studied the fractional Schrödinger equation

$$
\begin{equation*}
(-\Delta)^{s} u+V(x) u=K(x) f(u) \quad \text { in } \mathbb{R}^{n} \tag{1.2}
\end{equation*}
$$

where $0<s<1, n>2 s,(-\Delta)^{s}$ is the fractional Laplacian operator, which (up to normalization factors) may be defined as

$$
-(-\Delta)^{s} u(x)=\frac{1}{2} \int_{\mathbb{R}^{n}} \frac{u(x+y)+u(x-y)-2 u(x)}{|y|^{n+2 s}} d y
$$

By using a minimization argument and a quantitative deformation Lemma, Ambrosio and Isernia proved the existence of sign-changing solutions for (1.2).

Motivated by the above works, in this article, we study the non-local elliptic problems

$$
\begin{gather*}
-\mathcal{L}_{K} u=f(x, u), \quad x \in \Omega \\
u=0, \quad x \in \mathbb{R}^{n} \backslash \Omega \tag{1.3}
\end{gather*}
$$

where $f$ satisfies the following assumptions
(A1) $f \in C^{1}\left(\bar{\Omega} \times \mathbb{R}^{n}\right), \lim _{\tau \rightarrow 0} \frac{f(x, \tau)}{\tau}=0$, uniformly in $x \in \bar{\Omega}$;
(A2) $|f(x, \tau)| \leq C\left(1+|\tau|^{p-1}\right)$ for some $C>0$ and $p \in\left(2,2_{s}^{*}\right)$, where $2_{s}^{*}=\frac{2 n}{n-2 s}$;
(A3) There exists a constant $\mu>2$ such that

$$
0<\mu F(x, \tau) \leq \tau f(x, \tau), \quad \forall x \in \bar{\Omega}, \tau \in \mathbb{R} \backslash\{0\}
$$

where $F(x, \tau)=\int_{0}^{\tau} f(x, t) d t$;
(A4) For every $x \in \bar{\Omega}$ the function $\tau \mapsto \frac{f(x, \tau)}{|\tau|}$ is strictly increasing for all $|\tau|>0$. The non-local integrodifferential operator $\mathcal{L}_{K}$ is defined as follows

$$
\mathcal{L}_{K} u(x)=\frac{1}{2} \int_{\mathbb{R}^{n}}(u(x+y)+u(x-y)-2 u(x)) K(y) d y, \quad x \in \mathbb{R}^{n}
$$

where $K: \mathbb{R}^{n} \backslash\{0\} \rightarrow(0,+\infty)$ is a function with the properties
(A5) $\gamma K \in L^{1}\left(\mathbb{R}^{n}\right)$, where $\gamma(x)=\min \left\{|x|^{2}, 1\right\}$;
(A6) there exists $\lambda>0$ such that $K(x) \geq \lambda|x|^{-(n+2 s)}$ for any $x \in \mathbb{R}^{n} \backslash\{0\}$.
A typical example for $K$ is given by $K(x)=|x|^{-(n+2 s)}$. In this case $\mathcal{L}_{K}$ is the fractional Laplace operator $-(-\Delta)^{s}$.

To prove the existence of sign-changing solutions for (1.3), we argue exactly as in [2], where the authors deal with fractional Schrödinger equations.

We remark that the Dirichlet datum is given in $\mathbb{R}^{n} \backslash \Omega$ and not simply on $\partial \Omega$, consistent with the non-local character of the operator $\mathcal{L}_{K}$.

As is explained in [7, Remark 9.2], the minimax method of Berestycki and Lions strongly depends on a kind of nodal structure associated with equation 1.1), which is unknown for problem (1.3). The variational methods used in [4, 17] heavily rely on the following decomposition, for $u \in H_{0}^{1}(\Omega)$,

$$
\begin{gather*}
\left\langle\Phi^{\prime}(u), u^{+}\right\rangle=\left\langle\Phi^{\prime}\left(u^{+}\right), u^{+}\right\rangle, \quad\left\langle\Phi^{\prime}(u), u^{-}\right\rangle=\left\langle\Phi^{\prime}\left(u^{-}\right), u^{-}\right\rangle,  \tag{1.4}\\
\Phi(u)=\Phi\left(u^{+}\right)+\Phi\left(u^{-}\right),
\end{gather*}
$$

where $u^{+}:=\max \{u, 0\}, u^{-}:=\min \{u, 0\}$ and $\Phi$ is the energy functional of 1.1 given by

$$
\Phi(u)=\frac{1}{2} \int_{\Omega}|\nabla u|^{2} d x-\int_{\Omega} F(x, u) d x
$$

But for problem (1.3) and $u \in X_{0}$ (the space $X_{0}$ was introduced in [20, 21], and will also be defined in Section 2), we have

$$
\begin{gathered}
I(u)=I\left(u^{+}\right)+I\left(u^{-}\right)-\int_{\mathbb{R}^{2 n}}\left(u^{-}(x) u^{+}(y)+u^{-}(y) u^{+}(x)\right) K(x-y) d x d y \\
\left\langle I^{\prime}(u), u^{+}\right\rangle=\left\langle I^{\prime}\left(u^{+}\right), u^{+}\right\rangle-\int_{\mathbb{R}^{2 n}}\left(u^{-}(x) u^{+}(y)+u^{-}(y) u^{+}(x)\right) K(x-y) d x d y
\end{gathered}
$$

where $I$ is the energy functional of 1.3 given by

$$
I(u)=\frac{1}{2} \int_{\mathbb{R}^{2 n}}|u(x)-u(y)|^{2} K(x-y) d x d y-\int_{\Omega} F(x, u) d x
$$

Clearly, the functional $I$ does no longer satisfy the decomposition 1.4. In this article, motivated by [4, 5, 17, we try to get sign-changing solutions for (1.3) by seeking the minimizer of the energy functional $I$ over the following constraint:

$$
\mathcal{M}:=\left\{u \in \mathcal{N}: u^{ \pm} \not \equiv 0,\left\langle I^{\prime}(u), u^{+}\right\rangle=\left\langle I^{\prime}(u), u^{-}\right\rangle=0\right\}
$$

where the set $\mathcal{N}$ defined as

$$
\mathcal{N}:=\left\{u \in X_{0} \backslash\{0\}:\left\langle I^{\prime}(u), u\right\rangle=0\right\} .
$$

Since $\mathcal{L}_{K}$ is nonlocal, we need some technical analysis to show that $\mathcal{M} \neq \emptyset$. Since we look for sign-changing solution to (1.3), it is natural to seek functions $w \in \mathcal{M}$ such that

$$
I(w)=\inf _{v \in \mathcal{M}} I(v)
$$

As in [1, 3], we are able to prove the existence of a minimizer of $I$ on $\mathcal{M}$ and that it is a weak solution to 1.3 by using a suitable deformation argument. Now, we are ready to state the main results of this paper.
Theorem 1.1. Suppose that (A1)-(A4) hold. Then problem 1.3) admits a least energy sign-changing solution.

For equation (1.1), we can follow the argument of 4 to show that the least energy sign-changing solution has exactly two nodal domains. But in this framework, because there is a nonlocal term $\mathcal{L}_{K} u$, we can not get the same result.

The rest of this article is organized as follows. In Section 2, we prove some lemmas, which are crucial to investigate our main result. The proof of Theorem 1.1 is given in Section 3.

## 2. Preliminaries

Recall that the space $X$ introduced by Servadei and Valdinoci [20, 21, 22, 23, [24, 25] is defined as a linear space of Lebesgue measurable functions from $\mathbb{R}^{n}$ to $\mathbb{R}$, such that, any function $u$ restricted in $X$ belongs to $L^{2}(\Omega)$ and the map $(x, y) \mapsto$ $(u(x)-u(y)) \sqrt{K(x-y)}$ is in $L^{2}\left(\mathbb{R}^{2 n} \backslash(\mathcal{C} \Omega \times \mathcal{C} \Omega), d x d y\right)$, where $\mathcal{C} \Omega:=\mathbb{R}^{n} \backslash \Omega$. Moreover,

$$
X_{0}=\left\{u \in X: u=0 \text { a.e. in } \mathbb{R}^{n} \backslash \Omega\right\} .
$$

The function space $X$ is equipped with the norm

$$
\begin{equation*}
\|u\|_{X}=\left(\|u\|_{L^{2}(\Omega)}^{2}+\int_{Q}|u(x)-u(y)|^{2} K(x-y) d x d y\right)^{1 / 2} \tag{2.1}
\end{equation*}
$$

where $Q=\mathbb{R}^{2 n} \backslash(\mathcal{C} \Omega \times \mathcal{C} \Omega)$. For any $u \in X_{0}$, the space $X_{0}$ is endowed with the norm

$$
\|u\|_{X_{0}}=\left(\int_{Q}|u(x)-u(y)|^{2} K(x-y) d x d y\right)^{1 / 2}
$$

which is equivalent to the usual one defined in 2.1. In the following, we also denote the norm $\|\cdot\|_{X_{0}}$ as $\|\cdot\|$.

For the reader's convenience, we review the main embedding results for the space $X_{0}$.

Lemma 2.1 ([20, 21, 22, 23, 24, 25]). The embedding $X_{0} \hookrightarrow L^{r}(\Omega)$ is continuous for any $r \in\left[2,2_{s}^{*}\right]$, and compact for any $r \in\left[2,2_{s}^{*}\right)$.

Now, we collect some preliminary lemmas which will be used in the last section to prove our main result.
Lemma 2.2. Let $\left\{u_{j}\right\}$ be a sequence such that $u_{j} \rightharpoonup u$ in $X_{0}$, then, up to $a$ subsequence,

$$
\begin{equation*}
\lim _{j \rightarrow \infty} \int_{\Omega} f\left(x, u_{j}\right) u_{j} d x=\int_{\Omega} f(x, u) u d x \tag{i}
\end{equation*}
$$

$$
\begin{equation*}
\lim _{j \rightarrow \infty} \int_{\Omega} F\left(x, u_{j}\right) d x=\int_{\Omega} F(x, u) d x \tag{ii}
\end{equation*}
$$

$$
\begin{align*}
& \liminf _{j \rightarrow \infty} \int_{\mathbb{R}^{2 n}}\left(-u_{j}^{-}(x) u_{j}^{+}(y)-u_{j}^{-}(y) u_{j}^{+}(x)\right) K(x-y) d x d y  \tag{iii}\\
& \geq \int_{\mathbb{R}^{2 n}}\left(-u^{-}(x) u^{+}(y)-u^{-}(y) u^{+}(x)\right) K(x-y) d x d y \tag{2.2}
\end{align*}
$$

Proof. (i) By the compact embedding $X_{0} \hookrightarrow L^{p}(\Omega)\left(2 \leq p<2_{s}^{*}\right)$, taking if necessary a subsequence, we have $u_{j} \rightarrow u$ in $L^{p}(\Omega)$ and $u_{j}(x) \rightarrow u(x)$ a.e. on $\mathbb{R}^{n}$. By a standard discussion, there exists a function $g \in L^{p}(\Omega)$ such that

$$
|u(x)|,\left|u_{j}(x)\right| \leq g(x)
$$

By (A2) and $u \in L^{p}(\Omega)$, we have

$$
|f(x, u)|^{\frac{p}{p-1}} \leq C^{\frac{p}{p-1}}\left(1+|u|^{p-1}\right)^{\frac{p}{p-1}} \leq C^{\frac{p}{p-1}} 2^{\frac{p}{p-1}}\left(1+|u|^{p}\right) \in L^{1}(\Omega)
$$

it follows that $f(\cdot, u) \in L^{\frac{p}{p-1}}(\Omega)$. Since

$$
\left|f\left(x, u_{j}\right)-f(x, u)\right|^{\frac{p}{p-1}} \leq 2^{\frac{p}{p-1}} C^{\frac{p}{p-1}}\left(1+|g|^{p-1}\right)^{\frac{p}{p-1}} \in L^{1}(\Omega),
$$

it follows from the Lebesgue dominated convergence theorem that

$$
\int_{\Omega}\left|f\left(x, u_{j}\right)-f(x, u)\right|^{\frac{p}{p-1}} d x \rightarrow 0, \quad \text { as } j \rightarrow \infty
$$

By the Hölder inequality, we have

$$
\begin{aligned}
\int_{\Omega}\left(f\left(x, u_{j}\right)-f(x, u)\right) u_{j} d x & \leq\left(\int_{\Omega} \left\lvert\, f\left(x, u_{j}\right)-f(x, u)^{\frac{p}{p-1}} d x\right.\right)^{\frac{p-1}{p}}\left(\int_{\Omega}\left|u_{j}\right|^{p} d x\right)^{1 / p} \\
& \rightarrow 0, \quad \text { as } j \rightarrow \infty
\end{aligned}
$$

Thus

$$
\lim _{j \rightarrow \infty} \int_{\Omega} f\left(x, u_{j}\right) u_{j} d x=\lim _{n \rightarrow \infty} \int_{\Omega} f(x, u) u_{j} d x=\int_{\Omega} f(x, u) u d x
$$

(ii) By the mean value theorem, there exists $\lambda \in[0,1]$ such that

$$
\begin{aligned}
& \left|\int_{\Omega}\left(F\left(x, u_{j}\right)-F(x, u)\right) d x\right| \\
& =\left|\int_{\Omega} f\left(x, u+\lambda\left(u_{j}-u\right)\right)\left(u_{j}-u\right) d x\right| \\
& \leq \int_{\Omega} C\left(1+\left|u+\lambda\left(u_{j}-u\right)\right|^{p-1}\right)\left|u_{j}-u\right| d x \\
& \leq C \int_{\Omega}\left|u_{j}-u\right| d x+2^{p} C \int_{\Omega}|u|^{p-1}\left|u_{j}-u\right| d x+2^{p} C \int_{\Omega}\left|u_{j}-u\right|^{p} d x \\
& \leq C\left\|u_{j}-u\right\|_{1}+2^{p} C\|u\|_{p}^{p-1}\left\|u_{i}-u\right\|_{p}+2^{p} C\left\|u_{j}-u\right\|_{p} .
\end{aligned}
$$

Thus

$$
\lim _{j \rightarrow \infty} \int_{\Omega}\left(F\left(x, u_{j}\right)-F(x, u)\right) d x=0
$$

(iii) By using $u_{j}(x) \rightarrow u(x)$ a.e. on $\mathbb{R}^{n}$, Fatou's Lemma and

$$
\left(-u_{j}^{-}(x) u_{j}^{+}(y)-u_{j}^{-}(y) u_{j}^{+}(x)\right) K(x-y) \geq 0
$$

we have that 2.2 holds.
Lemma 2.3. (i) For all $u \in \mathcal{N}$ such that $\|u\| \rightarrow+\infty$, we have $I(u) \rightarrow+\infty$;
(ii) There exists $\rho>0$ such that $\|u\| \geq \rho$ for all $u \in \mathcal{N}$ and $\left\|w^{ \pm}\right\| \geq \rho$ for all $w \in \mathcal{M}$.

Proof. (i) By the definition of $\mathcal{N}$ and assumption (A3), we have

$$
\begin{aligned}
I(u) & =I(u)-\frac{1}{\mu}\left\langle I^{\prime}(u), u\right\rangle \\
& =\left(\frac{1}{2}-\frac{1}{\mu}\right)\|u\|^{2}-\int_{\Omega}\left(F(x, u)-\frac{1}{\mu} f(x, u) u\right) d x \\
& \geq\left(\frac{1}{2}-\frac{1}{\mu}\right)\|u\|^{2}
\end{aligned}
$$

Thus, $I(u) \rightarrow+\infty$ as $\|u\| \rightarrow+\infty$.
(ii) By assumptions (A1) and (A2), we have that for any $\varepsilon>0$, there exists a positive constant $C_{\varepsilon}$ such that

$$
\begin{equation*}
|f(t) t| \leq \varepsilon|t|^{2}+C_{\varepsilon}|t|^{p}, \quad \text { for all } t \in \mathbb{R} \tag{2.3}
\end{equation*}
$$

where $2<p<2_{s}^{*}$. Since $u \in \mathcal{N}$ we have $\left\langle I^{\prime}(u), u\right\rangle=0$, that is

$$
\|u\|=\int_{\Omega} f(x, u) u d x
$$

By (2.3), $X_{0} \hookrightarrow L^{2}(\Omega)$ and $X_{0} \hookrightarrow L^{p}(\Omega)$, we have

$$
\begin{equation*}
\|u\|^{2} \leq \varepsilon \int_{\Omega}|u|^{2} d x+C_{\varepsilon} \int_{\Omega}|u|^{p} d x \leq \varepsilon \gamma_{2}\|u\|^{2}+C_{\varepsilon} \gamma_{p}\|u\|^{p} \tag{2.4}
\end{equation*}
$$

where $\gamma_{2}$ and $\gamma_{p}$ are the embedding constants. We can choose $\varepsilon$ small enough in order to find $\rho>0$ such that $\|u\| \geq \rho$. Now, for $w \in \mathcal{M}$, we have that $\left\langle I^{\prime}(w), w^{ \pm}\right\rangle=$ 0 , so

$$
\left\langle I^{\prime}\left(w^{ \pm}\right), w^{ \pm}\right\rangle=\int_{\mathbb{R}^{2 n}}\left(u^{-}(x) u^{+}(y)+u^{-}(y) u^{+}(x)\right) K(x-y) d x d y \leq 0
$$

which gives

$$
\left\|w^{ \pm}\right\|^{2} \leq \int_{\Omega} w^{ \pm} f\left(x, w^{ \pm}\right) d x
$$

Then we can proceed as in the proof of (i).
Lemma 2.4. Let $\left\{w_{j}\right\} \subset \mathcal{M}$ such that $w_{j} \rightharpoonup w$ in $X_{0}$. Then $w^{ \pm} \neq 0$.
Proof. Firstly we observe that by Lemma 2.3 there exists $\rho>0$ such that

$$
\begin{equation*}
\left\|w_{j}^{ \pm}\right\| \geq \rho \quad \text { for all } j \in \mathbb{N} \tag{2.5}
\end{equation*}
$$

Since $w_{j} \in \mathcal{M}$, we have $\left\langle I^{\prime}\left(w_{j}\right), w_{j}^{ \pm}\right\rangle=0$; that is,

$$
\begin{align*}
& \left\|w_{j}^{ \pm}\right\|^{2}-\int_{\mathbb{R}^{2 n}}\left(w_{j}^{-}(x) w_{j}^{+}(y)+w_{j}^{-}(y) w_{j}^{+}(x)\right) K(x-y) d x d y  \tag{2.6}\\
& =\int_{\Omega} f\left(x, w_{j}^{ \pm}\right) w_{j}^{ \pm} d x
\end{align*}
$$

At this point, recalling that

$$
-\int_{\mathbb{R}^{2 n}}\left(w_{j}^{-}(x) w_{j}^{+}(y)+w_{j}^{-}(y) w_{j}^{+}(x)\right) K(x-y) d x d y \geq 0
$$

by 2.5 and 2.6 we deduce that

$$
\begin{equation*}
\rho^{2} \leq\left\|w_{j}^{ \pm}\right\|^{2} \leq \int_{\Omega} f\left(x, w_{j}^{ \pm}\right) w_{j}^{ \pm} d x \tag{2.7}
\end{equation*}
$$

Now, by the fact that $w_{j} \rightharpoonup w$ in $X_{0}$ and the compactly embedding $X_{0} \hookrightarrow L^{q}(\Omega)(q \in$ $\left[2,2_{s}^{*}\right)$ ), we know that $w_{j} \rightarrow w$ in $L^{q}(\Omega)$. Moreover, by using that $\left|t^{ \pm}-s^{ \pm}\right| \leq|t-s|$ for all $t, s \in \mathbb{R}$, we can deduce that $w_{j}^{ \pm} \rightarrow w^{ \pm}$in $L^{q}(\Omega)$, and for all $x \in \Omega$, we also have $w_{j}^{ \pm} \rightarrow w^{ \pm}$a.e. in $\Omega$. Similarly to Lemma 2.2 it is easy to see that

$$
\begin{equation*}
\int_{\Omega} f\left(x, w_{j}^{ \pm}\right) w_{j}^{ \pm} d x \rightarrow \int_{\Omega} f\left(x, w^{ \pm}\right) w^{ \pm} d x \tag{2.8}
\end{equation*}
$$

Putting together 2.7 and 2.8 we have

$$
0<\rho^{2} \leq \int_{\Omega} f\left(x, w^{ \pm}\right) w^{ \pm} d x
$$

showing that $w^{ \pm} \neq 0$.
Lemma 2.5. If $v \in X_{0}: v^{ \pm} \neq 0$, then there exist $s, t>0$ such that

$$
\left\langle I^{\prime}\left(t v^{+}+s v^{-}\right), v^{+}\right\rangle=0 \quad \text { and } \quad\left\langle I^{\prime}\left(t v^{+}+s v^{-}\right), v^{-}\right\rangle=0
$$

As a consequence $t v^{+}+s v^{-} \in \mathcal{M}$.
Proof. Let $W:(0,+\infty) \times(0,+\infty) \rightarrow \mathbb{R}^{2}$ be a continuous vector field given by

$$
W(t, s)=\left(\left\langle I^{\prime}\left(t v^{+}+s v^{-}\right), t v^{+}\right\rangle,\left\langle I^{\prime}\left(t v^{+}+s v^{-}\right), s v^{-}\right\rangle\right)
$$

for every $t, s \in(0,+\infty) \times(0,+\infty)$. By using 2.3), $X_{0} \hookrightarrow L^{2}(\Omega)$ and $X_{0} \hookrightarrow L^{p}(\Omega)$, we have

$$
\begin{align*}
& \left\langle I^{\prime}\left(t v^{+}+s v^{-}\right), t v^{+}\right\rangle \\
& =t^{2}\left\|v^{+}\right\|^{2}-s t \int_{\mathbb{R}^{2 n}}\left(v^{-}(x) v^{+}(y)+v^{-}(y) v^{+}(x)\right) K(x-y) d x d y \\
& \quad-\int_{\Omega} t v^{+} f\left(x, t v^{+}\right) d x  \tag{2.9}\\
& \geq t^{2}\left\|v^{+}\right\|^{2}-\int_{\Omega} t v^{+} f\left(x, t v^{+}\right) d x \geq t^{2}\left\|v^{+}\right\|^{2}-\varepsilon t^{2}\left\|v^{+}\right\|_{2}^{2}-C_{\varepsilon} t^{p}\left\|v^{+}\right\|_{p}^{p} \\
& \geq\left(1-\varepsilon \gamma_{2}\right) t^{2}\left\|v^{+}\right\|^{2}-C_{\varepsilon} \gamma_{p} t^{p}\left\|v^{+}\right\|_{p}^{p}
\end{align*}
$$

We choose $\varepsilon=\frac{1}{2 \gamma_{2}}$, then there exists $r>0$ small enough such that $\left\langle I^{\prime}\left(r v^{+}+\right.\right.$ $\left.\left.s v^{-}\right), r v^{+}\right\rangle>0$ for all $s>0$, and similarly there exists $\bar{r}>0$ small enough such that $\left\langle I^{\prime}\left(t v^{+}+\bar{r} v^{-}\right), \bar{r} v^{-}\right\rangle>0$ for all $t>0$. By assumption (A3) there exists a positive constant $C_{1}$ such that

$$
\begin{equation*}
F(x, t) \geq C_{1} t^{\mu}, \text { for every } t \text { sufficiently large, uniformly in } x \in \bar{\Omega} \tag{2.10}
\end{equation*}
$$

Hence, taking into account that $\mu>2$, we obtain

$$
\begin{aligned}
&\left\langle I^{\prime}\left(t v^{+}+s v^{-}\right), t v^{+}\right\rangle \\
&= t^{2}\left\|v^{+}\right\|^{2}-s t \int_{\mathbb{R}^{2 n}}\left(v^{-}(x) v^{+}(y)+v^{-}(y) v^{+}(x)\right) K(x-y) d x d y \\
&-\int_{\Omega} t v^{+} f\left(x, t v^{+}\right) d x \\
& \leq t^{2}\left\|v^{+}\right\|^{2}-s t \int_{\mathbb{R}^{2 n}}\left(v^{-}(x) v^{+}(y)+v^{-}(y) v^{+}(x)\right) K(x-y) d x d y \\
&-\mu \int_{\Omega} F\left(x, t v^{+}\right) d x \\
& \leq t^{2}\left\|v^{+}\right\|^{2}-s t \int_{\mathbb{R}^{2 n}}\left(v^{-}(x) v^{+}(y)+v^{-}(y) v^{+}(x)\right) K(x-y) d x d y \\
&-t^{\mu} C_{1} \mu \int_{\Omega}\left|v^{+}\right|^{\mu} d x \rightarrow-\infty \quad \text { as } t \rightarrow+\infty
\end{aligned}
$$

Then, there exists $R>0$ sufficiently large such that $\left\langle I^{\prime}\left(R v^{+}+s v^{-}\right), R v^{+}\right\rangle<0$ for all $s>0$ and similarly we can find $\bar{R}>0$ such that $\left\langle I^{\prime}\left(t v^{+}+\bar{R} v^{-}\right), \bar{R} v^{-}\right\rangle<0$ for all $t>0$. As a consequence, we have proved the existence of suitable $0<r<R$ such that, for all $t, s \in[r, R]$ it holds

$$
\begin{align*}
&\left\langle I^{\prime}\left(r v^{+}+s v^{-}\right), r v^{+}\right\rangle>0, \\
&\left\langle I^{\prime}\left(R v^{+}+s v^{\prime}\left(t v^{+}+\bar{r} v^{-}\right), \bar{r} v^{-}\right\rangle>0\right.  \tag{2.11}\\
& \hline
\end{align*}
$$

By applying Miranda's theorem [16] we have the conclution.
Definition 2.6. For each $v \in X_{0}$ with $v^{ \pm} \neq 0$, let us consider the function $h^{v}$ : $[0,+\infty) \times[0,+\infty) \rightarrow \mathbb{R}$ given by

$$
h^{v}(t, s)=I\left(t v^{+}+s v^{-}\right)
$$

and its gradient $\Phi^{v}:[0,+\infty) \times[0,+\infty) \rightarrow \mathbb{R}^{2}$ defined by

$$
\begin{align*}
\Phi^{v}(t, s) & =\left(\Phi_{1}^{v}(t, s), \Phi_{2}^{v}(t, s)\right)=\left(\frac{\partial h^{v}}{\partial t}(t, s), \frac{\partial h^{v}}{\partial s}(t, s)\right)  \tag{2.12}\\
& =\left(\left\langle I^{\prime}\left(t v^{+}+s v^{-}\right), v^{+}\right\rangle,\left\langle I^{\prime}\left(t v^{+}+s v^{-}\right), v^{-}\right\rangle\right)
\end{align*}
$$

Furthermore, we consider the Jacobian matrix of $\Phi^{v}$ :

$$
\left(\Phi^{v}\right)^{\prime}(t, s)=\left(\begin{array}{cc}
\frac{\partial \Phi_{1}^{v}}{\partial t}(t, s) & \frac{\partial \Phi_{1}^{v}}{\partial s}(t, s) \\
\frac{\partial \Phi_{2}^{v}}{\partial t}(t, s) & \frac{\partial \Phi_{1} 2^{v}}{\partial s}(t, s)
\end{array}\right)
$$

In the following we aim to prove that, if $w \in \mathcal{M}$, the function $h^{w}$ has a critical point and in particular a global minimum in $(t, s)=(1,1)$.

Lemma 2.7. If $w \in \mathcal{M}$, then
(i) $h^{w}(t, s)<h^{w}(1,1)=I(w)$, for all $t, s \geq 0$ such that $(t, s) \neq(1,1)$;
(ii) $\operatorname{det}\left(\Phi^{w}\right)^{\prime}(1,1)>0$.

Proof. (i) Since $w \in \mathcal{M}$, then $\left\langle I^{\prime}(w), w^{ \pm}\right\rangle=0$, that is

$$
\begin{aligned}
& \left\|w^{+}\right\|^{2}-\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y=\int_{\Omega} f\left(x, w^{+}\right) w^{+} d x \\
& \left\|w^{-}\right\|^{2}-\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y=\int_{\Omega} f\left(x, w^{-}\right) w^{-} d x
\end{aligned}
$$

From this and by the definition of $\Phi^{w}$, it follows that $(1,1)$ is a critical point of $h^{w}$. By (A3), there exists constants $C_{1}>0$ and $C_{2}>0$ such that

$$
F(x, \tau) \geq C_{1}|\tau|^{\mu}-C_{2}, \quad \forall x \in \bar{\Omega}
$$

It follows that

$$
\begin{aligned}
h^{w}(t, s)= & I\left(t w^{+}+s w^{-}\right) \\
\leq & \frac{1}{2}\left\|t w^{+}+s w^{-}\right\|^{2}-C_{1} \int_{\Omega}\left|t w^{+}+s w^{-}\right|^{\mu} d x+|\Omega| C_{2} \\
\leq & \frac{t^{2}+s^{2}}{2}\left[\left\|w^{+}\right\|^{2}+\left\|w^{-}\right\|^{2}\right. \\
& \left.-\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y\right] \\
& -C_{1} t^{\mu} \int_{\Omega}\left|w^{+}\right|^{\mu} d x-C_{1} s^{\mu} \int_{\Omega}\left|w^{-}\right|^{\mu} d x+|\Omega| C_{2} \\
\leq & \frac{t^{2}+s^{2}}{2}\left[\left\|w^{+}\right\|^{2}+\left\|w^{-}\right\|^{2}\right. \\
& \left.-\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y\right] \\
& -C_{1} t^{\mu} \int_{\Omega}\left|w^{+}\right|^{\mu} d x+|\Omega| C_{2} .
\end{aligned}
$$

Let us suppose that $|t| \geq|s|>0$, thus, by using $t^{2}+s^{2} \leq 2 t^{2}$ we see that

$$
\begin{aligned}
\frac{h^{w}(t, s)}{t^{2}+s^{2}} \leq & \frac{1}{2}\left[\left\|w^{+}\right\|^{2}+\left\|w^{-}\right\|^{2}-\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y\right] \\
& -C_{1} \frac{t^{\mu-2}}{2} \int_{\Omega}\left|w^{+}\right|^{\mu} d x+\frac{|\Omega| C_{2}}{t^{2}+s^{2}}
\end{aligned}
$$

Taking into account that $\mu>2$, we infer that

$$
\lim _{|(t, s)| \rightarrow+\infty} h^{w}(t, s)=-\infty
$$

By using the continuity of $h^{w}$ we deduce the existence of $(\bar{t}, \bar{s}) \in[0,+\infty) \times[0,+\infty)$ that is a global maximum point of $h^{w}$.

Now we prove that $\bar{t}, \bar{s}>0$. Suppose by contradiction that $\bar{s}=0$. Then $\left\langle I^{\prime}\left(\bar{t} w^{+}\right), \bar{t} w^{+}\right\rangle=0$, that is

$$
\begin{equation*}
\left\|w^{+}\right\|^{2}=\int_{\Omega}\left(w^{+}\right)^{2} \frac{f\left(x, \bar{t} w^{+}\right)}{\bar{t} w^{+}} d x \tag{2.13}
\end{equation*}
$$

Since

$$
\begin{aligned}
\left\langle I^{\prime}\left(w^{+}\right), w^{+}\right\rangle & =\left\langle I^{\prime}(w), w^{+}\right\rangle+\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y \\
& =\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y<0
\end{aligned}
$$

we obtain

$$
\begin{equation*}
\left\|w^{+}\right\|^{2}<\int_{\Omega}\left(w^{+}\right)^{2} \frac{f\left(x, w^{+}\right)}{w^{+}} d x \tag{2.14}
\end{equation*}
$$

Then, combining 2.13 and 2.14 we obtain

$$
0<\int_{\Omega}\left(w^{+}\right)^{2}\left[\frac{f\left(x, w^{+}\right)}{w^{+}}-\frac{f\left(x, \bar{t} w^{+}\right)}{\bar{t} w^{+}}\right] d x .
$$

So, by (A4), we deduce that $\bar{t}<1$. It follows from the condition (A4) that for every $x \in \bar{\Omega}$

$$
\begin{align*}
\tau \mapsto \frac{1}{2} \tau f(x, \tau)-F(x, \tau) \text { is strictly increasing for all } \tau>0 \\
\tau \mapsto \frac{1}{2} \tau f(x, \tau)-F(x, \tau) \text { is strictly decreasing for all } \tau<0 \tag{2.15}
\end{align*}
$$

Thus

$$
\begin{aligned}
& h^{w}(\bar{t}, 0)=I\left(\bar{t} w^{+}\right) \\
& =I\left(\bar{t} w^{+}\right)-\frac{1}{2}\left\langle I^{\prime}\left(\bar{t} w^{+}\right), \bar{t} w^{+}\right\rangle \\
& =\int_{\Omega}\left[\frac{1}{2} \bar{t} w^{+} f\left(x, \bar{t} w^{+}\right)-F\left(x, \bar{t} w^{+}\right)\right] d x \\
& \leq \int_{\Omega}\left[\frac{1}{2} \bar{t} w^{+} f\left(x, \bar{t} w^{+}\right)-F\left(x, \bar{t} w^{+}\right)\right] d x+\int_{\Omega}\left[\frac{1}{2} \bar{t} w^{-} f\left(x, \bar{t} w^{-}\right)-F\left(x, \bar{t} w^{-}\right)\right] d x \\
& <\int_{\Omega}\left[\frac{1}{2} w^{+} f\left(x, w^{+}\right)-F\left(x, w^{+}\right)\right] d x+\int_{\Omega}\left[\frac{1}{2} w^{-} f\left(x, w^{-}\right)-F\left(x, w^{-}\right)\right] d x \\
& =I(w)-\frac{1}{2}\left\langle I^{\prime}(w), w\right\rangle \\
& =I(w)=h^{w}(1,1) .
\end{aligned}
$$

So $h^{w}(\bar{t}, 0)<h^{w}(1,1)$, and this gives a contradiction because $(\bar{t}, 0)$ is a global maximum point. Similarly we can prove that $\bar{t}>0$.

Now we show that $\bar{s}, \bar{t} \leq 1$. Since $\left(h^{w}\right)^{\prime}(\bar{t}, \bar{s})=0$, we obtain

$$
\bar{t}^{2}\left\|w^{+}\right\|^{2}-\bar{s} \bar{t} \int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y
$$

$$
\begin{aligned}
& =\int_{\Omega} \bar{t} w^{+} f\left(x, \bar{t} w^{+}\right) d x \\
& \bar{s}^{2}\left\|w^{-}\right\|^{2}-\bar{s} \bar{t} \int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y \\
& =\int_{\Omega} \bar{s} w^{-} f\left(x, \bar{s} w^{-}\right) d x
\end{aligned}
$$

Assume that $\bar{t} \geq \bar{s}$. Since

$$
\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y \leq 0
$$

we have

$$
\begin{align*}
& \bar{t}^{2}\left\|w^{+}\right\|^{2}-\bar{t}^{2} \int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y \\
& \geq \int_{\Omega} \bar{t} w^{+} f\left(x, \bar{t} w^{+}\right) d x \tag{2.16}
\end{align*}
$$

Since $\left\langle I^{\prime}(w), w^{+}\right\rangle=0(w \in \mathcal{M})$, we deduce that

$$
\left\|w^{+}\right\|^{2}-\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y=\int_{\Omega} w^{+} f\left(x, w^{+}\right) d x
$$

which together with 2.16 gives

$$
\int_{\Omega}\left[\frac{f\left(x, \bar{t} w^{+}\right)}{\bar{t} w^{+}}-\frac{f\left(x, w^{+}\right)}{w^{+}}\right] d x \leq 0
$$

By (A4) we can infer that $\bar{t} \leq 1$.
Now we aim to prove that $h^{w}$ does not assume a global maximum in $[0,1] \times$ $[0,1] \backslash\{(1,1)\}$, namely $h^{w}(\bar{t}, \bar{s})<h^{w}(1,1)$ for every $(\bar{t}, \bar{s}) \in[0,1] \times[0,1] \backslash\{(1,1)\}$. By the definition of $h^{w}$ and 2.15 we have

$$
\begin{aligned}
h^{w}(\bar{t}, \bar{s})= & I\left(\bar{t} w^{+}+\bar{s} w^{-}\right)-\frac{1}{2}\left\langle I^{\prime}\left(\bar{t} w^{+}+\bar{s} w^{-}\right), \bar{t} w^{+}\right\rangle-\frac{1}{2}\left\langle I^{\prime}\left(\bar{t} w^{+}+\bar{s} w^{-}\right), \bar{s} w^{-}\right\rangle \\
= & \int_{\Omega \cap\{w \geq 0\}}\left[\frac{1}{2} \bar{t} w^{+} f\left(x, \bar{t} w^{+}\right)-F\left(x, \bar{t} w^{+}\right)\right] d x \\
& +\int_{\Omega \cap\{w \leq 0\}}\left[\frac{1}{2} \bar{s} w^{-} f\left(x, \bar{s} w^{-}\right)-F\left(x, \bar{s} w^{-}\right)\right] d x \\
< & \int_{\Omega \cap\{w \geq 0\}}\left[\frac{1}{2} w^{+} f\left(x, w^{+}\right)-F\left(x, w^{+}\right)\right] d x \\
& +\int_{\Omega \cap\{w \leq 0\}}\left[\frac{1}{2} w^{-} f\left(x, w^{-}\right)-F\left(x, w^{-}\right)\right] d x \\
= & \int_{\Omega}\left[\frac{1}{2} w f(x, w)-F(x, w)\right] d x=h^{w}(1,1) .
\end{aligned}
$$

(ii) Firstly, let us observe that

$$
\begin{gather*}
\frac{\partial \Phi_{1}^{w}}{\partial t}(t, s)=\left\|w^{+}\right\|^{2}-\int_{\Omega} f^{\prime}\left(x, t w^{+}\right)\left(w^{+}\right)^{2} d x \\
\frac{\partial \Phi_{2}^{w}}{\partial s}(t, s)=\left\|w^{-}\right\|^{2}-\int_{\Omega} f^{\prime}\left(x, t w^{-}\right)\left(w^{-}\right)^{2} d x  \tag{2.17}\\
\frac{\partial \Phi_{1}^{w}}{\partial s}(t, s)=\frac{\partial \Phi_{2}^{w}}{\partial t}(t, s) \\
=-\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y
\end{gather*}
$$

By (A4), it is easy to see that for every $x \in \bar{\Omega}$,

$$
\begin{equation*}
\tau^{2} f^{\prime}(x, \tau)-\tau f(x, \tau)>0 \quad \text { for all } \tau \neq 0 \tag{2.18}
\end{equation*}
$$

Then, by using the fact that $w \in \mathcal{M}, 2.17$ and 2.18 , we have

$$
\begin{aligned}
& \operatorname{det}\left(\Phi^{w}\right)^{\prime}(1,1) \\
&= {\left[\left\|w^{+}\right\|^{2}-\int_{\Omega} f^{\prime}\left(x, w^{+}\right)\left(w^{+}\right)^{2} d x\right]\left[\left\|w^{-}\right\|^{2}-\int_{\Omega} f^{\prime}\left(x, w^{-}\right)\left(w^{-}\right)^{2} d x\right] } \\
&-\left[\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y\right]^{2} \\
&= {\left[\int_{\Omega}\left(\left(w^{+}\right)^{2} f^{\prime}\left(x, w^{+}\right)-w^{+} f\left(x, w^{+}\right)\right) d x\right.} \\
&\left.-\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y\right] \\
& \quad \times\left[\int_{\Omega}\left(\left(w^{-}\right)^{2} f^{\prime}\left(x, w^{-}\right)-w^{-} f\left(x, w^{-}\right)\right) d x\right. \\
&\left.-\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y\right] \\
& \quad-\left[\int_{\mathbb{R}^{2 n}}\left(w^{-}(x) w^{+}(y)+w^{-}(y) w^{+}(x)\right) K(x-y) d x d y\right]^{2}>0
\end{aligned}
$$

## 3. Proof of Theorem 1.1

In this section, we prove Theorem 1.1 by following two steps.
Step 1. We show that there exists $w \in \mathcal{M}$ such that $I(w)=\inf _{v \in \mathcal{M}} I(v)$. By Lemma 2.3, there exists a minimizing sequence $\left\{w_{j}\right\} \subset \mathcal{M}$, bounded in $X_{0}$, such that

$$
\begin{equation*}
I\left(w_{j}\right) \rightarrow \inf _{v \in \mathcal{M}} I(v)=: c_{0}>0 \tag{3.1}
\end{equation*}
$$

By Lemma 2.1, up to a subsequence, we have

$$
\begin{gathered}
w_{j}^{ \pm} \rightharpoonup w^{ \pm} \quad \text { in } X_{0} \\
w_{j}^{ \pm} \rightarrow w^{ \pm} \quad \text { in } L^{p}(\Omega) \\
w_{j}^{ \pm} \rightarrow w^{ \pm} \quad \text { a.e. in } \mathbb{R}^{n} .
\end{gathered}
$$

From Lemma 2.4 we deduce that $w^{ \pm} \neq 0$, so $w=w^{+}+w^{-}$is sign-changing. By Lemma 2.5 there exist $s, t>0$ such that

$$
\begin{equation*}
\left\langle I^{\prime}\left(t w^{+}+s w^{-}\right), w^{+}\right\rangle=0, \quad\left\langle I^{\prime}\left(t w^{+}+s w^{-}\right), w^{-}\right\rangle=0 \tag{3.2}
\end{equation*}
$$

and $t w^{+}+s w^{-} \in \mathcal{M}$. Now, we prove that $s, t \leq 1$. Since $w_{j} \in \mathcal{M}$, we have $\left\langle I^{\prime}\left(w_{j}\right), w_{j}^{ \pm}\right\rangle=0$ or equivalently

$$
\begin{align*}
& \left\|w_{j}^{+}\right\|^{2}-\int_{\mathbb{R}^{2 n}}\left(w_{j}^{-}(x) w_{j}^{+}(y)+w_{j}^{-}(y) w_{j}^{+}(x)\right) K(x-y) d x d y \\
& =\int_{\Omega} w_{j}^{+} f\left(x, w_{j}^{+}\right) d x \\
& \left\|w_{j}^{-}\right\|^{2}-\int_{\mathbb{R}^{2 n}}\left(w_{j}^{-}(x) w_{j}^{+}(y)+w_{j}^{-}(y) w_{j}^{+}(x)\right) K(x-y) d x d y  \tag{3.3}\\
& =\int_{\Omega} w_{j}^{-} f\left(x, w_{j}^{-}\right) d x
\end{align*}
$$

The weak lower semicontinuity of the norm $\|\cdot\|$ in $X_{0}$ yields

$$
\begin{equation*}
\left\|w^{ \pm}\right\|^{2} \leq \liminf _{n \rightarrow \infty}\left\|w_{j}^{ \pm}\right\|^{2} \tag{3.4}
\end{equation*}
$$

By using (3.3), (3.4) and Lemma 2.2, we obtain

$$
\begin{equation*}
\left\langle I^{\prime}(w), w^{+}\right\rangle \leq 0 \quad \text { and } \quad\left\langle I^{\prime}(w), w^{-}\right\rangle \leq 0 . \tag{3.5}
\end{equation*}
$$

Then, we combine (3.2) with 3.5, and arguing similarly as in the proof of Lemma 2.7 (i) we obtain that $s, t \leq 1$.

Next, we show that $I\left(t w^{+}+s w^{-}\right)=c_{0}$ and $t=s=1$. By using $t w^{+}+s w^{-} \in$ $\mathcal{M}, w_{j} \in \mathcal{M}$, 2.15, 3.1), $s, t \in(0,1]$ and Lemma 2.2 we can see

$$
\begin{aligned}
c_{0} \leq & I\left(t w^{+}+s w^{-}\right)=I\left(t w^{+}+s w^{-}\right)-\frac{1}{2}\left\langle I^{\prime}\left(t w^{+}+s w^{-}\right), t w^{+}+s w^{-}\right\rangle \\
= & \int_{\Omega}\left[\frac{1}{2} f\left(x, t w^{+}+s w^{-}\right)\left(t w^{+}+s w^{-}\right)-F\left(x, t w^{+}+s w^{-}\right)\right] d x \\
= & \int_{\Omega \cap\{w \geq 0\}}\left[\frac{1}{2} t w^{+} f\left(x, t w^{+}\right)-F\left(x, t w^{+}\right)\right] d x \\
& +\int_{\Omega \cap\{w \leq 0\}}\left[\frac{1}{2} s w^{-} f\left(x, s w^{-}\right)-F\left(x, s w^{-}\right)\right] d x \\
\leq & \int_{\Omega}\left[\frac{1}{2} w^{+} f\left(x, w^{+}\right)-F\left(x, w^{+}\right)\right] d x \\
& +\int_{\Omega}\left[\frac{1}{2} w^{-} f\left(x, w^{-}\right)-F\left(x, w^{-}\right)\right] d x \\
= & \lim _{j \rightarrow \infty}\left[\int_{\Omega}\left(\frac{1}{2} w_{j}^{+} f\left(x, w_{j}^{+}\right)-F\left(x, w_{j}^{+}\right)\right) d x+\int_{\Omega}\left(\frac{1}{2} w_{j}^{-} f\left(x, w_{j}^{-}\right)-F\left(x, w_{j}^{-}\right)\right) d x\right] \\
= & \lim _{j \rightarrow \infty} \int_{\Omega}\left(\frac{1}{2} w_{j} f\left(x, w_{j}\right)-F\left(x, w_{j}\right)\right) d x=\lim _{j \rightarrow \infty}\left[I\left(w_{j}\right)-\frac{1}{2}\left\langle I^{\prime}\left(w_{j}\right), w_{j}\right\rangle\right] \\
= & \lim _{j \rightarrow \infty} I\left(w_{j}\right)=c_{0} .
\end{aligned}
$$

Thus, we have proved that there exist $t, s \in(0,1]$ such that $t w^{+}+s w^{-} \in \mathcal{M}$ and $I\left(t w^{+}+s w^{-}\right)=c_{0}$. Let us observe that by the above calculation we can infer that $t=s=1$, so $w=w^{+}+w^{-} \in \mathcal{M}$ and $I\left(w^{+}+w^{-}\right)=c_{0}$.

Step 2. Now, we prove that $I^{\prime}(w)=0$. We argue by contradiction. Then, we can find a positive constant $\beta>0$ and $v_{0} \in X_{0}$ with $\left\|v_{0}\right\|=1$, such that $\left\langle I^{\prime}(w), v_{0}\right\rangle=2 \beta>0$. By the continuity of $I^{\prime}$, we can choose a radius $R$ so that $\left\langle I^{\prime}(v), v_{0}\right\rangle=\beta>0$ for every $v \in B_{R}(w) \subset X_{0}$ with $v^{ \pm} \neq 0$.

Let $D:=(a, b) \times(a, b) \subset \mathbb{R}^{2}$ with $0<a<1<b$ such that
(i) $(1,1) \in D$ and $\Phi^{w}(t, s)=(0,0)$ in $\bar{D}$ if, and only if, $(t, s)=(1,1)$,
(ii) $c_{0} \notin h^{w}(\partial D)$,
(iii) $\left\{t w^{+}+s w^{-}:(t, s) \in \bar{D}\right\} \subset B_{R}(w)$,
where $h^{w}$ and $\Phi^{w}$ are defined by Definition 2.6, and satisfy Lemma 2.7. Then we can choose a radius $0<r<R$ such that

$$
\begin{equation*}
B=\bar{B}_{r}(w) \subset B_{R}(w) \quad \text { and } \quad B \cap\left\{t w^{+}+s w^{-}:(t, s) \in \partial D\right\}=\emptyset \tag{3.6}
\end{equation*}
$$

Now, let us define a continuous mapping $\rho: X_{0} \rightarrow[0,+\infty)$ such that $\rho(u):=$ $\operatorname{dist}\left(u, B^{c}\right)$ for all $u \in X_{0}$, and we consider a bounded Lipschitz vector field $\mathbb{V}$ : $X_{0} \rightarrow X_{0}$ given by $\mathbb{V}(u)=-\rho(u) v_{0}$. For every $u \in X_{0}$, denoting by $\eta(\tau)=\eta(\tau, u)$, we consider the following Cauchy problem

$$
\begin{gathered}
\eta^{\prime}(\tau)=\mathbb{V}(\eta(\tau)) \quad \text { for all } \tau>0 \\
\eta(0)=u
\end{gathered}
$$

We observe that there exist a continuous deformation $\eta(\tau, u)$ and $\tau_{0}>0$ such that for all $\tau \in\left[0, \tau_{0}\right]$ the following properties hold:
(a) $\eta(\tau, u)=u$ for all $u \notin B$,
(b) $\tau \rightarrow I(\eta(\tau, u))$ is decreasing for all $\eta(\tau, u) \in B$,
(c) $I(\eta(\tau, w)) \leq I(w)-\frac{r \beta}{2} \tau$.
(a) follows from the definition of $\rho$. Regarding (b), we observe that $\left\langle I^{\prime}(\eta(\tau)), v_{0}\right\rangle=$ $\beta>0$ for $\eta(\tau) \in B \subset B_{R}(w)$, and, by the definition of $\rho$, we have $\rho(\eta(\tau))>0$. Then

$$
\frac{d}{d \tau}(I(\eta(\tau)))=\left\langle I^{\prime}(\eta(\tau)), \eta^{\prime}(\tau)\right\rangle=-\rho(\eta(\tau))\left\langle I^{\prime}(\eta(\tau)), v_{0}\right\rangle=-\rho(\eta(\tau)) \beta<0
$$

for all $\eta(\tau) \in B$, that is $I(\eta(\tau, u))$ is decreasing with respect to $\tau$. Now we prove (c). Being $\tau_{0}>0$ such that $\eta(\tau, u) \in B$ for every $\tau \in\left[0, \tau_{0}\right]$, we can assume that

$$
\|\eta(\tau, w)-w\| \leq \frac{r}{2}, \quad \text { for any } \tau \in\left[0, \tau_{0}\right]
$$

Since $\rho(\eta(\tau, w))=\operatorname{dist}\left(\eta(\tau, w), B^{c}\right) \geq \frac{r}{2}$, we deduce that

$$
\frac{d}{d \tau} I(\eta(\tau, w)) \leq-\rho(\eta(\tau, w)) \beta \leq-\frac{r \beta}{2}
$$

and, integrating on $\left[0, \tau_{0}\right]$, we obtain

$$
I\left(\eta\left(\tau_{0}, w\right)\right)-I(w) \leq-\frac{r \beta}{2} \tau_{0}
$$

Now, we consider a suitable deformed path $\bar{\eta}_{0}: \bar{D} \rightarrow X_{0}$ defined by

$$
\bar{\eta}_{0}(t, s):=\eta\left(\tau_{0}, t w^{+}+s w^{-}\right), \quad \text { for all }(t, s) \in \bar{D}
$$

We note that

$$
\max _{(t, s) \in \bar{D}} I\left(\bar{\eta}_{0}(t, s)\right)<c_{0}
$$

Indeed, by (b) and the fact that $\eta(0, u)=u$, we have

$$
I\left(\bar{\eta}_{0}(t, s)\right)=I\left(\eta\left(\tau_{0}, t w^{+}+s w^{-}\right)\right) \leq I\left(\eta\left(0, t w^{+}+s w^{-}\right)\right)
$$

$$
\begin{aligned}
& =I\left(t w^{+}+s w^{-}\right)=h^{w}(t, s) \\
& <c_{0}, \quad \forall(t, s) \in \bar{D} \backslash\{(1,1)\},
\end{aligned}
$$

and for $(t, s)=(1,1)$, by (c) we have

$$
\begin{aligned}
I\left(\bar{\eta}_{0}(1,1)\right) & =I\left(\eta\left(\tau_{0}, w^{+}+w^{-}\right)\right)=I\left(\eta\left(\tau_{0}, w\right)\right) \\
& =\leq I(w)-\frac{r \beta}{2} \tau_{0}<I(w)=c_{0}
\end{aligned}
$$

Then, $\bar{\eta}_{0}(t, s) \cap \mathcal{M}=\emptyset$; that is,

$$
\begin{equation*}
\bar{\eta}_{0}(t, s) \notin \mathcal{M} \quad \text { for all }(t, s) \in \bar{D} \tag{3.7}
\end{equation*}
$$

On the other hand, defining $\Psi_{\tau_{0}}: D \rightarrow \mathbb{R}^{2}$ such that

$$
\Psi_{\tau_{0}}:=\left(\frac{1}{t}\left\langle I^{\prime}\left(\bar{\eta}_{0}(t, s)\right),\left(\bar{\eta}_{0}(t, s)\right)^{+}\right\rangle, \frac{1}{s}\left\langle I^{\prime}\left(\bar{\eta}_{0}(t, s)\right),\left(\bar{\eta}_{0}(t, s)\right)^{-}\right\rangle\right)
$$

We see that, for all $(t, s) \in \partial D$, by (3.6 and (a) for $\tau=\tau_{0}$, it holds

$$
\Psi_{\tau_{0}}(t, s)=\left(\left\langle I^{\prime}\left(t w^{+}+s w^{-}\right), w^{+}\right\rangle,\left\langle I^{\prime}\left(t w^{+}+s w^{-}\right), w^{-}\right\rangle\right)=\Phi^{w}(t, s)
$$

Then, by using Brouwer's topological degree, we have

$$
\operatorname{deg}\left(\Psi_{\tau_{0}}, D,(0,0)\right)=\operatorname{deg}\left(\Phi^{w}, D,(0,0)\right)=\operatorname{sgn}\left(\operatorname{det}\left(\Phi^{w}\right)^{\prime}(1,1)\right)=1
$$

so we deduce that $\Psi_{\tau_{0}}$ has a zero $(\bar{t}, \bar{s}) \in D$, that is

$$
\left\langle I^{\prime}\left(\bar{\eta}_{0}(\bar{t}, \bar{s})\right),\left(\bar{\eta}_{0}(\bar{t}, \bar{s})\right)^{ \pm}\right\rangle=0
$$

Therefore, there exists $(\bar{t}, \bar{s}) \in D$ such that $\bar{\eta}_{0}(\bar{t}, \bar{s}) \in \mathcal{M}$ and this contradicts (3.7).

## References

[1] C. O. Alves, M. A. S. Souto; Existence of least energy nodal solution for a SchrödingerPoisson system in bounded domains, Z. Angew. Math. Phys., 65 (2014), 1153-1166.
[2] V. Ambrosio, T. Isernia; Sign-changing solutions for a class of fractional Schrödinger equations with vanishing potentials, arXiv:1609.09003.
[3] S. Barile, G. M. Figueiredo; Existence of least energy positive, negative and nodal solutions for a class of pq-problems with potentials vanishing at infinity, J. Math. Anal. Appl., 427 (2015), 1205-1233.
[4] T. Bartsch, Z. Liu, T. Weth; Sign changing solutions of superlinear Schrödinger equations, Commun. Part. Diff. Eq., 29 (2004), 25-42.
[5] T. Bartsch, T. Weth; Three nodal solutions of singularly perturbed elliptic equations on domains without topology, Ann. Inst. H. Poincaré Anal. Non Linéaire, 22 (2005), 259-281.
[6] T. Bartsch, T. Weth, M. Willem; Partial symmetry of least energy nodal solutions to some variational problems, J. Anal. Math., 96 (2005), 1-18.
[7] H. Berestycki, P. L. Lions; Nonlinear scalar field equations, II, Existence of infinitely many solutions, Arch. Rational Mech. Anal., 82 (1983), 347-375.
[8] L. Caffarelli, J. M. Roquejoffre, Y. Sire; Variational problems for free boundaries for the fractional Laplacian, J. Eur. Math. Soc., 12 (2010), 1151-1179.
[9] L. Caffarelli, S. Salsa, L. Silvestre; Regularity estimates for the solution and the free boundary of the obstacle problem for the fractional Laplacian, Invent. Math., 171 (2008), 425-461.
[10] A. Castro, J. Cossio, J. Neuberger; A sign-changing solution for a superlinear Dirichlet problem, Rocky Mountain J. Math., 27 (1997), 1041-1053.
[11] S. T. Chen, Y. B. Li, X. H. Tang; Sign-changing solutions for asymptotically linear Schrodinger equation in bounded domains, Electron. J. Differ. Eq. 317 (2016), 1-9.
[12] E. Di Nezza, G. Palatucci, E. Valdinoci; Hitchhiker's guide to the fractional sobolev spaces, Bull. Sci. Math. 136 (2012), 521-573.
[13] S. Dipierro, M. Medina, E. Valdinoci; Fractional Elliptic Problems with Critical Growth in the Whole of $\mathbb{R}^{n}$, Lecture Notes, Scuola Normale Superiore, Vol. 15, 2017, pp. VIII+158.
[14] Z. Gao, X. H. Tang, W. Zhang; Least energy sign-changing solutions for nonlinear problems involving fractional laplacian, Electron. J. Differential Eq., 238 (2016), 1-6.
[15] Z. L. Liu, J. X. Sun; Invariant Sets of Descending Flow in Critical Point Theory with Applications to Nonlinear Differential Equations, J. Diff. Eqns. 172 (2001), 257-299.
[16] C. Miranda; Un'osservazione sul teorema di Brouwer, Boll. Unione Mat. Ital., 3 (2) (1940), 5-7.
[17] E. S. Noussair, J. Wei; On the effect of the domain geometry on the existence and profile of nodal solution of some singularly perturbed semilinear Dirichlet problem, Indiana Univ. Math. J., 46 (1997), 1255-1271.
[18] M. Schechter, Z. Q. Wang, W. Zou; New Linking Theorem and Sign-Changing Solutions, Commun. Part. Diff. Eq. 29 (2005), 471-488.
[19] M. Schechter, W. Zou; Sign-changing critical points from linking type theorems, Trans. Amer. Math. Soc., 358 (2006), 5293-5318.
[20] R. Servadei, E. Valdinoci; Variational methods for non-local operators of elliptic type, Discrete Contin. Dyn. Syst., 33 (2013), 2105-2137.
[21] R. Servadei, E. Valdinoci; Mountain Pass solutions for non-local elliptic operators, J. Math. Anal. Appl., 389 (2012), 887-898.
[22] R. Servadei, E. Valdinoci; The Brezis-Nirenberg result for the fractional Laplacian, Trans. Amer. Math. Soc., 367 (2015), 67-102.
[23] R. Servadei, E. Valdinoci; A Brezis-Nirenberg result for non-local critical equations in low dimension, Commun. Pure Appl. Anal. 12 (2013), 2445-2464.
[24] R. Servadei, E. Valdinoci; Weak and viscosity solutions of the fractional Laplace equation, Publ. Mat. 58 (2014), 133-154.
[25] R. Servadei, E. Valdinoci; Fractional Laplacian equations with critical Sobolev exponent, Rev. Mat. Complut. 28 (2015), 655-676.
[26] L. Silvestre; Regularity of the obstacle problem for a fractional power of the Laplace operator, Comm. Pure Appl. Math., 60 (2006), 67-112.
[27] K. Teng, K. Wang, R. Wang; A sign-changing solution for nonlinear problems involving the fractional Laplacian, Electron. J. Differential. Eq. 109 (2015), 1-12.
[28] Z. Wang, H. S. Zhou; Radial sign-changing solution for fractional Schrödinger equation, Discrete Contin. Dyn. Syst., 36 (2016), 499-508.

Huxiao Luo
School of Mathematics and Statistics, Central South University, Changsha, Hunan 410083, China

E-mail address: luohuxiao1989@163.com


[^0]:    2010 Mathematics Subject Classification. 35R11, 58E30.
    Key words and phrases. Brouwer's degree theory; sign-changing solutions;
    non-local elliptic equations; deformation Lemma.
    © 2017 Texas State University.
    Submitted March 30, 2017. Published July 14, 2017.

