

## EXISTENCE OF SOLUTIONS TO NONLINEAR FRACTIONAL SCHRÖDINGER EQUATIONS WITH SINGULAR POTENTIALS

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ABSTRACT. We study the eigenvalue problem

$$(-\Delta)^s u(x) + V(x)u(x) - K(x)|u|^{p-2}u(x) = \lambda u(x) \quad \text{in } \mathbb{R}^N,$$

where  $s \in (0, 1)$ ,  $N > 2s$ ,  $2 < p < 2^* = \frac{2N}{N-2s}$ ,  $V(x)$  is indefinite and allowed to be unbounded from below, and  $K(x)$  is nonnegative and allowed to be unbounded from above. When  $\lambda < \lambda_0 = \inf \sigma((-\Delta)^s + V(x))$  (the lowest spectrum of the operator  $(-\Delta)^s + V(x)$ ), we obtain a positive ground state solution by using the constrained minimization method. Also we discuss the regularity of solutions.

### 1. INTRODUCTION AND STATEMENT OF MAIN RESULTS

In this article, we consider standing waves of the nonlinear fractional Schrödinger equation

$$i\psi_t = (-\Delta)^s \psi + V(x)\psi - K(x)|\psi|^{p-2}\psi \quad \text{in } \mathbb{R}^N, \quad (1.1)$$

where  $(x, t) \in \mathbb{R}^N \times (0, \infty)$ ,  $0 < s < 1$ ,  $V(x)$  and  $K(x)$  are some real functions. The operator  $(-\Delta)^s$  is the fractional Laplacian of order  $s$ .

This equation was introduced by Laskin [8, 9], and comes from fractional quantum mechanics for the study of particles on stochastic fields modelled by Lévy process. The Lévy process is widely used to model a variety of processes, such as turbulence, financial dynamics, biology and physiology, see [7, 11, 19]. When  $s = 1$ , the Lévy process becomes the Brownian motion, and the equation (1.1) reduces to the classical Schrödinger equation

$$i\psi_t = -\Delta\psi + V(x)\psi - K(x)|\psi|^{p-2}\psi \quad \text{in } \mathbb{R}^N. \quad (1.2)$$

Standing wave solutions to this equation are solutions of the form  $\psi(x, t) = e^{-i\lambda t}u(x)$  where  $u(x)$  satisfies the equation

$$-\Delta u + (V(x) - \lambda)u - K(x)|u|^{p-2}u = 0 \quad \text{in } \mathbb{R}^N, \quad (1.3)$$

which has been extensively studied in the past 20 years. We mention some earlier work here. Oh [12] studied positive multi-lump bound states, and it was assumed that  $K(x) \equiv \gamma$  for some  $\gamma > 0$ , and  $V(x)$  belongs to a class of potentials  $(V)_a$  for some  $a$  and  $\lambda < a$  ( $V \in (V)_a$  if either  $V(x) \equiv a$  or  $V(x) > a$  for all  $x \in \mathbb{R}^N$  and  $(V(x) - a)^{-1/2} \in Lip(\mathbb{R}^N)$ ). Rabinowitz [15] investigated the ground state solutions

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of the problem (1.3) under the condition  $\inf_{\mathbb{R}^N} V(x) > \lambda$  and after this Byeon and Wang [1] considered the case  $\inf_{\mathbb{R}^N} V(x) = \lambda$  which they call it critical frequency case.

Our goal is to look for standing wave solutions of the form  $\psi(x, t) = e^{-i\lambda t}u(x)$  to equation (1.1) for fractional order  $s \in (0, 1)$ . Precisely, we will investigate the problem.

$$\begin{aligned} (-\Delta)^s u(x) + (V(x) - \lambda)u(x) - K(x)|u|^{p-2}u(x) &= 0 \quad \text{in } \mathbb{R}^N, \\ u(x) &\in H^s(\mathbb{R}^N). \end{aligned} \quad (1.4)$$

Where  $s \in (0, 1)$ ,  $2 < p < 2^* = \frac{2N}{N-2s}$ ,  $N > 2s$ ,  $\lambda \in \mathbb{R}$ ,  $V(x)$  and  $K(x)$  are real functions satisfying the following conditions:

- (A1)  $V(x) : \mathbb{R}^N \rightarrow \mathbb{R}$ ,  $V(x) \in L^{\frac{N}{2s}}(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$ ;
- (A2) for any  $\epsilon > 0$ , the Lebesgue measure  $|\{x : |V(x)| > \epsilon\}| < \infty$ .
- (A3)  $K(x) \geq 0$ ,  $K(x) \not\equiv 0$ ,  $K(x) \in L^{\frac{2^*}{2^*-p}}(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$ ;
- (A4) for any  $\epsilon > 0$ , the Lebesgue measure  $|\{x : |K(x)| > \epsilon\}| < \infty$ .
- (A5)  $V(x) \in L^{\tilde{q}}(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$ ,  $K(x) \in L^{\tilde{r}}(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  for  $\tilde{q} > \frac{N}{2s}$  and  $\tilde{r} > \frac{2^*}{2^*-p}$ .

We remark that in [9], Laskin investigated the fractional Hydrogen-like atom where  $V(x) = -\frac{Ze^2}{|x|}$  (for  $N = 3$  and  $1/2 < s < 1$ ), and evaluated the corresponding energy spectrum. It is easy to check that such potential satisfies condition (A1).

In recent years, there have been a few results for nonlinear fractional Schrödinger equations like (1.4). Teng [18] investigated multiple solutions of the equation

$$(-\Delta)^s u + V(x)u = f(x, u), \quad (1.5)$$

for  $V(x) \in C(\mathbb{R}^N)$ ,  $\text{ess inf } V(x) > 0$ , and  $f \in C(\mathbb{R}^N \times \mathbb{R})$ . Secchi [16] studied the ground state solutions of (1.5) for the case that  $V \in C^1(\mathbb{R}^N)$ ,  $\inf_{x \in \mathbb{R}^N} V(x) = V_0 > 0$ , and  $f \in C^1(\mathbb{R}^N \times \mathbb{R})$  satisfying Ambrosetti-Rabinowitz condition. In [3], ground states and bound states of (1.5) are obtained by assuming that  $V(x) > 1$  and  $\lim_{|x| \rightarrow +\infty} V(x) = +\infty$ , and the nonlinearity is  $f(t) = |t|^{p-1}t$ . Chang [2] investigated the ground state solutions for asymptotically linear fractional Schrödinger equations. In particular, Felmer [6] studied the existence of positive solutions of (1.5) for  $V(x) \equiv 1$  and  $f(x, u)$  is superlinear and has subcritical growth with respect to  $u$  such that there exist  $1 < p < (N + 2s)/(N - 2s)$ , so that

$$f(x, \xi) \leq C(1 + |\xi|)^p \quad \text{for all } \xi \in \mathbb{R} \text{ and a.e. } x \in \mathbb{R}^N. \quad (1.6)$$

Furthermore, they discuss the regularity, decay and symmetry properties of solutions.

The nonlinearity  $K(x)|u|^{p-2}u(x)$  in this paper is quite different from (1.6), since  $K(x)$  may not be bounded by a constant  $C$ . For example,  $K(x) = \frac{1}{|x-x_0|^\alpha}$  for  $0 < \alpha < \frac{(2^*-p)N}{2^*}$ , satisfies (A3), (A4), and has singular point  $x_0 \in \mathbb{R}^N$ . On the other hand, since  $V(x)$  is indefinite, it is hard to use usual mountain pass arguments to obtain ground state solutions([6, 16, 2]), here we will use the constrained minimization method to obtain the ground state solutions.

We say that  $u \in H^s(\mathbb{R}^N)$  is a weak solution of (1.4), if for any  $\phi \in H^s(\mathbb{R}^N)$ ,

$$\int_{\mathbb{R}^N} (-\Delta)^{s/2} \bar{u} \cdot (-\Delta)^{s/2} \phi \, dx + \int_{\mathbb{R}^N} (V(x) - \lambda) \bar{u} \cdot \phi \, dx = \int_{\mathbb{R}^N} K(x) |u|^{p-2} \bar{u} \cdot \phi \, dx,$$

where  $\bar{u}(x)$  is conjugation of  $u(x)$  in the complex space  $H^s(\mathbb{R}^N)$ .

Solutions of (1.4) correspond to the critical points of the energy functional

$$\mathcal{I}(u) = \frac{1}{2} \left[ \int_{\mathbb{R}^N} |(-\Delta)^{s/2} u|^2 dx + \int_{\mathbb{R}^N} (V(x) - \lambda) |u|^2 dx \right] - \frac{1}{p} \int_{\mathbb{R}^N} K(x) |u|^p dx. \quad (1.7)$$

And a *ground state* of (1.4) is a solution that minimizes the energy functional on the Nehari manifold

$$\mathcal{N} = \left\{ u \in H^s(\mathbb{R}^N) \setminus \{0\} : \int_{\mathbb{R}^N} |(-\Delta)^{s/2} u|^2 + (V(x) - \lambda) |u|^2 dx = \int_{\mathbb{R}^N} K(x) |u|^p dx \right\}. \quad (1.8)$$

Now we state our main result.

**Theorem 1.1.** *Let  $s \in (0, 1)$ ,  $2 < p < 2^*$  ( $2^* = \frac{2N}{N-2s}$ ),  $N > 2s$ . Assume that (A1)–(A4) are satisfied. Let*

$$\begin{aligned} \lambda_0 &= \inf \sigma((-\Delta)^s + V(x)) \\ &= \inf \left\{ \int_{\mathbb{R}^N} |(-\Delta)^{s/2} \psi|^2 + V(x) |\psi|^2 dx : \psi \in H^s(\mathbb{R}^N), \|\psi\|_{L^2} = 1 \right\}, \end{aligned}$$

and assume that  $\lambda \leq 0$  and  $\lambda < \lambda_0$ . Then (1.4) admits at least one nonnegative weak solution such that this solution is a ground state.

To prove the positive property of nonnegative weak solutions, we need to take advantage of the representation formula

$$u = \mathcal{K}^\mu * f = \int_{\mathbb{R}^N} K(x - \xi) f(\xi) d\xi,$$

for some  $\mu > 0$ , and that  $u$  satisfies the equation

$$(-\Delta)^s u + \mu u = f \quad \text{in } \mathbb{R}^N,$$

where  $\mathcal{K}^\mu$  is the Bessel kernel

$$\mathcal{K}^\mu = \mathcal{F}^{-1} \left( \frac{1}{\mu + |\xi|^{2s}} \right).$$

We have the following positive property.

**Theorem 1.2.** *Under the assumptions of Theorem 1.1, let  $w(x)$  be a nonnegative ground state solution obtain in Theorem 1.1. If we further assume that  $V(x)$  is bound from above, then  $w(x)$  can be chosen positive in  $\mathbb{R}^N$ .*

The next step is to prove regularity of the weak solutions. Inspired by ideas in [6], We also use the representation formula above to discuss the regularity. We have the following result.

**Theorem 1.3.** *Let  $u(x) \in H^s(\mathbb{R}^N)$  be a solution of (1.4), assume that  $\lambda < 0$  and (A5) holds, i.e.,  $V(x) \in L^{\tilde{q}}(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$ ,  $K(x) \in L^{\tilde{r}}(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  for  $\tilde{q} > \frac{N}{2s}$  and  $\tilde{r} > \frac{2^*}{2^* - p}$ . Then  $u \in C^{0,\alpha}(\mathbb{R}^N)$  for some  $\alpha \in (0, 1)$ . Moreover,  $u(x) \rightarrow 0$  as  $|x| \rightarrow \infty$ .*

We remark that having the regularities above, by the same arguments as in [6, Theorem 1.5], it is easy to show that the positive ground state solutions  $u(x)$  behave at infinity like  $\frac{1}{|x|^{N+2s}}$ .

The rest of the article is originated as follows. In section 2 we give some preliminary and show some properties of the operator  $(-\Delta)^s + V(x)$ . In section 3 we

will show that weak convergence in  $H^s(\mathbb{R}^N)$  implies strong convergence on finite measure sets, which is important to prove our result. In section 4 we prove that weak continuity of the potential energies. In section 5 we prove the Theorem 1.1. In section 6 we give the proof of Theorem 1.2 and 1.3.

**Notation.** To coincide with the book [10], the Banach spaces  $L^p(\mathbb{R}^N)$ ,  $H^s(\mathbb{R}^N)$  used here are complex Banach spaces. And the inner product is defined by

$$(f(x), g(x)) = \int_{\mathbb{R}^N} \overline{f(x)}g(x) dx, \quad \text{for any } f(x), g(x) \in L^2(\mathbb{R}^N), \quad (1.9)$$

where  $\overline{f}$  denotes conjugation of  $f(x)$ .

$\widehat{u}$  denotes the Fourier transform of  $u \in L^2(\mathbb{R}^N)$ .

$$L^q(\mathbb{R}^N) + L^\infty(\mathbb{R}^N) := \{u = u_0 + u_1 : u_0 \in L^q(\mathbb{R}^N), u_1 \in L^\infty(\mathbb{R}^N)\}.$$

$\rightharpoonup$  denotes weakly converge.  $C^{0,\alpha}(\mathbb{R}^N)$  denotes Hölder continuous with exponent  $\alpha \in (0, 1)$ .

## 2. PRELIMINARIES

The fractional Laplacian  $(-\Delta)^s$  of a rapidly decaying test function  $u$  is defined as

$$(-\Delta)^s u(x) = C_{N,s} \text{P. V.} \int \frac{u(x) - u(y)}{|x - y|^{N+2s}} dx dy, \quad (2.1)$$

where P.V. denotes the principal value of the singular integral, and  $C_{N,s}$  is a constant.

We recall that the fractional Sobolev space  $W^{s,p}(\mathbb{R}^N)$  (e.g., see [16]) is defined for any  $p \in [1, \infty)$  and  $s \in (0, 1)$  as

$$W^{s,p}(\mathbb{R}^N) = \left\{ u \in L^p(\mathbb{R}^N) : \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{sp+N}} dx dy < \infty \right\},$$

endowed with the norm

$$\|u\|_W^{s,p} = \left( \int_{\mathbb{R}^N} |u|^p dx + \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{sp+N}} dx dy \right)^{1/p}.$$

When  $p = 2$ , these spaces are also denoted by  $H^s(\mathbb{R}^N)$ .

When  $p = 2$ , there is an equivalent definition of fractional Sobolev spaces based on Fourier analysis that

$$H^s(\mathbb{R}^N) = \left\{ u \in L^2(\mathbb{R}^N) : \int (1 + |\xi|^{2s}) |\widehat{u}(\xi)|^2 d\xi < \infty \right\},$$

$$\widehat{(-\Delta)^s u} = |\xi|^{2s} \widehat{u}, \quad \text{for } u \in H^s(\mathbb{R}^N),$$

and the norm can be equivalently written

$$\|u\|_{H^s} = \left( \|u\|_{L^2}^2 + \int |\xi|^{2s} |\widehat{u}(\xi)|^2 d\xi \right)^{1/2} = \left( \|u\|_{L^2}^2 + \|(-\Delta)^{s/2} u\|_{L^2}^2 \right)^{1/2}$$

Therefore, we see that  $H^s(\mathbb{R}^N)$  is just  $L^2(\mathbb{R}^N, d\mu)$ , where  $\mu$  is a measure defined by

$$\mu(dx) = (1 + |x|^{2s}) dx.$$

A sequence  $f^j(x)$  converges weakly to  $f(x)$  (we write  $f^j \rightharpoonup f$ ) in  $H^s(\mathbb{R}^N)$  in the following sense (see [10, §7.18] or [4]): for any  $g(x) \in H^s(\mathbb{R}^N)$ , when  $j \rightarrow \infty$ , one has

$$\frac{1}{(2\pi)^{N/2}} \int_{\mathbb{R}^N} [\widehat{f^j}(\xi) - \widehat{f}(\xi)]\widehat{g}(\xi)(1 + |\xi|^{2s}) d\xi \rightarrow 0. \tag{2.2}$$

The following lemma is obvious. There are more details about  $H^{1/2}(\mathbb{R}^N)$  in [10], so the case for general  $s \in (0, 1)$  is just the same arguments to  $H^{1/2}(\mathbb{R}^N)$ .

**Lemma 2.1.** *If a sequence  $f^j$  converges weakly to  $f$  in  $H^s(\mathbb{R}^N)$ . Then, there exists a constant  $M$  independent of number  $j$ , such that*

$$\|f^j\|_{H^s} \leq M, \quad \|f\|_{H^s} \leq M. \tag{2.3}$$

*Proof.* Since  $H^s(\mathbb{R}^N)$  is just  $L^2(\mathbb{R}^N, d\mu)$ , thus by uniform boundedness principle and Lower semicontinuity of  $L^p$ -norm respectively, we obtain (2.3).  $\square$

Now we review the Sobolev inequality and Sobolev-Gagliardo-Nirenberg inequality for fractional Sobolev spaces, we only show the case for  $H^s(\mathbb{R}^N)$ .

**Lemma 2.2** (Sobolev inequality [17]). *Let  $s \in (0, 1)$  be such that  $N > 2s$ . Then*

$$\|u\|_{L^{2^*}} \leq S_{N,s} \|(-\Delta)^{s/2} u\|_{L^2}$$

for every  $u \in H^s(\mathbb{R}^N)$ , where  $S_{N,s}$  is sharp constants depending only on  $N, s$ , and

$$2^* = \frac{2N}{N - 2s}$$

is the fractional critical exponent.

**Lemma 2.3** (Sobolev-Gagliardo-Nirenberg inequality [16]). *Let  $q \in (2, 2^*)$ . Then there exists a constant  $C > 0$  such that*

$$\|u\|_{L^q}^q \leq C \|u\|_{H^s}^{\frac{(q-2)N}{2s}} \|u\|_{L^2}^{q - \frac{(q-2)N}{2s}}$$

for every  $u \in H^s(\mathbb{R}^N)$ .

Next we show some properties of fractional Schrödinger operator  $(-\Delta)^s + V(x)$ . For any  $\psi(x) \in H^s(\mathbb{R}^N)$ , let  $\lambda_0$  be defined in Theorem 1.1, define

$$\mathcal{E}(\psi) := \|(-\Delta)^{s/2} \psi\|_{L^2}^2 + \int_{\mathbb{R}^N} V(x) |\psi|^2 dx. \tag{2.4}$$

then  $\lambda_0 = \inf\{\mathcal{E}(\psi) : \psi \in H^s(\mathbb{R}^N), \|\psi\|_{L^2} = 1\}$ . We have the following theorem.

**Theorem 2.4.** *For  $s \in (0, 1)$ ,  $N > 2s$ , if  $V(x) \in L^{\frac{N}{2s}}(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$ , then*

- (i)  $\lambda_0$  is finite.
- (ii)  $\|(-\Delta)^{s/2} \psi\|_{L^2}^2 \leq C\mathcal{E}(\psi) + D\|\psi\|_{L^2}^2$  for  $\psi \in H^s(\mathbb{R}^N)$  and suitable constants  $C$  and  $D$ .

*Proof.* Since  $V(x) \in L^{\frac{N}{2s}}(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$ , we can write  $V(x) = v(x) + w(x)$  with  $v(x) \in L^{\frac{N}{2s}}(\mathbb{R}^N)$  and  $w(x) \in L^\infty(\mathbb{R}^N)$ .

First we claim that we can choose  $v(x)$  satisfying  $\|v(x)\|_{L^{\frac{N}{2s}}} \leq \frac{1}{2}(S_{N,s})^{-2}$ . In fact, for  $M > 0$ , define  $S_v(M)$  by

$$S_v(M) = \{x \in \mathbb{R}^N : |v(x)| > M\},$$

then by Chebyshev inequality (see [5])

$$|S_v(M)| \leq \left( \frac{C \|v\|_{L^{\frac{N}{2s}}}}{M} \right)^{N/(2s)}. \quad (2.5)$$

Let  $\chi_A$  be the characteristic function on subset  $A \subset \mathbb{R}^N$ . Decompose  $v(x)$  into

$$v(x) = \chi_{S_v(M)} v(x) + (1 - \chi_{S_v(M)}) v(x).$$

Let  $v_1 = \chi_{S_v(M)} v(x)$ ,  $v_2 = (1 - \chi_{S_v(M)}) v(x)$ , then  $v_1 \in L^{\frac{N}{2s}}(\mathbb{R}^N)$ ,  $v_2 \in L^\infty(\mathbb{R}^N)$ , and by (2.5) we have  $\|v_1\|_{L^{\frac{N}{2s}}} < \frac{1}{2}(S_{N,s})^{-2}$  for large enough  $M$ . Replace  $v(x)$  by  $v_1$ , then the claim holds.

For any function  $\psi \in H^s(\mathbb{R}^N)$ , combing with  $\|v(x)\|_{L^{\frac{N}{2s}}} \leq \frac{1}{2}(S_{N,s})^{-2}$  and using Hölder inequality and Sobolev inequality (Lemma 2.2), we have

$$\begin{aligned} \left| \int_{\mathbb{R}^N} v(x) |\psi|^2 dx \right| &\leq \|v(x)\|_{L^{\frac{N}{2s}}} \|\psi\|_{L^{2s}}^2 \\ &\leq S_{N,s}^2 \|v(x)\|_{L^{\frac{N}{2s}}} \|(-\Delta)^{s/2} \psi\|_{L^2}^2 \\ &\leq \frac{1}{2} \|(-\Delta)^{s/2} \psi\|_{L^2}^2, \end{aligned}$$

it follows that

$$\begin{aligned} \mathcal{E}(\psi) &= \|(-\Delta)^{s/2} \psi\|_{L^2}^2 + \int_{\mathbb{R}^N} v(x) |\psi|^2 dx + \int_{\mathbb{R}^N} w(x) |\psi|^2 dx \\ &\geq \frac{1}{2} \|(-\Delta)^{s/2} \psi\|_{L^2}^2 - \|w(x)\|_{L^\infty} \|\psi\|_{L^2}^2 \geq -\|w(x)\|_{L^\infty} \|\psi\|_{L^2}^2, \end{aligned} \quad (2.6)$$

and we see that  $-\|w(x)\|_{L^\infty}$  is a lower bound to  $\lambda_0$ , i.e., (i) holds. Furthermore, the first inequality of (2.6) implies

$$\|(-\Delta)^{s/2} \psi\|_{L^2}^2 \leq 2(\mathcal{E}(\psi) + \|w(x)\|_{L^\infty} \|\psi\|_{L^2}^2),$$

i.e., (ii) holds.  $\square$

### 3. WEAK CONVERGENCE IMPLIES STRONG CONVERGENCE ON SMALL SETS

Consider the semigroup  $\{e^{-(\Delta)^{s,t}}\}_{t>0}$ . We know that, for any function  $f(x) \in H^s(\mathbb{R}^N)$ ,

$$e^{-\widehat{(-\Delta)^{s,t}} f(\xi)} = e^{-|\xi|^{2s} t} \hat{f}(\xi).$$

Now we define the heat kernel for  $s \in (0, 1)$ ,  $t > 0$ , and  $x \in \mathbb{R}^N$  as

$$\mathcal{H}(x, t) = \frac{1}{(2\pi)^{N/2}} \int_{\mathbb{R}^N} e^{i x \cdot \xi - t |\xi|^{2s}} d\xi, \quad (3.1)$$

and we know that

$$e^{-(\Delta)^{s,t}} f(x) = \int_{\mathbb{R}^N} \mathcal{H}(x - y, t) f(y) dy. \quad (3.2)$$

It is well known that  $\mathcal{H}(x, t)$  has the following properties, see [6, Appendix A] and references therein.

**Lemma 3.1.**  $\mathcal{H}(x, t)$  is radially symmetric in  $x$ , and there exists two constants  $c_1$  and  $c_2$  such that

$$c_1 \min \left\{ t^{-\frac{N}{2s}}, t|x|^{-N-2s} \right\} \leq \mathcal{H}(x, t) \leq c_2 \min \left\{ t^{-\frac{N}{2s}}, t|x|^{-N-2s} \right\}. \quad (3.3)$$

Now we use the properties of semigroup with respect to  $(-\Delta)^s$  to prove that weak convergence in  $H^s(\mathbb{R}^N)$  implies strong convergence on any finite measure set (not just on a bounded domain  $\Omega \in \mathbb{R}^N$ , see compact embeddings in [13, 14]). This result can also be found in [4], here we give a different proof along the ideas in [10, Theorem8.6].

**Theorem 3.2.** *Let  $\{f^j\} \subset H^s(\mathbb{R}^N)$  such that  $f^j$  converges weakly to  $f$  in  $H^s(\mathbb{R}^N)$ . Let  $A \subset \mathbb{R}^N$  be any set of finite Lebesgue measure, i.e.,  $|A| < \infty$ , and let  $\chi_A$  be its characteristic function. Then*

$$\chi_A f^j \rightarrow \chi_A f \quad \text{strongly in } L^q(\mathbb{R}^N)$$

for  $1 \leq q < 2^* = \frac{2N}{N-2s}$ , when  $N > 2s$ .

*Proof.* We take three steps to prove the theorem.

**Step 1.** We claim that, for any  $f \in H^s(\mathbb{R}^N)$ ,

$$\|f - e^{-(\Delta)^s t} f\|_{L^2} \leq \|(-\Delta)^{s/2} f\|_{L^2} \sqrt{t}. \tag{3.4}$$

In fact, we know that

$$1 - \exp[-(|\xi|)^{2s} t] \leq \min \{1, (|\xi|)^{2s} t\} \leq |\xi|^s \sqrt{t},$$

and it follows that

$$\begin{aligned} \|f - e^{-(\Delta)^s t} f\|_{L^2}^2 &= \int_{\mathbb{R}^N} |\hat{f}(\xi)|^2 (1 - \exp[-(|\xi|)^{2s} t])^2 d\xi \\ &\leq \int_{\mathbb{R}^N} |\hat{f}(\xi)|^2 (|\xi|^s \sqrt{t})^2 d\xi = \|(-\Delta)^{s/2} f\|_{L^2}^2 t, \end{aligned}$$

this proves (3.4).

**Step 2.** We first prove that  $\chi_A f^j \rightarrow \chi_A f$  strongly in  $L^2(\mathbb{R}^N)$ . Let  $g^j := e^{-(\Delta)^s t} f^j$ , by Lemma 2.1, we note that

$$\|(-\Delta)^{s/2} f^j\|_{L^2} \leq \|f^j\|_{H^s} \leq M, \tag{3.5}$$

$$\|(-\Delta)^{s/2} f\|_{L^2} \leq \|f\|_{H^s} \leq M, \tag{3.6}$$

where  $M$  is a constant independent of  $j$ . Then by (3.4) we have

$$\|f^j - g^j\|_{L^2} = \|f^j - e^{-(\Delta)^s t} f^j\|_{L^2} \leq M\sqrt{t},$$

$$\|f - g\|_{L^2} = \|f - e^{-(\Delta)^s t} f\|_{L^2} \leq M\sqrt{t}.$$

Simply note that

$$\begin{aligned} \|\chi_A(f^j - f)\|_{L^2} &\leq \|\chi_A(f^j - g^j)\|_{L^2} + \|\chi_A(g^j - g)\|_{L^2} + \|\chi_A(g - f)\|_{L^2} \\ &\leq 2M\sqrt{t} + \|\chi_A(g^j - g)\|_{L^2}. \end{aligned}$$

For  $\epsilon > 0$  given, first choose  $t > 0$  (depending on  $\epsilon$ ) such that  $2M\sqrt{t} < \epsilon/2$  and if for  $j$  (depending on  $\epsilon$ ) we have  $\|\chi_A(g^j - g)\|_{L^2} < \epsilon/2$ , then we have  $\|\chi_A(f^j - f)\|_{L^2} < \epsilon$ . Therefore, it remains to prove that  $\chi_A g^j \rightarrow \chi_A g$  strongly in  $L^2(\mathbb{R}^N)$ .

To prove  $\chi_A g^j \rightarrow \chi_A g$  strongly in  $L^2(\mathbb{R}^N)$ , first we note that, if  $|y - x| \geq t^{\frac{1}{2s}}$ , we have  $t|x - y|^{-N-2s} \leq t^{-\frac{N}{2s}}$ , and then by Lemma 3.1, we have

$$\begin{aligned} \mathcal{H}(x - y, t) &\leq c_2 t|x - y|^{-N-2s} = 2c_2 \frac{t}{2|x - y|^{N+2s}} \\ &\leq 2c_2 \frac{t}{t^{\frac{N+2s}{2s}} + |x - y|^{N+2s}}. \end{aligned} \tag{3.7}$$

Then, for every fix  $x$ , we have  $\mathcal{H}(x - y, t) \in L^{(2^*)'}(\mathbb{R}^N)$ , where  $(2^*)' = 2N/(N + 2s)$ , is dual index to  $2^*$ . In fact, let  $B(x, t^{\frac{1}{2s}})$  denote a ball center at  $x$  and has radius  $t^{\frac{1}{2s}}$ ,

$$\begin{aligned} & \int_{\mathbb{R}^N} (\mathcal{H}(x - y, t))^{(2^*)'} dy \\ &= \int_{B(x, t^{\frac{1}{2s}})} (\mathcal{H}(x - y, t))^{(2^*)'} dy + \int_{\mathbb{R}^N \setminus B(x, t^{\frac{1}{2s}})} (\mathcal{H}(x - y, t))^{(2^*)'} dy \\ &\leq \int_{B(x, t^{\frac{1}{2s}})} (c_2 t^{-\frac{N}{2s}})^{(2^*)'} dy + \int_{\mathbb{R}^N \setminus B(x, t^{\frac{1}{2s}})} (c_2 t |x - y|^{-N-2s})^{(2^*)'} dy \\ &\leq M_1 + \int_{\mathbb{R}^N} (2c_2 \frac{t}{t^{\frac{N+2s}{2s}} + |x - y|^{N+2s}})^{(2^*)'} dy \\ &\leq M_1 + \int_{\mathbb{R}^N} (2c_2 \frac{t}{t^{\frac{N+2s}{2s}} + |y|^{N+2s}})^{(2^*)'} dy \leq M_2, \end{aligned}$$

where  $M_2$  is a constant independent of  $x$ .

Since for every  $x$ , we have  $\mathcal{H}(x - y, t) \in L^{(2^*)'}(\mathbb{R}^N)$ , by Hölder inequality

$$\chi_A |g^j(x)| \leq \|\mathcal{H}(x - y, t)\|_{(2^*)'} \|f^j\|_{2^*} \chi_A(x).$$

Using Lemma 2.2 and (3.5),  $\|f^j\|_{2^*} \leq S_{N,s} \|(-\Delta)^{s/2} f^j\|_{L^2} \leq S_{N,s} M$ . Hence  $\chi_A g^j$  is dominated by a constant multiple of the square integrable function  $\chi_A(x)$ . On the other hand, if  $g^j(x)$  converges pointwise to  $g(x)$  for every  $x \in \mathbb{R}^N$ , Then by general dominated convergence theorem, we have  $\chi_A g^j \rightarrow \chi_A g$  strongly in  $L^2(\mathbb{R}^N)$ . Next we shall prove  $g^j(x)$  converges pointwise for every  $x \in \mathbb{R}^N$ . We note that, for fixed  $x$ ,

$$H(\widehat{x - y}, t)(\xi) = (e^{-ix \cdot \xi}) e^{-t|\xi|^{2s}},$$

and

$$\begin{aligned} g^j(x) &= e^{-(\Delta)^s t} f^j(x) = \int_{\mathbb{R}^N} \mathcal{H}(x - y, t) f^j(y) dy \\ &= \int_{\mathbb{R}^N} H(\widehat{x - y}, t)(\xi) \hat{f}^j(\xi) d\xi = \int_{\mathbb{R}^N} (e^{-ix \cdot \xi}) e^{-t|\xi|^{2s}} \hat{f}^j(\xi) d\xi \\ &= \int_{\mathbb{R}^N} \frac{(e^{-ix \cdot \xi}) e^{-t|\xi|^{2s}}}{1 + |\xi|^{2s}} \hat{f}^j(\xi) (1 + |\xi|^{2s}) d\xi. \end{aligned}$$

Let  $h(y)$  be a function satisfying  $\hat{h}(\xi) = \frac{(e^{-ix \cdot \xi}) e^{-t|\xi|^{2s}}}{1 + |\xi|^{2s}}$ , it is easy to see that  $h(y) \in H^s(\mathbb{R}^N)$ . Since  $f^j$  converges weakly to  $f$  in  $H^s(\mathbb{R}^N)$ , by (2.2), then we have  $g^j(x)$  converges pointwise to  $g(x)$  for every  $x \in \mathbb{R}^N$ . Hence we complete the proof of Step 2.

**Step 3.** The inequality

$$\|\chi_A(f^j - f)\|_{L^q} \leq \|\chi_A\|_{L^r} \|\chi_A(f - f^j)\|_{L^2}$$

for  $1/q = 1/r + 1/2$  proves the theorem for  $1 \leq q \leq 2$ . Again by Hölder inequality, Lemma 2.2,

$$\begin{aligned} \|\chi_A(f^j - f)\|_{L^q} &\leq \|\chi_A(f^j - f)\|_{L^2}^\alpha \|\chi_A(f - f^j)\|_{L^{2^*}}^{1-\alpha} \\ &\leq \|\chi_A(f^j - f)\|_{L^2}^\alpha \|f - f^j\|_{L^{2^*}}^{1-\alpha} \\ &\leq \|\chi_A(f^j - f)\|_{L^2}^\alpha (S_{N,s})^{1-\alpha} \|(-\Delta)^{s/2}(f - f^j)\|_{L^2}^{1-\alpha} \end{aligned}$$

$$\leq \|\chi_A(f^j - f)\|_{L^2}^\alpha (2M S_{N,s})^{1-\alpha},$$

where  $\alpha = (1/q - 1/2^*)(1/2 - 1/2^*)$ , and this proves the theorem for  $2 \leq q < 2^*$ . The proof is complete.  $\square$

4. WEAK CONTINUITY OF THE POTENTIAL ENERGIES

**Lemma 4.1.** *Let  $2 \leq q < 2^*$ ,  $F_\psi := \int_{\mathbb{R}^N} F(x)|\psi|^q dx$ ,  $F(x)$  be a real function on  $\mathbb{R}^N$  such that  $F(x) \in L^{\frac{2^*}{2^*-q}}(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  and  $|\{x : |F(x)| > \epsilon\}| < \infty$  for any  $\epsilon > 0$ . Then  $F_\psi$  is weakly continuous in  $H^s(\mathbb{R}^N)$ , i.e., if  $\psi^j \rightharpoonup \psi$  as  $j \rightarrow \infty$  in  $H^s(\mathbb{R}^N)$ , then  $F_{\psi^j} \rightarrow F_\psi$  as  $j \rightarrow \infty$ .*

*Proof.* Note that by assumption,  $\|\psi^j\|_{H^s}$  is uniformly bounded, i.e., there is a constant  $M > 0$  independent of  $j$  such that  $\|\psi^j\|_{H^s} \leq M$  for all  $j$ . For any  $\delta > 0$ , define

$$F^\delta(x) = \begin{cases} F(x) & \text{if } |F(x)| \leq \frac{1}{\delta}, \\ 0 & \text{if } |F(x)| > \frac{1}{\delta}. \end{cases}$$

First we claim that  $F - F^\delta \in L^{\frac{2^*}{2^*-q}}(\mathbb{R}^N)$ . Indeed, let  $\Omega = \{x : |F(x)| > \frac{1}{\delta}\}$ , by assumption above we know  $|\Omega| < \infty$ . Writing  $F(x) = f_1(x) + f_2(x)$  with  $f_1 \in L^{\frac{2^*}{2^*-q}}(\mathbb{R}^N)$  and  $f_2(x) \in L^\infty(\mathbb{R}^N)$ , then we have

$$F - F^\delta = \chi_\Omega F(x) = \chi_\Omega f_1(x) + \chi_\Omega f_2(x),$$

where  $\chi_\Omega$  be the characteristic function on  $\Omega$ . Since  $|\Omega| < \infty$ , by Hölder inequality, we have  $\chi_\Omega f_2(x) \in L^{\frac{2^*}{2^*-q}}(\Omega)$ . It follows that  $\chi_\Omega f_2(x) \in L^{\frac{2^*}{2^*-q}}(\mathbb{R}^N)$ , thus the claim holds.

Moreover,  $F - F^\delta \rightarrow 0$  in  $L^{\frac{2^*}{2^*-q}}(\mathbb{R}^N)$  as  $\delta \rightarrow 0$  (by dominated convergence). Since  $\|\psi^j\|_{H^s} \leq M$ , by Sobolev inequality (Lemma 2.2)

$$\|\psi^j\|_{L^{2^*}} \leq S_{N,s} \|(-\Delta)^{s/2} \psi^j\|_{L^2} \leq S_{N,s} \|\psi^j\|_{H^s} \leq C_1. \tag{4.1}$$

By Hölder inequality, we have

$$\int (F - F^\delta)|\psi^j|^q \leq \|F - F^\delta\|_{\frac{2^*}{2^*-q}} \|\psi^j\|_{2^*}^q \leq C_1 \|F - F^\delta\|_{\frac{2^*}{2^*-q}} = C_\delta,$$

with  $C_\delta$  independent of  $j$ , moreover,  $C_\delta \rightarrow 0$  as  $\delta \rightarrow 0$ . Thus, our goal of showing that  $F_{\psi^j} \rightarrow F_\psi$  as  $j \rightarrow \infty$  would be achieved if we can prove that  $F_{\psi^j}^\delta \rightarrow F_\psi^\delta$  as  $j \rightarrow \infty$  for each  $\delta > 0$ .

To prove that  $F_{\psi^j}^\delta \rightarrow F_\psi^\delta$  as  $j \rightarrow \infty$ , now fix  $\delta$  and define the set

$$A_\epsilon = \{x : |F^\delta(x)| > \epsilon\}$$

for  $\epsilon > 0$ . By assumption,  $|A_\epsilon| < \infty$ . Then

$$F_{\psi^j}^\delta = \int_{A_\epsilon} F^\delta |\psi^j|^q + \int_{A_\epsilon^c} F^\delta |\psi^j|^q. \tag{4.2}$$

Since  $2 \leq q < 2^*$ ,  $\|\psi^j\|_2 \leq \|\psi^j\|_{H^s} \leq M$ , by Lemma 2.3, we have

$$\|\psi^j\|_{L^q} \leq C \|\psi^j\|_{H^s}^{\frac{(k-2)N}{2sq}} \|\psi^j\|_{L^2}^{1 - \frac{(q-2)N}{2sq}} \leq CM, \tag{4.3}$$

by weak lower semicontinuity of the norm, we also have

$$\|\psi\|_{L^q} \leq \liminf_{j \rightarrow \infty} \|\psi^j\|_{L^q} \leq CM. \tag{4.4}$$

Then

$$\int_{A_\epsilon^c} F^\delta |\psi^j|^q \leq \epsilon \int_{\mathbb{R}^N} |\psi^j|^q \leq \epsilon (CM)^q,$$

i.e., the last term of (4.2) tend to zero as  $\epsilon \rightarrow 0$ , and hence it suffices to show that the first term of (4.2) converges to  $\int_{A_\epsilon} F^\delta |\psi|^q$ .

This is accomplished as follows. By Theorem 3.2 (in the Appendix below), on any finite measure set (that we take to be  $A_\epsilon$ )  $\psi^j \rightarrow \psi$  strongly in  $L^r(A_\epsilon)$ , for  $r \in [1, 2^*)$ . Here we can choose  $r = q$ . Since  $q \geq 2$ , using the inequality

$$||\psi^j|^q - |\psi|^q| \leq C_q (|\psi^j|^{q-1} + |\psi|^{q-1}) |\psi^j - \psi|,$$

where  $C_q$  is a constant only dependent of  $q$ , and by (4.3), (4.4) and Hölder inequality, we have

$$\begin{aligned} \int_{A_\epsilon} ||\psi^j|^q - |\psi|^q| &\leq \int_{A_\epsilon} C_q (|\psi^j|^{q-1} + |\psi|^{q-1}) |\psi^j - \psi| \\ &\leq C_q ||\psi^j|^{q-1} + |\psi|^{q-1}||_{L^{\frac{q}{q-1}}(A_\epsilon)} ||\psi^j - \psi||_{L^q(A_\epsilon)} \\ &\leq C_3 ||\psi^j - \psi||_{L^q(A_\epsilon)}, \end{aligned}$$

so  $|\psi^j|^q \rightarrow |\psi|^q$  strongly in  $L^1(A_\epsilon)$ . Since  $F^\delta \in L^\infty(\mathbb{R}^N)$  (see the definition above), we conclude that

$$\int_{A_\epsilon} F^\delta |\psi^j|^q \rightarrow \int_{A_\epsilon} F^\delta |\psi|^q, \quad \text{as } j \rightarrow \infty.$$

This completes the proof. □

### 5. PROOF OF THEOREM 1.1

We will give the proof by a series of lemmas. Firstly, for any  $\beta > 0$ , we set

$$\Sigma_\beta := \{u \in H^s(\mathbb{R}^N) : \int_{\mathbb{R}^N} K(x)|u|^p dx = \beta\}.$$

**Lemma 5.1.** *Assume that  $K(x)$  satisfies (A3), then  $\Sigma_\beta$  is not empty.*

*Proof.* Since  $K(x) \geq 0$  and  $K(x) \not\equiv 0$ , for any fixed  $u \in H^s(\mathbb{R}^N) \setminus \{0\}$ , we have

$$\int_{\mathbb{R}^N} K(x)|u|^p dx > 0.$$

Write  $K(x) = K_1 + K_2$  with  $K_1 \in L^{\frac{2^*}{2^*-p}}(\mathbb{R}^N)$  and  $K_2 \in L^\infty(\mathbb{R}^N)$ . For any fixed  $u \in H^s(\mathbb{R}^N) \setminus \{0\}$ , Since  $2 < p < 2^*$ , by Hölder inequality, Lemma 2.2, Lemma 2.3 we have

$$\begin{aligned} \int_{\mathbb{R}^N} K(x)|u|^p dx &\leq \|K_1\|_{L^{\frac{2^*}{2^*-p}}} \|u\|_{L^{2^*}}^p + \|K_2\|_{L^\infty} \|u\|_{L^p}^p \\ &\leq C_1 \|(-\Delta)^{s/2} u\|_{L^2}^p + C_2 \|u\|_{H^s}^{\frac{(p-2)N}{2s}} \|u\|_{L^2}^{p - \frac{(p-2)N}{2s}} < \infty, \end{aligned}$$

where  $C_1$  and  $C_2$  are some constants. Then we can choose  $t > 0$  such that  $tu(x) \in \Sigma_\beta$ , where

$$t = \left( \frac{\beta}{\int_{\mathbb{R}^N} K(x)|u|^p dx} \right)^{1/p}.$$

□

Let  $\mathcal{I}(u)$  be the energy functional defined by (1.7), we want to consider the minimizing problem

$$\inf_{\Sigma_\beta} \mathcal{I}(u) = \frac{1}{2} \inf_{\Sigma_\beta} \left\{ \int_{\mathbb{R}^N} (|(-\Delta)^{s/2}u|^2 + (V(x) - \lambda)|u|^2) dx \right\} - \frac{1}{p}\beta.$$

Let

$$\mathcal{J}(u) = \int_{\mathbb{R}^N} (|(-\Delta)^{s/2}u|^2 + (V(x) - \lambda)|u|^2) dx, \tag{5.1}$$

with  $m_\beta = \inf_{u \in \Sigma_\beta} \mathcal{J}(u)$ , so we have

$$\inf_{u \in \Sigma_\beta} \mathcal{I}(u) = \frac{1}{2}m_\beta - \frac{1}{p}\beta.$$

Thus minimizing  $\mathcal{I}(u)$  on  $\Sigma_\beta$  is equivalent to considering just  $m_\beta$ .

**Lemma 5.2.** *With the assumptions of Theorem 1.1, let  $\{u_k\}_k \subset \Sigma_\beta$  be a minimizing sequence for  $m_\beta$ . Then  $\{u_k\}$  is bounded in  $H^s(\mathbb{R}^N)$ .*

*Proof.* Since  $\{u_k\}_k$  is a minimizer sequence for  $m_\beta$ , it follows that

$$\lim_{k \rightarrow \infty} \mathcal{J}(u_k) = m_\beta.$$

Then,  $\mathcal{J}(u_k)$  is bounded by a constant independent of  $k$ , i.e.,  $\mathcal{J}(u_k) \leq M$ . Since  $V(x) \in L^{\frac{N}{2s}}(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  by Theorem 2.4, we know that  $\lambda_0$  is finite. By the assumption  $\lambda < \lambda_0$  in Theorem 1.1, we have

$$\mathcal{J}(u_k) \geq \lambda_0 \int_{\mathbb{R}^N} |u_k|^2 dx - \lambda \int_{\mathbb{R}^N} |u_k|^2 dx = (\lambda_0 - \lambda) \|u_k\|_2^2,$$

it follows that  $\|u_k\|_2 \leq M/(\lambda_0 - \lambda)$ . i.e.,  $\{u_k\}_k$  is bounded in  $L^2(\mathbb{R}^N)$ .

Since  $\lambda \leq 0$ , by (ii) of Theorem 2.4, we have

$$\begin{aligned} \|(-\Delta)^{s/2}u_k\|_{L^2}^2 &\leq C\mathcal{E}(u_k) + D\|u_k\|_{L^2}^2 \\ &\leq C(\mathcal{E}(u_k) - \lambda\|u_k\|_{L^2}^2) + D\|u_k\|_{L^2}^2 \\ &= C\mathcal{J}(u_k) + D\|u_k\|_{L^2}^2 \\ &\leq CM + DM/(\lambda_0 - \lambda). \end{aligned}$$

Therefore,  $\{u_k\}_k$  is bounded in  $H^s(\mathbb{R}^N)$ . □

**Lemma 5.3.** *With the assumptions of Theorem 1.1, for every  $\beta > 0$ ,  $m_\beta$  is attained by a nonnegative function, namely there exists  $u_0 \in \Sigma_\beta$ ,  $u_0(x) \geq 0$  a.e. in  $\mathbb{R}^N$ , such that*

$$m_\beta = \mathcal{J}(u_0).$$

Moreover,  $m_\beta > 0$ .

*Proof.* Let  $\{u_k\}_k \subset \Sigma_\beta$  be a minimizing sequence for  $m_\beta$ . In Section 2, we know that  $\|(-\Delta)^{s/2}u_k\|_{L^2}^2$  is equivalent to

$$\int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_k(x) - u_k(y)|^2}{|x - y|^{2s+N}} dx dy,$$

it follows that  $\|(-\Delta)^{s/2}|u_k|\|_{L^2}^2 \leq \|(-\Delta)^{s/2}u_k\|_{L^2}^2$ , hence the sequence  $\{|u_k|\}_k$  is still a minimizing sequence and we can assume from the beginning that  $u_k \geq 0$  a.e. in  $\mathbb{R}^N$  for all  $k$ . By Lemma 5.2, this minimizing sequence is bounded in  $H^s(\mathbb{R}^N)$ , so up to subsequences,

$$u_k \rightharpoonup u_0 \quad \text{in } H^s(\mathbb{R}^N),$$

by Lemma 4.1, we have

$$\int_{\mathbb{R}^N} K(x)|u_k|^p dx \rightarrow \int_{\mathbb{R}^N} K(x)|u_0|^p dx,$$

and then

$$\int_{\mathbb{R}^N} K(x)|u_0|^p dx = \int_{\mathbb{R}^N} K(x)|u_k|^p dx = \beta,$$

thus  $u_0 \in \Sigma_\beta$ .

Applying Lemma 4.1 for  $p = 2$ , we have

$$\int_{\mathbb{R}^N} V(x)|u_k|^2 dx \rightarrow \int_{\mathbb{R}^N} V(x)|u_0|^2 dx.$$

Since  $\lambda \leq 0$ , by weak lower semicontinuity of the norm, it follows that

$$\begin{aligned} & \int_{\mathbb{R}^N} |(-\Delta)^{s/2}u_0|^2 dx + \int_{\mathbb{R}^N} V(x)|u_0|^2 dx - \lambda \int_{\mathbb{R}^N} |u_0|^2 dx \\ &= \|(-\Delta)^{s/2}u_0\|_{L^2}^2 + (-\lambda)\|u_0\|_{L^2}^2 + \int_{\mathbb{R}^N} V(x)|u_0|^2 dx \\ &\leq \liminf_{k \rightarrow \infty} [\|(-\Delta)^{s/2}u_k\|_{L^2}^2 + (-\lambda)\|u_k\|_{L^2}^2 + \int_{\mathbb{R}^N} V(x)|u_k|^2 dx] = m_\beta, \end{aligned}$$

together with  $u_0 \in \Sigma_\beta$ , this shows that

$$m_\beta = \int_{\mathbb{R}^N} |(-\Delta)^{s/2}u_0|^2 dx + \int_{\mathbb{R}^N} V(x)|u_0|^2 dx - \lambda \int_{\mathbb{R}^N} |u_0|^2 dx = \mathcal{J}(u_0).$$

Note that  $u_0 \in \Sigma_\beta$  implies that  $u_0 \not\equiv 0$ , then from the definition of  $\lambda_0$  given in Theorem 1.1 it follows that

$$m_\beta = \mathcal{J}(u_0) \geq (\lambda_0 - \lambda)\|u_0\|_{L^2}^2 > 0.$$

This completes the proof. □

**Lemma 5.4.** *With the assumptions of Theorem 1.1, let  $u_0$  be a minimizer for  $m_\beta$ . Then  $u_0$  satisfies*

$$\begin{aligned} & \int_{\mathbb{R}^N} (-\Delta)^{s/2}\overline{u_0} \cdot (-\Delta)^{s/2}v dx + \int_{\mathbb{R}^N} (V(x) - \lambda)\overline{u_0} \cdot v dx \\ &= \frac{m_\beta}{\beta} \int_{\mathbb{R}^N} K(x)|u_0|^{p-2}\overline{u_0} \cdot v dx \end{aligned} \tag{5.2}$$

for all  $v \in H^s(\mathbb{R}^N)$ .

*Proof.* Let  $\mathcal{J}(u_0)$  be energy functional defined by (5.1). Fix  $v(x) \in H^s(\mathbb{R}^N)$ , for  $\varepsilon \in \mathbb{R}$  small enough, when  $r \in (-\varepsilon, \varepsilon)$ , the function  $u_0 + rv$  is not identically zero. Therefore there exists a function  $t(r) : (-\varepsilon, \varepsilon) \rightarrow (0, \infty)$  such that

$$\int_{\mathbb{R}^N} |K(x)t(r)(u_0 + rv)|^p dx = \beta.$$

Precisely,

$$t(r) = \left( \frac{\beta}{\int_{\mathbb{R}^N} |K(x)(u_0 + rv)|^p dx} \right)^{1/p}.$$

Note that the map  $r \mapsto t(r)(u_0 + rv)$  defines a curve on  $\Sigma_\beta$  that passes through  $u_0$  when  $r = 0$ . The function  $t(r)$  is differentiable on  $(-\varepsilon, \varepsilon)$ ,

$$t'(r) = -\beta^{1/p} \left( \int_{\mathbb{R}^N} |K(x)(u_0 + rv)|^p dx \right)^{-\frac{1}{p}-1} \operatorname{Re}(K(x)|u_0 + rv|^{p-2}(u_0 + rv), v),$$

where  $\operatorname{Re}$  denotes real part of inner product  $(\cdot, \cdot)$  (defined in (1.9)), and

$$\operatorname{Re}(|u_0 + rv|^{p-2}(u_0 + rv), v) = \operatorname{Re} \int_{\mathbb{R}^N} K(x) |u_0 + rv|^{p-2} \overline{(u_0 + rv)} \cdot v \, dx.$$

Then we have

$$t(0) = 1 \quad \text{and} \quad t'(0) = -\beta^{-1} \operatorname{Re}(K(x)|u_0|^{p-2}u_0, v). \quad (5.3)$$

We define  $\gamma : (-\varepsilon, \varepsilon) \rightarrow \mathbb{R}$  as

$$\begin{aligned} \gamma(r) &= \mathcal{J}(t(r)(u_0 + rv)) = t^2(r)\mathcal{J}(u_0 + rv) \\ &= t^2(r)((-\Delta)^{s/2}(u_0 + rv), (-\Delta)^{s/2}(u_0 + rv)) \\ &\quad + t^2(r)((V(x) - \lambda)(u_0 + rv), (u_0 + rv)). \end{aligned}$$

Since  $t(r)(u_0 + rv) \in \Sigma_\beta$  for every  $r \in (-\varepsilon, \varepsilon)$ , the point  $r = 0$  is a local minimum for  $\gamma$ , such that

$$\gamma(0) = \mathcal{J}(u_0) = m_\beta. \quad (5.4)$$

The function  $\gamma$  is differentiable and

$$\begin{aligned} \gamma'(r) &= 2t(r)t'(r)\mathcal{J}(u_0 + rv) \\ &\quad + 2t^2(r) \operatorname{Re}[( (-\Delta)^{s/2}(u_0 + rv), (-\Delta)^{s/2}v) + ((V(x) - \lambda)(u_0 + rv), v)]. \end{aligned}$$

by (5.3), (5.4), then

$$\begin{aligned} 0 = \gamma'(0) &= 2t(0)t'(0)\mathcal{J}(u_0) + 2t^2(0) \operatorname{Re} [((-\Delta)^{s/2}u_0, (-\Delta)^{s/2}v) \\ &\quad + ((V(x) - \lambda)u_0, v)] \\ &= -2\beta^{-1} \operatorname{Re}(K(x)|u_0|^{p-2}u_0, v)m_\beta \\ &\quad + 2 \operatorname{Re} [((-\Delta)^{s/2}u_0, (-\Delta)^{s/2}v) + ((V(x) - \lambda)u_0, v)]. \end{aligned} \quad (5.5)$$

Since  $v$  is an arbitrary complex function in  $H^s(\mathbb{R}^N)$ , it follows that

$$-\beta^{-1}(K(x)|u_0|^{p-2}u_0, v)m_\beta + [((-\Delta)^{s/2}u_0, (-\Delta)^{s/2}v) + ((V(x) - \lambda)u_0, v)] = 0,$$

i.e. (5.2) holds.  $\square$

Let  $u_0$  be a minimizer for  $m_\beta$ . Set  $u_0(x) = c w(x)$ , where  $c \in \mathbb{R}$  will be determined later. By Lemma 5.4,  $w(x)$  satisfies

$$c[(-\Delta)^{s/2}w, (-\Delta)^{s/2}v] + ((V(x) - \lambda)w, v) = \frac{m_\beta}{\beta} c^{p-1}(K(x)|w|^{p-2}w, v)$$

for all  $v \in H^s(\mathbb{R}^N)$ . Choosing  $c = (\frac{\beta}{m_\beta})^{\frac{1}{p-2}}$ , we see that  $w(x)$  is nonnegative by Lemma 5.3 and satisfies

$$((-\Delta)^{s/2}w, (-\Delta)^{s/2}v) + ((V(x) - \lambda)w, v) = (K(x)|w|^{p-2}w, v) \quad \forall v \in H^s(\mathbb{R}^N),$$

namely  $w(x)$  is a weak (nonzero) solution of (1.4), such that

$$u_0(x) = \left(\frac{\beta}{m_\beta}\right)^{\frac{1}{p-2}} w(x). \quad (5.6)$$

Thus we obtain the existence of the solution.

Let  $\mathcal{N}$  be the Nehari manifold defined by (1.8), note that  $w(x) \in \mathcal{N}$ . We mention in Section 1 that a ground state of (1.4) is a solution that minimizes the energy

functional  $\mathcal{I}(u)$  on the Nehari manifold  $\mathcal{N}$ , next we will prove that  $w(x)$  is a ground state, that is, we need to prove that

$$\mathcal{I}(w) \leq \mathcal{I}(\phi), \quad \text{for any } \phi \in \mathcal{N}. \quad (5.7)$$

For any function  $\phi \in \mathcal{N}$ , then by the definition of  $\mathcal{N}$  we have

$$\mathcal{I}(\phi) = \left(\frac{1}{2} - \frac{1}{p}\right)\mathcal{J}(\phi), \quad (5.8)$$

where  $\mathcal{J}(\phi)$  is energy functional defined in (5.1).

Fix any  $\phi \in \mathcal{N}$  and let  $\theta := \int_{\mathbb{R}^N} K(x)|\phi|^p dx$ , then  $\phi \in \Sigma_\theta$ . Let  $v_0 = \tilde{c}w(x)$  with  $\tilde{c} = \left(\frac{\theta}{m_\theta}\right)^{\frac{1}{p-2}}$ , we claim that  $v_0$  is a minimizer for  $m_\theta$ . Indeed, for any  $u \in \Sigma_\beta$ , the scaling  $v = \left(\frac{\theta}{\beta}\right)^{1/p}u \in \Sigma_\theta$ , then  $\mathcal{J}(v) = \left(\frac{\theta}{\beta}\right)^{2/p}\mathcal{J}(u)$ , it follows that

$$\frac{m_\beta}{\beta^{2/p}} = \frac{m_\theta}{\theta^{2/p}}, \quad \text{for any } \theta > 0 \text{ such that } \theta \neq \beta. \quad (5.9)$$

Note that, by (5.6) we know  $w(x) = \left(\frac{m_\beta}{\beta}\right)^{\frac{1}{p-2}}u_0$ , then by (5.9) we have

$$v_0 = \tilde{c}w(x) = \left(\frac{\theta}{m_\theta}\right)^{\frac{1}{p-2}}\left(\frac{m_\beta}{\beta}\right)^{\frac{1}{p-2}}u_0 = \left(\frac{\theta}{\beta}\right)^{1/p}u_0. \quad (5.10)$$

Since  $u_0$  is the minimizer for  $m_\beta$ , it follows that  $u_0 \in \Sigma_\beta$ , and that that  $v_0 \in \Sigma_\theta$ . Moreover, using (5.9) again,

$$\mathcal{J}(v_0) = \left(\frac{\theta}{\beta}\right)^{2/p}\mathcal{J}(u_0) = \left(\frac{\theta}{\beta}\right)^{2/p}m_\beta = m_\theta,$$

thus  $v_0$  is the minimizer for  $m_\theta$ .

Since  $w \in \mathcal{N}$ ,  $v_0, \phi \in \Sigma_\theta$  and  $v_0$  is the minimizer for  $m_\theta$ , by (5.8) we have

$$\begin{aligned} \mathcal{I}(w) &= \left(\frac{1}{2} - \frac{1}{p}\right)\mathcal{J}(w) = \left(\frac{1}{2} - \frac{1}{p}\right)\mathcal{J}(\tilde{c}^{-1}v_0) \\ &= \left(\frac{1}{2} - \frac{1}{p}\right)\tilde{c}^{-2}\mathcal{J}(v_0) = \left(\frac{1}{2} - \frac{1}{p}\right)\left(\frac{m_\theta}{\theta}\right)^{\frac{2}{p-2}}\mathcal{J}(v_0) \\ &\leq \left(\frac{1}{2} - \frac{1}{p}\right)\left(\frac{m_\theta}{\theta}\right)^{\frac{2}{p-2}}\mathcal{J}(\phi) = \left(\frac{m_\theta}{\theta}\right)^{\frac{2}{p-2}}\mathcal{I}(\phi), \end{aligned}$$

hence to prove  $\mathcal{I}(w) \leq \mathcal{I}(\phi)$ , it is sufficient to show that  $\frac{m_\theta}{\theta} \leq 1$ . Since  $\phi \in \mathcal{N} \cap \Sigma_\theta$ , we obtain

$$\mathcal{J}(\phi) = \int_{\mathbb{R}^N} |(-\Delta)^{s/2}\phi|^2 dx + \int_{\mathbb{R}^N} (V(x) - \lambda)|\phi|^2 dx = \int_{\mathbb{R}^N} K(x)|\phi|^p dx = \theta.$$

Thus

$$m_\theta = \inf_{u \in \Sigma_\theta} \mathcal{J}(u) \leq \mathcal{J}(\phi) = \theta,$$

i.e.,  $\frac{m_\theta}{\theta} \leq 1$ . Thus  $w(x)$  is a ground state of (1.4). This completes the proof of Theorem 1.1.

## 6. PROOF OF THEOREMS 1.2 AND 1.3

In this section we prove that weak solutions of (1.4) are of class  $C^{0,\alpha}(\mathbb{R}^N)$  for some  $\alpha \in (0, 1)$ . First we give some properties of  $L^q(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  which will be used below.

**Proposition 6.1.** *The space  $L^q(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  has following properties.*

- (i)  $L^r(\mathbb{R}^N) \subset L^q(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  for any  $1 \leq q \leq r \leq \infty$ .

(ii)  $L^r(\mathbb{R}^N) + L^\infty(\mathbb{R}^N) \subset L^q(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  for any  $1 \leq q \leq r \leq \infty$ .

*Proof.* (i) Let  $f(x) \in L^r(\mathbb{R}^N)$ , for a given constant  $M > 0$  we have  $f(x) = f_0 + f_1$ , where

$$f_0 = \chi_{\{|f(x)| > M\}} f(x), \quad f_1 = \chi_{\{|f(x)| \leq M\}} f(x).$$

by the Chebyshev inequality [5]

$$|\{x : |f(x)| > M\}| \leq \left(\frac{C\|f\|_{L^r}}{M}\right)^r < \infty.$$

Since  $q \leq r$ , then  $L^r(\{x : |f(x)| > M\}) \subset L^q(\{x : |f(x)| > M\})$ , then  $f_0 \in L^q(\mathbb{R}^N)$ . It is obvious to see that  $f_1 \in L^\infty(\mathbb{R}^N)$ . Then  $f(x) \in L^q(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$ . Therefore, the case (i) holds.

The case (ii) is easy to obtain from case (i). □

Recall that the definition of fractional Sobolev spaces (e.g. see [6]) for  $p \geq 1$  and  $\beta > 0$ :

$$\mathcal{L}^{\beta,p} = \{u \in L^p(\mathbb{R}^N) | \mathcal{F}^{-1}[(1 + |\xi|^2)^{\beta/2} \hat{u}] \in L^p(\mathbb{R}^N)\},$$

and associated to the fractional Laplacian, the space

$$\mathcal{W}^{\beta,p} = \{u \in L^p(\mathbb{R}^N) | \mathcal{F}^{-1}[(1 + |\xi|^\beta) \hat{u}] \in L^p(\mathbb{R}^N)\}.$$

The following two theorems are basic results for these spaces which can be found in [6].

**Theorem 6.2** ([6]). *Assume that  $p \geq 1$  and  $\beta > 0$ . The following hold:*

- (i)  $\mathcal{L}^{\beta,p} = \mathcal{W}^{\beta,p}$ , and  $\mathcal{L}^{n,p} = W^{n,p}(\mathbb{R}^N)$  for all  $n \in \mathbb{N}$ , where  $W^{n,p}$  is the usual Sobolev space.
- (ii) For  $\alpha \in (0, 1)$  and  $2\alpha < \beta$ , we have  $(-\Delta)^\alpha : W^{\beta,p} \rightarrow W^{\beta-2\alpha,p}$ .
- (iii) For  $\alpha, \gamma \in (0, 1)$  and  $0 < \mu \leq \gamma - 2\alpha$ , we have

$$(-\Delta)^\alpha : C^{0,\gamma}(\mathbb{R}^N) \rightarrow C^{0,\mu}(\mathbb{R}^N) \quad \text{if } 2\alpha < \gamma,$$

and, for  $0 \leq \mu \leq 1 + \gamma - 2\alpha$ ,

$$(-\Delta)^\alpha : C^{1,\gamma}(\mathbb{R}^N) \rightarrow C^{0,\mu}(\mathbb{R}^N) \quad \text{if } 2\alpha > \gamma.$$

**Theorem 6.3** ([6]). (i) *If  $0 \leq \alpha$ , and either  $1 < p \leq q \leq Np/(N - \alpha p) < \infty$  or  $p = 1$  and  $1 \leq q < N/(N - \alpha)$ , then  $\mathcal{L}^{\alpha,p}$  is continuously embedded in  $L^q(\mathbb{R}^N)$ .*

- (ii) *Assume that  $0 \leq \alpha \leq 2$  and  $\alpha > N/p$ . If  $\alpha - N/p > 1$  and  $0 < \mu \leq \alpha - N/p - 1$ , then  $\mathcal{L}^{\alpha,p}$  is continuously embedded in  $C^{1,\mu}(\mathbb{R}^N)$ . If  $\alpha - N/p < 1$  and  $0 < \mu \leq \alpha - N/p$ , then  $\mathcal{L}^{\alpha,p}$  is continuously embedded in  $C^{0,\mu}(\mathbb{R}^N)$ .*

Let  $\mathcal{H}(x, t)$  be defined in (3.1) (in the Appendix below), then we define the kernel  $\mathcal{K}, \mathcal{K}^\mu$  with  $\mu > 0$  as

$$\mathcal{K}(x) = \int_0^\infty e^{-t} \mathcal{H}(x, t) dt, \quad \mathcal{K}^\mu(x) = \int_0^\infty e^{-\mu t} \mathcal{H}(x, t) dt. \tag{6.1}$$

By the rescaling property of  $\mathcal{H}(x, t)$ ,

$$\mathcal{H}\left(x, \frac{t}{\mu}\right) = \mu^{\frac{N}{2s}} \mathcal{H}\left(\mu^{\frac{1}{2s}} x, t\right),$$

we have

$$\mathcal{K}^\mu(x) = \mu^{\frac{N}{2s} - 1} \mathcal{K}\left(\mu^{\frac{1}{2s}} x\right). \tag{6.2}$$

On the other hand, In the Appendix of [6], we know that  $\mathcal{K}(x) = \mathcal{F}^{-1}\left(\frac{1}{1+|\xi|^{2s}}\right)$ , then in the same way, we have

$$\mathcal{K}^\mu(x) = \mathcal{F}^{-1}\left(\frac{1}{\mu + |\xi|^{2s}}\right). \tag{6.3}$$

The following theorem can be found in [6].

**Theorem 6.4** ([6]). *Let  $N \geq 2$  and  $s \in (0, 1)$ . Then we have the following:*

- (i)  $\mathcal{K}$  is positive, radically symmetric and smooth in  $\mathbb{R}^N \setminus \{0\}$ . Moreover, it is nonincreasing as a function of  $r = |x|$ .
- (ii) For appropriate constants  $C_1$  and  $C_2$ ,

$$\begin{aligned} \mathcal{K}(x) &\leq \frac{C_1}{|x|^{N+2s}} \quad \text{if } |x| \geq 1, \\ \mathcal{K}(x) &\leq \frac{C_2}{|x|^{N-2s}} \quad \text{if } |x| \leq 1. \end{aligned} \tag{6.4}$$

**Corollary 6.5.** *For  $N \geq 2$  and  $s \in (0, 1)$ , we have  $\mu > 0$  and  $\mathcal{K}^\mu$  satisfies Theorem 6.4 (i)-(ii).*

Since (6.2) holds, then it is easy to verify the above corollary.

*Proof of Theorem 1.2.* Since  $V(x)$  is bound from above, then there exists a constant  $M > 0$  such that  $V(x) \leq M$ . Note that  $u(x) \in H^s(\mathbb{R}^N)$  is a nonnegative solution of (1.4) satisfying

$$(-\Delta)^s u(x) + V(x)u(x) - K(x)|u|^{p-2}u(x) = \lambda u(x),$$

then

$$(-\Delta)^s u(x) + (M - \lambda)u(x) = (M - V(x))u(x) + K(x)|u|^{p-2}u(x).$$

Let  $\mu_0 = M - \lambda$ , since  $\lambda \leq 0$ , we have  $\mu_0 > 0$ . Let  $h(x) = (M - V(x))u(x) + K(x)|u|^{p-2}u(x)$ , then we have

$$u(x) = \mathcal{K}^{\mu_0} * h(x).$$

Note that  $u(x)$  is nonnegative and nontrivial,  $V(x) \leq M$ ,  $K(x) \neq 0$ , we have  $h(x) \geq 0$  such that  $h(x) \neq 0$ . By the corollary 6.5, we know that  $\mathcal{K}^{\mu_0}$  is positive, it follows that  $u(x)$  is positive in  $\mathbb{R}^N$ . The proof is complete.  $\square$

To discuss the regularity of the weak solution (1.4), first we discuss the following result about liner equations.

**Theorem 6.6.** *Let  $s \in (0, 1)$ , assume that  $u \in H^s(\mathbb{R}^N)$ ,  $N > 2s$  such that*

$$(-\Delta)^s u(x) + \mu u(x) = V(x)u(x) \quad \text{in } \mathbb{R}^N, \tag{6.5}$$

*for  $\mu > 0$ ,  $V(x) \in L^q(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  with  $q > \frac{N}{2s}$ . Then  $u \in C^{0,\alpha}(\mathbb{R}^N)$  for some  $\alpha \in (0, 1)$ . Moreover,  $u(x) \rightarrow 0$  as  $|x| \rightarrow \infty$ .*

*Proof.* First we know that  $u \in H^s(\mathbb{R}^N) = \mathcal{W}^{s,2}$ . Let  $1 = r_0 > r_1 > r_2 > \dots$ , and consider  $B_i = B(0, r_i)$ , the ball of radius  $r_i$  and centered at the origin. We define  $h(x) = V(x)u(x)$ , since  $V(x) \in L^q(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$ , we have  $V(x) = V_1 + V_2$  such that  $V_1 \in L^q(\mathbb{R}^N)$  and  $V_2 \in L^\infty(\mathbb{R}^N)$ , then  $h(x) = h_1 + h_2$  with  $h_1 = V_1 u(x)$  and  $h_2 = V_2 u(x)$ . Since  $u \in H^s(\mathbb{R}^N)$ , by Sobolev inequality we have  $u \in L^{2^*}(\mathbb{R}^N)$  with  $2^* = 2N/(N - 2s)$ . Since  $V_1 \in L^q(\mathbb{R}^N)$ , by Hölder inequality, then we have  $h_1 \in$

$L^{k_0}(\mathbb{R}^N)$  with  $k_0 = (1/q + 1/2^*)^{-1}$ . Therefore,  $h(x) = h_1 + h_2$  with  $h_1 \in L^{k_0}(\mathbb{R}^N)$  and  $h_2 \in L^{2^*}(\mathbb{R}^N)$ .

Now let  $\eta_1 \in C^\infty$  with  $0 \leq \eta_1 \leq 1$ , with support in  $B_0$  and such that  $\eta_1 \equiv 1$  in  $B_{1/2}$ , where  $B_{1/2} = B(0, r_{1/2})$  with  $r_1 < r_{1/2} < r_0$ . Let  $u_1$  be the solution of the equation

$$(-\Delta)^s u_1 + \mu u_1 = \eta_1 h(x) \quad \text{in } \mathbb{R}^N, \tag{6.6}$$

then

$$(-\Delta)^s(u - u_1) + \mu(u - u_1) = (1 - \eta_1)h(x) \quad \text{in } \mathbb{R}^N, \tag{6.7}$$

so that

$$u - u_1 = \mathcal{K}^\mu * \{(1 - \eta_1)h\}. \tag{6.8}$$

Using the Hölder inequality and (6.4) we have

$$\begin{aligned} &|u(x) - u_1(x)| \\ &\leq C \{ \|\mathcal{K}^\mu\|_{L^{l_0}(B_{1/2}^c)} \|(1 - \eta_1)h_1\|_{L^{k_0}} + \|\mathcal{K}^\mu\|_{L^{l_1}(B_{1/2}^c)} \|(1 - \eta_1)h_2\|_{L^{2^*}} \}, \end{aligned} \tag{6.9}$$

for all  $x \in B_1$ , where  $l_0 = k_0/(k_0 - 1)$ ,  $k_0$  is given above, and  $l_1 = 2^*/(2^* - 1)$ . In view of this inequality we have to concentrate our attention in  $u_1(x)$ .

Since  $B_0$  is bound and  $V(x) \in L^q(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$ , we obtain that  $\eta_1 V(x) \in L^q(B_0)$ . With the assumption  $q > \frac{N}{2s}$ , we have  $\eta_1 V(x) \in L^{q_0}(B_0)$  for  $\frac{N}{2s} < q_0 \leq \min\{q, \frac{N}{s}\}$ . Since  $u \in L^{2^*}(\mathbb{R}^N)$ , by Hölder inequality, we have  $\eta_1 V(x)u \in L^{k_1}(\mathbb{R}^N)$ , for  $k_1 = (1/q_0 + 1/2^*)^{-1}$  such that  $k_1 > 1$ . Since  $\eta_1$  has support in  $B_0$ , we have  $\eta_1 V(x)u \in L^{p_1}(\mathbb{R}^N)$ , for any  $1 < p_1 < \min\{k_1, N/(2s)\}$ . Note that  $u_1$  satisfies (6.6), thus by the definition of the space  $\mathcal{W}^{2s, p_1}$ , we have  $u_1 \in \mathcal{W}^{2s, p_1}$ . Then, using Sobolev embedding of the Theorem 6.3 (i) and (6.9), we have  $u \in L^{q_1}(B_1)$  for  $q_1 = p_1 N / (N - 2s p_1)$ .

Now we repeat the procedure, but consider a smooth function  $\eta_2$  such that  $0 \leq \eta_2 \leq 1$ , with support in  $B_1$  and  $\eta_2 \equiv 1$  in  $B_{3/2}$ , where  $B_{3/2} = B(0, r_{3/2})$  with  $r_2 < r_{3/2} < r_1$ . We also have  $\eta_2 V(x) \in L^{q_0}(B_1)$  for any  $\frac{N}{2s} < q_0 \leq \min\{q, \frac{N}{s}\}$ , we can set  $\frac{1}{q_0} = \frac{2s}{N} - \epsilon$  with  $0 < \epsilon \leq \frac{s}{N}$ . By Hölder inequality again, we have  $\eta_2 V(x)u \in L^{p_2}(B_2)$  for any

$$1 \leq p_2 < p_1 / (1 - \epsilon) \quad \text{where } p_2 = (1/q_0 + 1/q_1)^{-1}.$$

Proceeding as above, with the obvious changes we obtain that

$$u_2 = \mathcal{K}^\mu * (\eta_2 h(x)),$$

satisfying  $u_2 \in \mathcal{W}^{2s, p_2}$ . Then we have  $u \in L^{q_2}(B_2)$  for  $q_2 = p_2 N / (N - 2s p_2)$ .

Repeating the argument, for sequences  $\eta_j$ ,  $p_j$  and  $q_j = p_j N / (N - 2s p_j)$ , we have  $\eta_j V(x)u \in L^{p_j}(B_j)$  for any

$$1 \leq p_j < p_{j-1} / (1 - \epsilon) \quad \text{where } p_j = (1/q_0 + 1/q_j)^{-1}.$$

It follows that for some finite  $j$ ,  $\eta_j V(x)u \in L^{p_j}(B_j)$  such that  $p_j > N/(2s)$ . Then by Theorem 6.3(ii), we have  $u_j \in C^{0, \alpha}(\mathbb{R}^N)$  for some  $\alpha \in (0, 1)$ . Since  $u_j$  satisfies the inequality that similar to (6.9), we have  $u \in C^{0, \alpha}(B_{j+1})$ .

The ball  $B_j$  is centered at the origin, but we may arbitrarily move it around  $\mathbb{R}^N$ . Covering  $\mathbb{R}^N$  with these balls, we obtain that  $u \in C^{0, \alpha}(\mathbb{R}^N)$ . Finally, the fact that  $u \in L^{2^*}(\mathbb{R}^N) \cap C^{0, \alpha}(\mathbb{R}^N)$  implies that  $u(x) \rightarrow 0$  as  $|x| \rightarrow \infty$ , completing the proof.  $\square$

*Proof of Theorem 1.3.* Note that  $u(x)$  satisfies

$$(-\Delta)^s u(x) + V(x)u(x) - K(x)|u|^{p-2}u(x) = \lambda u(x),$$

for  $2 < p < 2^*$ . Let  $\tilde{V}(x) = -V(x) + K(x)|u|^{p-2}$ , then the equation becomes

$$(-\Delta)^s u(x) - \lambda u(x) = \tilde{V}(x)u(x).$$

We claim that  $\tilde{V}(x) \in L^l(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  for some  $l > \frac{N}{2s}$ .

Since the condition (A5) holds,  $K(x) \in L^{\tilde{r}}(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  for  $\tilde{r} > \frac{2^*}{2^*-p}$ , then  $K(x) = K_1 + K_2$  with  $K_1 \in L^{\tilde{r}}(\mathbb{R}^N)$  and  $K_2 \in L^\infty(\mathbb{R}^N)$ . Then Since  $u \in L^{2^*}(\mathbb{R}^N)$ , we have  $K_2|u|^{p-2} \in L^{r_0}(\mathbb{R}^N)$  for  $r_0 = \frac{2^*}{p-2} > \frac{2^*}{2^*-2} = \frac{N}{2s}$ . By Hölder inequality, we have  $K_1|u|^{p-2} \in L^{r_1}(\mathbb{R}^N)$  with  $r_1 = (\frac{1}{\tilde{r}} + \frac{p-2}{2^*})^{-1}$  such that  $r_1 > \frac{N}{2s}$ . Then by Proposition 6.1 (i), we have  $K(x)|u|^{p-2} \in L^{l_1}(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  with  $l_1 = \min\{r_0, r_1\}$  such that  $l_1 > \frac{N}{2s}$ . Then by Proposition 6.1 (ii), we have  $\tilde{V}(x) \in L^l(\mathbb{R}^N) + L^\infty(\mathbb{R}^N)$  with  $l = \min\{\tilde{q}, l_1\}$  (where  $\tilde{q}$  given in (A5)), such that  $l > \frac{N}{2s}$ . Then by Theorem 6.6, we obtain the regular result of Theorem 1.3.  $\square$

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