

## CRITICAL EXPONENT FOR THE ASYMPTOTIC BEHAVIOR OF RESCALED SOLUTIONS TO THE POROUS MEDIUM EQUATION

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ABSTRACT. In this article, we find that  $\mu_c \equiv 2N/(N(m-1)+2)$  is the critical exponent for the asymptotic behavior of rescaled solutions  $t^{\mu/2}u(t^\beta x, t)$  for the porous medium equation.

### 1. INTRODUCTION

In this article, we consider the asymptotic behavior of solutions to the Cauchy problem of the porous medium equation

$$\frac{\partial u}{\partial t} - \Delta u^m = 0 \quad \text{in } \mathbb{R}^N \times (0, \infty), \quad (1.1)$$

$$u(x, t) = u_0(x) \quad \text{in } \mathbb{R}^N. \quad (1.2)$$

Here the initial value satisfies

$$u_0 \in C_0^+(\mathbb{R}^N) \equiv \{\varphi \in C(\mathbb{R}^N); \lim_{|x| \rightarrow \infty} \varphi(x) = 0 \text{ and } \varphi(x) \geq 0\}$$

and  $m > 1$  is a physical constant.

Asymptotic behavior of solutions for the porous medium equation has attracted much attention of mathematicians for a long time and many interesting results have been obtained, see [2, 8, 9, 10, 12, 13, 14, 15, 18, 19, 20, 21].

Friedman and Kamin [9] first revealed the fact that if the nonnegative initial value  $u_0 \in L^1(\mathbb{R}^N)$ , then the solution  $u(x, t)$  of problem (1.1)–(1.2) satisfies

$$\lim_{t \rightarrow \infty} t^{\frac{N}{N(m-1)+2}} \|u(\cdot, t) - U_M(\cdot, t)\|_{L^\infty(\mathbb{R}^N)} = 0,$$

where  $U_M(x, t)$  is the source-type solution with the same mass  $M$  as that of  $u_0$ ; see also [10, 13].

This result means that if  $0 \leq u_0 \in L^1(\mathbb{R}^N)$ , then the  $\omega$ -limit set of rescaled solutions  $t^{\mu/2}u(t^\beta x, t)$  with  $\mu = \frac{2N}{N(m-1)+2}$  and  $\beta = \frac{1}{N(m-1)+2}$  contains one point; that is, the rescaled solutions  $t^{\frac{N}{N(m-1)+2}}u(t^{\frac{1}{N(m-1)+2}}x, t)$  possess the simple asymptotic behavior (KV point in Figure 1). However, for  $u_0 \in L^\infty(\mathbb{R}^N)$ , in 2002, Vázquez and Zuazua [14] found that the  $\omega$ -limit set of the rescaled solutions  $t^{\mu/2}u(t^\beta x, t)$

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2010 *Mathematics Subject Classification.* 35B40, 35K65.

*Key words and phrases.* Complexity; asymptotic behavior; porous medium equation.

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Submitted November 17, 2016. Published January 10, 2017.

of problem (1.1)–(1.2) with  $\mu = 0$  and  $\beta = 1/2$  may contain infinite points, i.e.,  $u(t^{1/2}, t)$  (VZ point in Figure 1) possess complicated asymptotic behavior.

Such phenomena that the different exponents of rescaled solutions  $t^{\mu/2}u(t^\beta x, t)$  show different asymptotic behaviors for the porous medium equation have been studied in [2, 14, 15, 20, 21], for other evolution equations, one can see [3, 4, 5, 6, 7, 11].

For the  $\omega$ -limit set of the rescaled solutions  $t^{\mu/2}u(t^\beta, t)$  of problem (1.1)–(1.2) in  $C_0(\mathbb{R}^N)$ , we showed in our previous paper [20] that if  $(\mu, \beta) \in I$  ( $0 < \mu < \frac{2N}{N(m-1)+2}$  and  $\beta > \beta(\mu) = \frac{2-\mu(m-1)}{4}$ ), then there exists  $u_0 \in C_0^+(\mathbb{R}^N)$  such that this  $\omega$ -limit set contains infinite points; see Figure 1). In another paper [21], we revealed that if  $\mu$  and  $\beta$  in the line segment  $\beta(\mu) = \frac{2-\mu(m-1)}{4}$  ( $0 < \mu < \frac{2N}{N(m-1)+2}$ , see Figure 1), then there also exists  $u_0 \in C_0^+(\mathbb{R}^N)$  such that this  $\omega$ -limit set contains infinite points. While in this paper, we will reveal the different fact that if  $(\mu, \beta) \in II$  ( $\mu \geq \frac{2N}{N(m-1)+2}$ ,  $\beta > 0$ ), then for any  $u_0 \in C_0^+(\mathbb{R}^N)$ , this  $\omega$ -limit set contains at most one point, see Figure 1), i.e., the complicated asymptotic behavior of the rescaled solutions cannot happen.

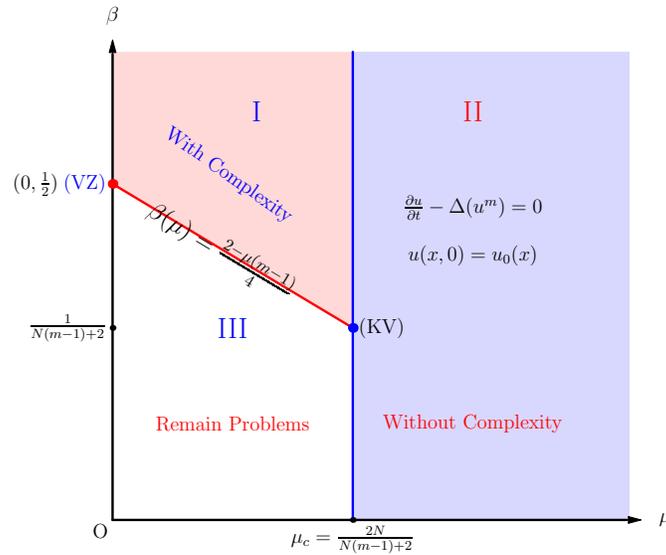


FIGURE 1. The  $\mu$ - $\beta$  Parameters Plane

**Remark 1.1.** From the above results, we can find that  $\mu_c = 2N/(N(m-1)+2)$  is the critical exponent of  $\mu$  on the asymptotic behavior of the rescaled solutions  $t^{\mu/2}u(t^\beta x, t)$ . It is not clear whether the rescaled solutions  $t^{\mu/2}u(t^\beta x, t)$  with  $(\mu, \beta) \in III$  ( $0 < \mu < 2N/(N(m-1)+2)$  and  $0 < \beta < (2-\mu(m-1))/4$ , see Figure 1) possess complicated asymptotic behavior, so the problem of the critical exponent for  $\beta$  still has not been solved.

The rest of this article is organized as following. In the next section, we introduce some definitions and concepts to give a series of lemmas. In the last of this paper, we give and prove our results.

## 2. PRELIMINARIES

Before introducing the main results of this paper, we give some concepts as in [1, 16, 17]. For  $f \in L^1_{\text{loc}}(\mathbb{R}^N)$  and  $r > 0$ , let

$$\|f\|_r = \sup_{R \geq r} R^{-\frac{N(m-1)+2}{m-1}} \int_{|x| \leq R} |f(x)| dx.$$

Then we define the space  $X = X(\mathbb{R}^N)$  by

$$X \equiv \{f \in L^1_{\text{loc}}(\mathbb{R}^N); \|f\|_1 < \infty\},$$

and equip this space with the norm  $\|\cdot\|_1$ . Hence it is a Banach space, and any norm  $\|\cdot\|_r$ ,  $r > 0$ , is an equivalent norm. For  $f \in X$ , we define

$$\ell(f) = \lim_{r \rightarrow \infty} \|f\|_r.$$

The space  $X_0 = X_0(\mathbb{R}^N)$  is defined by

$$X_0 \equiv \{f \in X; \ell(f) = 0\}.$$

Notice that  $L^1(\mathbb{R}^N) \subset X_0 \subset X \subset L^1_{\text{loc}}(\mathbb{R}^N)$  with continuous inclusions. Similarly,  $L^\infty(\mathbb{R}^N) \subset X_0$  with continuous inclusion. We now give the definition of solutions for problem (1.1)–(1.2) with the initial value  $u_0 \in X_0$ .

**Definition 2.1.** A nonnegative measurable function  $u = u(x, t)$  defined in  $S_T = [0, T) \times \mathbb{R}^N$ ,  $T > 0$ , is a solution of (1.1)–(1.2) if

- (I)  $u \in C([0, T); L^1_{\text{loc}}(\mathbb{R}^N)) \cap L^\infty(0, T; X)$ ;
- (II)  $u^m \in L^1((0, T) \times B_r(0))$  for any  $B_r(0) \equiv \{x \in \mathbb{R}^N; |x| < r, r > 0\}$ ;
- (III) for every test function  $\phi \in C_c^{2,1}(S_T)$ , it holds

$$\iint_{S_T} (u\phi_t + u^m \Delta \phi) dx dt + \int_{\mathbb{R}^N} u_0(x)\phi(x, 0) dx = 0.$$

For any  $u_0 \in X_0$ , the existence and uniqueness of the solution is well established in [1, 16, 17]. Moreover, problem (1.1)–(1.2) generates a bounded continuous semigroup in the space  $X_0$  given by

$$S(t) : u_0 \rightarrow u(x, t); \tag{2.1}$$

that is,  $S(t)u_0 \in C([0, \infty); X_0)$ , see [16, 17]. We now introduce the definitions of scalings and present the commutative relations between the semigroup operators and the dilation operators as in [20, 21]. For  $\lambda, \mu, \beta > 0$  and  $u_0 \in X_0$ , the space-time dilation  $\Gamma_\lambda^{\mu, \beta}$  is defined as following:

$$\Gamma_\lambda^{\mu, \beta}[u_0](x) \equiv D_\lambda^{\mu, \beta}[S(\lambda^2 t)u_0(x)] = \lambda^\mu u(\lambda^{2\beta} x, \lambda^2 t),$$

where the dilation  $D_\lambda^{\mu, \beta}$  is defined as

$$D_\lambda^{\mu, \beta} w(x) \equiv \lambda^\mu w(\lambda^{2\beta} x)$$

and  $S(t)$  is the PME semigroup given by (2.1). From the definitions of  $D_\lambda^{\mu, \beta}$  and  $S(t)$ , we can get the following commutative relations between the semigroup operators  $S(t)$  and the dilation operators  $D_\lambda^{\mu, \beta}$ ,

$$\Gamma_\lambda^{\mu, \beta} u_0(x) = D_\lambda^{\mu, \beta}[S(\lambda^2 t)u_0(x)] = S(\lambda^{2-4\beta-\mu(m-1)} t)[D_\lambda^{\mu, \beta} u_0](x).$$

In particular,

$$\Gamma_{\sqrt{t}}^{\mu, \beta} u_0(x) = S(t^{\frac{2-4\beta-\mu(m-1)}{2}})[D_{\sqrt{t}}^{\mu, \beta} u_0](x), \tag{2.2}$$

see details in [20, 21]. The set of functions

$$\omega^{\mu,\beta}(u_0) \equiv \{f \in C_0^+(\mathbb{R}^N); \exists t_n \rightarrow \infty \text{ s.t. } D_{\sqrt{t_n}}^{\mu,\beta}[S(t_n)u_0](\cdot) \xrightarrow{t_n \rightarrow \infty} f \text{ in } L^\infty(\mathbb{R}^N)\}$$

is called  $\Omega$ -limit set. We also introduce the following symbol to denote the positive set of  $u(x, t)$  at time  $t$ ,

$$\Omega(t) \equiv \{x \in \mathbb{R}^N; u(x, t) > 0\}.$$

The  $\rho$ -neighborhood of the set  $\Omega(t)$  is defined as

$$\Omega_\rho(t) \equiv \{x \in \mathbb{R}^N; d(x, \Omega(t)) \leq \rho\},$$

where  $d(x, \Omega(t))$  is the distance from  $x$  to  $\Omega(t)$ . We now list some important properties of the solutions.

**Lemma 2.2** ([16]). *If  $0 \leq u_0 \in L^1(\mathbb{R}^N)$ , then the solution  $u(x, t)$  satisfies the  $L^1$ - $L^\infty$  smoothing effect: for every  $t > 0$ ,*

$$\|u(\cdot, t)\|_{L^\infty(\mathbb{R}^N)} \leq C_1 \|u_0\|_{L^1(\mathbb{R}^N)}^{\frac{2}{N(m-1)+2}} t^{-\frac{N}{N(m-1)+2}},$$

where  $C_1$  is a constant dependent on  $m$  and  $N$ .

The following lemma was proved in [20], we give here a different proof for the sake of completeness.

**Lemma 2.3** ([20]). *Let  $u(x, t)$  be a nonnegative solution of (1.1)–(1.2) with the initial value  $u_0$  such that  $0 \leq u_0 \in L^1(\mathbb{R}^N)$ . Then for any  $0 \leq t_1 < t_2 < \infty$ ,*

$$\Omega(t_2) \subset \Omega_{\rho(t_2-t_1)}(t_1),$$

where

$$\rho(t_2 - t_1) = C_2 (t_2 - t_1)^{\frac{1}{N(m-1)+2}} \|u_0\|_{L^1(\mathbb{R}^N)}^{\frac{m-1}{N(m-1)+2}}$$

and  $C_2$  is a constant dependent on  $m$  and  $N$ .

*Proof.* To prove this lemma, we need the fact that if  $u(x, t)$  is a nonnegative solution of (1.1)–(1.2) with the initial data  $u_0$  satisfying

$$0 \leq u_0 \in L^\infty(\mathbb{R}^N),$$

then

$$\Omega(t_2) \subset \Omega_{\rho(t_2-t_1)}(t_1) \quad \text{for } 0 \leq t_1 < t_2 < \infty, \tag{2.3}$$

where

$$\rho(t_2 - t_1) = C(t_2 - t_1)^{1/2} \|u_0\|_{L^\infty(\mathbb{R}^N)}^{\frac{m-1}{2}}.$$

In fact, for any given  $x_0 \in \mathbb{R}^N$  with  $d(x_0) > 0$ , if  $R \geq d(x_0)$ , then

$$\begin{aligned} R^{-\frac{N(m-1)+2}{m-1}} \int_{B_R(x_0)} u_0(y) dy &\leq C \|u_0\|_{L^\infty(\mathbb{R}^N)} R^{-\frac{N(m-1)+2}{m-1}} R^N \\ &= C \|u_0\|_{L^\infty(\mathbb{R}^N)} R^{-\frac{2}{m-1}} \\ &\leq C \|u_0\|_{L^\infty(\mathbb{R}^N)} d(x_0)^{-\frac{2}{m-1}}; \end{aligned}$$

or if  $R < d(x_0)$ , then

$$\int_{B_R(x_0)} u_0(y) dy = 0,$$

where  $B_R(x_0) = \{y; |x_0 - y| < R\}$ . So

$$B(x_0) \equiv \sup_{R \geq d(x_0)} R^{-\frac{N(m-1)+2}{m-1}} \int_{B_R(x_0)} u_0(y) dy \leq C \|u_0\|_{L^\infty(\mathbb{R}^N)} d(x_0)^{-\frac{2}{m-1}}. \quad (2.4)$$

The condition  $0 \leq u_0 \in L^\infty(\mathbb{R}^N) \subset X_0$  implies that if  $|x| \leq R$  and  $r \leq R$ , then

$$u(x, t) \leq C t^{-\frac{N}{N(m-1)+2}} R^{\frac{2}{m-1}} \|u_0\|_{L^\infty(\mathbb{R}^N)}^{\frac{2}{N(m-1)+2}} \quad \text{for } 0 < t < \infty,$$

see [1, 16]. This result and (2.4) imply that

$$u(x_0, t) = 0 \text{ for all } 0 \leq t \leq C \|u_0\|_{L^\infty(\mathbb{R}^N)}^{-(m-1)} d(x_0)^2.$$

This implies  $\Omega(t) \subset \Omega_{\rho(t)}(0)$ , where

$$\rho(t) = C \|u_0\|_{L^\infty(\mathbb{R}^N)}^{\frac{m-1}{2}} t^{1/2}.$$

From this, we can get the desired result.

We now discuss the case that  $0 \leq u_0 \in L^1(\mathbb{R}^N)$  to complete the proof. Without loss of generality, we can restrict our consideration to the case of  $t_1 = 0$ . For any  $0 < t < \infty$ , we select a sequence of times

$$t_k = 2^{-k}t \rightarrow 0 \quad \text{as } k \rightarrow \infty.$$

We then consider the evolution in the time intervals  $I_k = [t_k, t_{k-1}]$ ; that is, we will estimate the increase of the support in these time intervals. From the  $L^1$ - $L^\infty$  smoothing effect, at each initial time  $t = t_k$ , we have

$$\|u(t_k)\|_{L^\infty(\mathbb{R}^N)} \leq C(p, N) \|u_0\|_{L^1(\mathbb{R}^N)}^{\frac{2}{N(m-1)+2}} t_k^{-\frac{N}{N(m-1)+2}}. \quad (2.5)$$

Therefore, we can deduce from (2.3) that

$$\Omega(t_{k-1}) \subset \Omega_{\rho(t_{k-1}-t_k)}(t_k),$$

where  $\rho(t_{k-1} - t_k) = C \|u(t_k)\|_{L^\infty(\mathbb{R}^N)}^{\frac{m-1}{2}} (t_{k-1} - t_k)^{1/2}$ . Iterating, we have

$$\Omega(t) \subset \Omega_{\rho(t)}(0),$$

where

$$\begin{aligned} \rho(t) &= C \sum_{k=1}^{\infty} \|u(t_k)\|_{L^\infty(\mathbb{R}^N)}^{\frac{m-1}{2}} (t_{k-1} - t_k)^{1/2} \leq C \sum_{k=1}^{\infty} \|u_0\|_{L^1(\mathbb{R}^N)}^{\frac{m-1}{N(m-1)+2}} t_k^{\frac{1}{N(m-1)+2}} \\ &= C \|u_0\|_{L^1(\mathbb{R}^N)}^{\frac{m-1}{N(m-1)+2}} t^{\frac{1}{N(m-1)+2}} \sum_{k=1}^{\infty} 2^{-\frac{k}{N(m-1)+2}} \leq C \|u_0\|_{L^1(\mathbb{R}^N)}^{\frac{m-1}{N(m-1)+2}} t^{\frac{1}{N(m-1)+2}}. \end{aligned}$$

Here we have used the estimates (2.5). The proof is complete.  $\square$

The next lemma is called Aleksandrov’s reflection (see [16]). We introduce some notation to give this principle. Any  $H$ , hyperplane of  $\mathbb{R}^N$ , divides  $\mathbb{R}^N$  into two half spaces  $\Omega_1(H)$  and  $\Omega_2(H)$ . We denote by  $\pi = \pi_H$  the specular symmetry that maps a point  $x \in \Omega_1(H)$  into its symmetric image with respect to  $H$ ,  $\pi_H(x) \in \Omega_2(H)$ .

**Lemma 2.4** (Aleksandrov’s Reflection Principle [16]). *Let  $u \geq 0$  be a solution of problem (1.1)–(1.2) with initial value  $u_0 \in X_0$ . Suppose that for a given hyperplane  $H$  and all  $x \in \Omega_1(H)$ ,*

$$u_0(\pi_H(x)) \leq u_0(x).$$

Then, for all times  $0 \leq t < \infty$ ,

$$u(\pi_H(x), t) \leq u(x, t), \quad x \in \Omega_1(H).$$

The following lemma depends on Lemma 2.3 and 2.4.

**Lemma 2.5.** *Suppose  $u(x, t)$  is a non-negative solution of (1.1)–(1.2) with initial-value  $u_0 \in C_0^+(\mathbb{R}^N)$  and  $u_0 \not\equiv 0$ . Let*

$$M(t) = \int_{|x| \leq t^{\frac{1}{2N(m-1)+4}}} u_0(x) dx.$$

Then there exists a  $0 < t_0 < \infty$  such that for  $t \geq t_0$ ,

$$u(0, t) \geq Ct^{-\frac{N}{N(m-1)+2}} M(t)^{\frac{2}{N(m-1)+2}}.$$

*Proof.* Since the nonnegative initial value  $u_0 \not\equiv 0$  and  $u_0 \in C(\mathbb{R}^N)$ , then there exist constants  $t_1, C_3 > 0$  such that

$$\int_{B_{t_1}} u_0(x) dx \geq C_3.$$

Now let

$$t_2 = C_2^{-\frac{2}{N(m-1)+2}} C_3^{-2m+2},$$

$$t_3 = (2^{N+1} C_1 |B_1|)^{\frac{2N(m-1)+4}{N}} C_3^{-2m+2}$$

where  $C_1, C_2$  are the constants given in Lemma 2.2 and Lemma 2.3 respectively. Let  $t_0 = \max(t_1, t_2, t_3)$ . Then for any  $t \geq t_0$ , using comparison principle, we can suppose that  $u_0$  is supported in the ball  $B_t = \{x; |x| \leq t^{\frac{1}{2N(m-1)+4}}\}$ . In fact, for general  $u_0$ , suppose  $\eta_t(x)$  is a cut-off function compactly supported in  $B_t$  and less than one with

$$\int_{B_t} \eta_t(x) u_0(x) dx \geq \frac{1}{2} M(t),$$

then  $u_0 \eta_t$  is lesser than  $u_0$ . Therefore, if  $v$  is the solution with initial data  $u_0 \eta_t$ , then

$$v(x, s) \leq u(x, s) \quad \text{for all } s > 0.$$

Hence, if this lemma holds for  $v(x, t)$ , then

$$u(0, t) \geq v(0, t) \geq C \left(\frac{1}{2} M(t)\right)^{\frac{2}{N(m-1)+2}} t^{-\frac{N}{N(m-1)+2}}.$$

Therefore, in the next part of this proof, we assume that  $\text{supp } u_0 \subset B_t$ . So,

$$M(t) = \int_{\mathbb{R}^N} u_0(x) dx \geq C_3.$$

The  $L^1$ - $L^\infty$  smoothing effect implies that for any  $s > 0$ ,

$$0 \leq u(x, s) \leq C_1 M(t)^{\frac{2}{N(m-1)+2}} s^{-\frac{N}{N(m-1)+2}}.$$

The conservation of mass means that for all  $s \geq 0$ ,

$$\begin{aligned} \int_{\mathbb{R}^N} u_0(x) dx &= \int_{\mathbb{R}^N} u(x, s) dx \\ &= \int_{|x| \geq 2t^{\frac{1}{2N(m-1)+4}}} u(x, s) dx + \int_{|x| \leq 2t^{\frac{1}{2N(m-1)+4}}} u(x, s) dx, \end{aligned}$$

the last term can be estimated as

$$\int_{|x| \leq 2t^{\frac{1}{2N(m-1)+4}}} u(x, s) dx \leq 2^N C_1 |B_1| M(t)^{\frac{2}{N(m-1)+2}} s^{-\frac{N}{N(m-1)+2}} t^{\frac{N}{2N(m-1)+4}}, \quad (2.6)$$

where  $|B_1|$  is the measure of the unit ball  $B_1$  in  $\mathbb{R}^N$ . Since  $\text{supp } u_0 \subset B_t$ , then Lemma 2.3 indicates that for all  $s > 0$ ,

$$\text{supp } u(x, s) \subset B_{R_1(s)},$$

where  $R_1(s) = t^{\frac{1}{2N(m-1)+4}} + C_2 M(t)^{\frac{m-1}{N(m-1)+2}} s^{\frac{1}{N(m-1)+2}}$ . Let  $s = t$  and

$$R(t) = 4C_2 M(t)^{\frac{m-1}{N(m-1)+2}} t^{\frac{1}{N(m-1)+2}}.$$

Notice that  $t \geq t_0 \geq t_2 = C_2^{-\frac{2}{N(m-1)+2}} C_3^{-2m+2}$  and  $M(t) \geq C_3$ . So

$$R(t) > 2R_1(t) \geq 4t^{\frac{1}{2N(m-1)+4}}. \quad (2.7)$$

The hypothesis  $\text{supp } u_0 \subset B_t$  implies, via the Aleksandrov reflection principle (Lemma 2.4), that for all  $|x| \geq 2t^{\frac{1}{2N(m-1)+4}}$  and  $s \geq 0$ ,

$$u(0, s) \geq u(x, s).$$

So, from (2.7), we have

$$\begin{aligned} u(0, t)R(t)^N &\geq u(0, t)(R(t)^N - 2^N t^{\frac{N}{2N(m-1)+4}}) \\ &= \frac{1}{|B_1|} \int_{2t^{\frac{1}{2N(m-1)+4}} \leq |x| \leq R(t)} u(0, t) dx \\ &\geq \frac{1}{|B_1|} \int_{2t^{\frac{1}{2N(m-1)+4}} \leq |x| \leq R(t)} u(x, t) dx \\ &= \frac{1}{|B_1|} \int_{|x| \geq 2t^{\frac{1}{2N(m-1)+4}}} u(x, t) dx \\ &= \frac{1}{|B_1|} \int_{\mathbb{R}^N} u(x, t) dx - \frac{1}{|B_1|} \int_{|x| < 2t^{\frac{1}{2N(m-1)+4}}} u(x, t) dx. \end{aligned}$$

Now using estimate (2.6) and  $t \geq t_0 \geq t_3$ , we obtain

$$u(0, t)R(t)^N \geq \frac{1}{|B_1|} [M(t) - 2^N C_1 |B_1| M(t)^{\frac{2}{N(m-1)+2}} t^{-\frac{N}{2N(m-1)+4}}] \geq \frac{1}{2|B_1|} M(t).$$

It follows from the definition of  $R(t)$  that

$$u(0, t) \geq C t^{-\frac{N}{N(m-1)+2}} M(t)^{\frac{2}{N(m-1)+2}}.$$

The proof is complete. □

### 3. RESULTS AND THEIR PROOFS

**Theorem 3.1.** *Let  $u_0 \in C_0^+(\mathbb{R}^N)$ ,  $u_0 \not\equiv 0$ . If there exist  $0 \not\equiv v \in C_0(\mathbb{R}^N)$ ,  $\mu_0 \geq \frac{2N}{N(m-1)+2}$ ,  $\beta_0 > 0$  and a sequence  $\{t_n\}_{n=1}^\infty$  with  $\lim_{n \rightarrow \infty} t_n = +\infty$  such that*

$$\Gamma_{\sqrt{t_n}}^{\mu_0, \beta_0} u_0 = t_n^{\frac{\mu_0}{2}} [S(t_n)u_0](t_n^{\beta_0 \cdot}) \xrightarrow{t_n \rightarrow \infty} v \quad \text{in } C_0(\mathbb{R}^N), \quad (3.1)$$

then

$$u_0 \in L^1(\mathbb{R}^N), \quad \mu_0 = \frac{2N}{N(m-1)+2},$$

$$\beta_0 = \frac{2 - \mu_0[m - 1]}{4} = \frac{1}{N(m - 1) + 2}.$$

In other words, if  $\mu > \frac{2N}{N(m-1)+2}$ , or if  $\mu = \frac{2N}{N(m-1)+2}$  and  $\beta \neq \frac{1}{N(m-1)+2}$ , then

$$\omega(u_0) = \emptyset, \quad \text{or} \quad \omega(u_0) = \{0\}.$$

*Proof.* It follows from (3.1) and Lemma 2.5 that if  $n$  sufficiently large, then

$$v(0) + 1 \geq [\Gamma^{\mu_0, \beta_0} u_0](0) = t_n^{\frac{\mu_0}{2}} [S(t_n)u_0](0) \geq Ct_n^{\frac{\mu_0 - \frac{2N}{N(m-1)+2}}{2}} M(t_n)^{\frac{2}{N(m-1)+2}}. \quad (3.2)$$

Here  $M(t)$  is given by Lemma 2.5. Letting  $n \rightarrow \infty$ , we conclude that

$$\mu_0 = \frac{2N}{N(m - 1) + 2}$$

and  $u_0 \in L^1(\mathbb{R}^N)$ . Notice also that  $u_0 \geq 0$ . This gives

$$D^{\frac{2N}{N(m-1)+2}, \frac{1}{N(m-1)+2}} S(t)u_0(x) = t^{\frac{N}{N(m-1)+2}} u(t^{\frac{1}{N(m-1)+2}} x, t) \rightarrow U_M(x, 1) \quad (3.3)$$

uniformly on  $\mathbb{R}^N$  as  $t \rightarrow \infty$ . Here  $U_M(x, t)$  is the source-type solution with the same mass as that of  $u_0$ , where  $M = \int_{\mathbb{R}^N} u_0(x) dx$ , see [10, 13]. Therefore,

$$D^{\frac{2N}{N(m-1)+2}, \beta_0} S(t_n)u_0(x) - U_M(x t_n^{\beta_0 - \frac{1}{N(m-1)+2}}, 1) \xrightarrow{n \rightarrow \infty} 0 \quad (3.4)$$

uniformly on  $\mathbb{R}^N$ . The expression of the source-type solution clearly means

$$\text{supp}(U_M(x, 1)) \subset \{x; |x| \leq CM^{\frac{m-1}{N(m-1)+2}}\},$$

so that if  $\beta_0 > \frac{1}{N(m-1)+2}$ , then

$$U_M(x t_n^{\beta_0 - \frac{1}{N(m-1)+2}}, 1) \rightarrow 0 \quad \text{for all } x \neq 0$$

as  $t_n \rightarrow \infty$ . Notice also that  $v \not\equiv 0$ , so (3.4) is compatible with (3.1) only if

$$\beta_0 \leq \frac{1}{N(m - 1) + 2}.$$

On the other hand, from (3.1) and (3.3) we deduce that

$$D^{\frac{2N}{N(m-1)+2}, \frac{1}{N(m-1)+2}} S(t_n)u_0(x) - v(t_n^{\frac{1}{N(m-1)+2} - \beta_0} x) \rightarrow 0 \quad (3.5)$$

uniformly on  $\mathbb{R}^N$  as  $t_n \rightarrow \infty$ . The hypothesis that  $v \in C_0(\mathbb{R}^N)$  clearly implies that if  $\beta_0 < \frac{1}{N(m-1)+2}$ , then

$$v(t_n^{\frac{1}{N(m-1)+2} - \beta_0} x) \rightarrow 0 \quad \text{for all } x \neq 0$$

as  $t_n \rightarrow \infty$ . Recall that  $u_0 \not\equiv 0$ , so  $U_M \not\equiv 0$ . Therefore, (3.5) is compatible with (3.3) only if

$$\beta_0 \geq \frac{1}{N(m - 1) + 2}.$$

Hence

$$\beta_0 = \frac{1}{N(m - 1) + 2}.$$

So that  $\omega^{\mu, \beta}(u_0) = \emptyset$  if  $\mu > \frac{2N}{N(m-1)+2}$ , or if  $\mu = \frac{2N}{N(m-1)+2}$  and  $\beta \neq \frac{1}{N(m-1)+2}$ . This completes the proof.  $\square$

**Theorem 3.2.** *Let*

$$\mu = \frac{2N}{N(m-1)+2}, \quad \text{and} \quad \beta = \frac{1}{N(m-1)+2}.$$

If  $u_0 \in C_0^+(\mathbb{R}^N)$ , then

$$\omega^{\mu,\beta}(u_0) = \emptyset, \quad \text{or} \quad \omega^{\mu,\beta}(u_0) = \{U_M(x, 1)\},$$

where  $U_M(x, t)$  is source-type solution with the same mass  $M$  as that of  $u_0$ .

*Proof.* If  $u_0 \in C_0^+(\mathbb{R}^N)$ , then  $u_0 \in L^1(\mathbb{R}^N)$ , or else  $u_0 \in L_{\text{loc}}^1(\mathbb{R}^N)$  with  $\|u_0\|_{L^1(\mathbb{R}^N)} = \infty$ . If  $u_0 \in L^1(\mathbb{R}^N)$ , then

$$\lim_{t \rightarrow \infty} t^{\frac{N}{N(m-1)+2}} u(t^{\frac{1}{N(m-1)+2}} x, t) = U_M(x, 1) \quad \text{in } L^\infty(\mathbb{R}^N). \quad (3.6)$$

So

$$\omega^{\mu,\beta}(u_0) = \{U_M(x, 1)\}.$$

If  $u_0 \in L_{\text{loc}}^1(\mathbb{R}^N)$  and  $\|u_0\|_{L^1(\mathbb{R}^N)} = \infty$ , approximating  $u_0$  by an increasing sequence of integrable data  $u_{0n}$ , applying (3.6) and passing to the limit, we have

$$\lim_{t \rightarrow \infty} t^{\frac{N}{N(m-1)+2}} u(t^{\frac{1}{N(m-1)+2}} x, t) = \infty \quad \text{in } L^\infty(\mathbb{R}^N).$$

Hence  $\omega^{\mu,\beta}(u_0) = \emptyset$ . The proof is complete.  $\square$

**Remark 3.3.** As we had showed in [20, 21] that for  $0 < \mu < 2N/(N(m-1)+2)$ , if  $\beta = (2 - \mu(m-1))/4$ , then there exists an initial value  $u_0 \in C_0^+(\mathbb{R}^N)$  such that the  $\Omega$ -limit set  $\omega^{\mu,\beta}(u_0)$  contains the set

$$S(1)C_0^+(\mathbb{R}^N) \equiv \{S(1)\varphi; \varphi \in C_0^+(\mathbb{R}^N)\},$$

or if  $\beta > \frac{2-\mu(m-1)}{4}$ , then there also exists an initial value  $u_0 \in C_0^+(\mathbb{R}^N)$  such that the  $\Omega$ -limit set  $\omega^{\mu,\beta}(u_0)$  contains the set

$$C_0^{+,0}(\mathbb{R}^N) \equiv \{\varphi \in C_0^+(\mathbb{R}^N); \varphi(0) = 0\}.$$

Therefore,

$$\mu_c = \frac{2N}{N(m-1)+2}$$

is the critical exponent of  $\mu$  on the asymptotic behavior of the rescaled solutions  $t^{\mu/2}u(t^\beta \cdot, t)$ .

**Acknowledgements.** This research was supported by the NSFC (11071099 and 11371153), Natural Science Foundation Project of CQ (cstc2016jcyjA0596), Scientific and Technological Research Program of Chongqing Municipal Education Commission (KJ1401003, KJ1601006), and Innovation Team Building at Institutions of Higher Education in Chongqing (CXTDX201601035).

#### REFERENCES

- [1] P. Bénilan, M. G. Crandall, M. Pierre; *Solutions of the porous medium in  $\mathbb{R}^N$  under optimal conditions on the initial-values*, Indiana Univ. Math. J., **33** (1984), 51–87.
- [2] J. A. Carrillo, J. L. Vázquez; *Asymptotic complexity in filtration equations*, J. Evol. Equ., **7** (2007), 471–495.
- [3] T. Cazenave, F. Dickstein, F. B. Weissler; *Universal solutions of a nonlinear heat equation on  $\mathbb{R}^N$* , Ann. Scuola Norm. Sup. Pisa Cl. Sci., **5** (2003), 77–117.

- [4] T. Cazenave, F. Dickstein, F.B. Weissler; *Multiscale asymptotic behavior of a solution of the heat equation in  $\mathbb{R}^N$* , In *Nonlinear Differential Equations: A Tribute to D. G. de Figueiredo*, *Progress in Nonlinear Differential Equations and their Applications*, Birkhäuser Verlag, Basel, **66**(5) (2005), 185–194.
- [5] T. Cazenave, F. Dickstein, F. B. Weissler; *Chaotic behavior of solutions of the Navier-Stokes system in  $\mathbb{R}^N$* , *Adv. Differential Equations*, **10** (2005), 361–398.
- [6] T. Cazenave, F. Dickstein, F. B. Weissler; *A solution of the constant coefficient heat equation on  $\mathbb{R}$  with exceptional asymptotic behavior: an explicit constuction*, *J. Math. Pures Appl.*, **85**(1) (2006), 119–150.
- [7] T. Cazenave, F. Dickstein, F. B. Weissler; *Nonparabolic asymptotic limits of solutions of the heat equation on  $\mathbb{R}^N$* , *J. Dynam. Differential Equations*, **19** (2007), 789–818.
- [8] E. DiBenedetto; *Degenerate parabolic equations*, Springer-Verlag, New York, (1993).
- [9] A. Friedman, S. Kamin; *The asymptotic behavior of gas in an  $N$ -dimensional porous medium*, *Trans. Amer. Math. Soc.*, **262** (1980), 551–563.
- [10] S. Kamin, J. L. Vázquez; *Fundamental solutions and asymptotic behaviour for the  $p$ -Laplacian equation*, *Rev. Mat. Iberoamericana*, **4** (1988), 339–354.
- [11] H. Mouajria, S. Tayachi, F.B. Weissler; *The heat semigroup on sectorial domains, highly singular initial values and applications*, *J. Evol. Equ.*, **16**(2) (2016), 341–364.
- [12] G. Toscani; *Entropy dissipation and the rate of convergence to equilibrium for the Fokker-Planck equation*, *Quart. Appl. Math.*, **67** (1999), 521–541.
- [13] J. L. Vázquez; *Asymptotic behavior for the porous medium equation in the whole space*, *J. Evol. Equ.*, **3** (2003), 67–118.
- [14] J. L. Vázquez, E. Zuazua; *Complexity of large time behaviour of evolution equations with bounded data*, *Chin. Ann. Math. Ser. B*, **23** (2002), 293–310.
- [15] J. L. Vázquez, M. Winkler; *Highly time-oscillating solutions for very fast diffusion equations*, *J. Evol. Equ.*, **11** (3) (2011), 725–742.
- [16] J. L. Vázquez; *The Porous Medium Equation, Mathematical Theory*, *Oxford Mathematical Monographs*, The Clarendon Press/Oxford University Press, Oxford/New York, (2007).
- [17] J. L. Vázquez; *Smoothing and decay estimates for nonlinear parabolic equations, Equations of porous medium type*, Oxford University Press, (2006).
- [18] L. W. Wang, J. X. Yin, C. H. Jin;  *$\omega$ -limit sets for porous medium equation with initial data in some weighted spaces*, *Discrete Contin. Dyn. Syst. Ser. B*, **18** (2013), 223–236.
- [19] Z.Q. Wu, J. N. Zhao, J. X. Yin, H. L. Li; *Nonlinear Diffusion Equations*, World Scientific, Singapore, (2001).
- [20] J. X. Yin, L. W. Wang, R. Huang; *Complexity of asymptotic behavior of the porous medium equation in  $\mathbb{R}^N$* , *J. Evol. Equ.*, **11** (2011), 429–455.
- [21] J. X. Yin, L. W. Wang, R. Huang; *Complexity of asymptotic behavior of Solutions for the porous medium equation with Absorption*, *Acta Math. Sci. Ser. B Engl. Ed.*, **6** (2010) 1865–1880.
- [22] J. N. Zhao, H. J. Yuan; *Lipschitz continuity of solutions and interfaces of the evolution  $p$ -Laplacian equation*, *Northeast. Math. J.*, **8** (1) (1992), 21–37.

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