

**SUFFICIENT CONDITIONS FOR HADAMARD
WELL-POSEDNESS OF A COUPLED
THERMO-CHEMO-POROELASTIC SYSTEM**

TETYANA MALYSHEVA, LUTHER W. WHITE

ABSTRACT. This article addresses the well-posedness of a coupled parabolic-elliptic system modeling fully coupled thermal, chemical, hydraulic, and mechanical processes in porous formations that impact drilling and borehole stability. The underlying thermo-chemo-poroelastic model is a system of time-dependent parabolic equations describing thermal, solute, and fluid diffusions coupled with Navier-type elliptic equations that attempt to capture the elastic behavior of rock around a borehole. An existence and uniqueness theory for a corresponding initial-boundary value problem is an open problem in the field. We give sufficient conditions for the well-posedness in the sense of Hadamard of a weak solution to a fully coupled parabolic-elliptic initial-boundary value problem describing homogeneous and isotropic media.

1. INTRODUCTION

In this article, we are concerned with the well-posedness of a coupled parabolic-elliptic system arising in petroleum- and geothermal-related applications of rock mechanics. This work is motivated by the problems of drilling and borehole stability in porous formations that involve the modeling of fully coupled thermal, chemical, hydraulic, and mechanical (elastic deformation) processes.

1.1. Literature review. The poroelasticity theory describing the coupled processes of elastic deformation and pore fluid diffusion in fluid-saturated isothermal porous media can be traced back to the pioneering works of von Terzaghi [26, 27] and Biot [3]. Biot [4] pointed out a complete mathematical analogy between poroelasticity and thermoelasticity with the temperature playing the same role as the fluid pressure and heat conduction corresponding to fluid flow. A complete well-posedness analysis of a general initial-boundary value problem for a system of coupled partial differential equations that describes the Biot consolidation model [3] in poroelasticity, as well as a coupled quasi-static problem in thermoelasticity, has been carried out by Showalter [22, 23]. Based on the theory of linear degenerate evolution equations in Hilbert spaces, the existence and uniqueness of strong and

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weak solutions, as well as regularity theory for the initial-boundary value problem, were developed.

The Biot theory for isothermal systems was first extended by Schiffman [21] to account for the effects of thermal expansion of both the pore fluid and the elastic matrix. Since then, a substantial literature on thermo-poroelasticity theory and the modeling of the coupled hydro-thermo-mechanical behavior of a fluid-saturated porous media has been developed for rock mechanics, including petroleum and geothermal borehole stability [1, 2, 5, 7, 8, 9, 14, 15, 16, 18, 19, 28]. Considerable attention has been recently placed on the impact of chemical processes in porous media on drilling and borehole stability [6, 11, 20, 29]. Due to the complexity of cross-coupling mechanisms involved in thermo-poroelastic and chemo-thermo-poroelastic models, the general question of existence and uniqueness of solutions to the corresponding initial-boundary value problems remains open and very few analytical solutions are currently available. Typically, the solutions are derived under the assumptions that some of the couplings can be neglected [5, 11, 14, 15, 18, 20]. An exact unique analytical solution for a specific case of the fully coupled thermo-hydro-mechanical response of a fluid-saturated porous sphere under mechanical pulse load was developed by Belotserkovets and Prevost [2] using the Laplace transformation and the residue theorem.

The model that we study is based on equations derived by Diek [6] and constitutes the general theory of fully coupled chemical thermo-poroelasticity for porous media saturated by a compressible fluid. The theory satisfies the first and second laws of thermodynamics and is based on concepts of irreversible thermodynamics, a novel rock constitutive relation, and Onsager's transport phenomenology [6]. Our objective here is to establish the well-posedness theory for the fully coupled thermo-chemo-poroelastic (TCP) system describing homogeneous and isotropic fluid-saturated porous media. The main result of this work is Theorem 4.1 which gives a sufficient condition for the well-posedness in the sense of Hadamard of a weak solution to the coupled linear parabolic-elliptic initial-boundary value problem describing the fully coupled TCP model.

1.2. Notation. Let $\Omega \subset \mathbb{R}^n$, $n = 2, 3$, be a bounded open domain with a sufficiently smooth boundary Γ and $\bar{\Omega} = \Omega \cup \Gamma$. We write \mathbf{x} for $(x, y, z) \in \bar{\Omega} \subset \mathbb{R}^3$ or $(x, y) \in \bar{\Omega} \subset \mathbb{R}^2$. Let $\Gamma_F \subset \Gamma$ and Γ_F have a nonempty interior relative to Γ . The following notation will be used for the spaces of vector-valued functions: $\mathbb{H}^n = L^2(\Omega)^n$, $\mathbb{V}^n = H^1(\Omega)^n$, $\mathbb{V}_0^n = H_0^1(\Omega)^n$ with the norm

$$\|\mathbf{u}\|_{\mathbb{V}_0^n} = \left[\sum_{k=1}^n \int_{\Omega} |\nabla u_k|^2 d\Omega \right]^{1/2}, \quad \mathbf{u} = [u_1 \dots u_n]^T$$

and $\tilde{\mathbb{V}}_0^n = \{\boldsymbol{\varphi} \in \mathbb{V}^n : \boldsymbol{\varphi}|_{\Gamma_F} = \mathbf{0}\}$ with the norm inherited from \mathbb{V}^n or from \mathbb{V}_0^n , $n \in \mathbb{N}$, as above. We denote by $L^2(a, b; X)$ the space of L^2 -integrable functions from $[a, b] \subset \mathbb{R}$ into a Hilbert space X with the norm

$$\|u\|_{L^2(a, b; X)} = \left[\int_a^b \|u(t)\|_X^2 dt \right]^{1/2}$$

We also introduce the Hilbert space

$$\mathcal{W}(a, b; X) = \{u : u \in L^2(a, b; X), \dot{u} \in L^2(a, b; X')\}$$

with the norm

$$\|u\|_{\mathcal{W}} = \left[\|u\|_{L^2(a,b;X)}^2 + \|\dot{u}\|_{L^2(a,b;X')}^2 \right]^{1/2}$$

where the superscript dot $(\dot{\cdot})$ denotes a time derivative and X' is the dual of X .

Let $\mathbf{f} = [f_1 \dots f_n]^T \in \mathbb{V}^n$. We will use the notation $\nabla \mathbf{f} = [\nabla f_1 \dots \nabla f_n]^T$ and

$$(\nabla \mathbf{f}, \nabla \mathbf{g})_{\mathbb{H}^n} = \sum_{k=1}^n (\nabla f_k, \nabla g_k)_{\mathbb{H}^n}, \quad \mathbf{f}, \mathbf{g} \in \mathbb{V}^n$$

Throughout this article, the symbol c will be used to denote the Poincaré constant: $c = c(\Omega) > 0$, $\|\mathbf{u}\|_{\mathbb{H}^n}^2 \leq c\|\mathbf{u}\|_{\mathbb{V}_0^n}^2$, $\mathbf{u} \in \mathbb{V}_0^n$.

2. THE COUPLED TCP MODEL: UNDERLYING EQUATIONS

The underlying TCP model is a system of time-dependent parabolic partial differential equations coupled with Navier-type elliptic equations with time, $t \in (0, t_f)$, as a parameter. The parabolic equations represent heat, solute, and fluid diffusions, and the Navier-type elliptic equations attempt to capture the elastic behavior of rock, while incorporating thermal, chemical, and porous media effects. The equations are developed in terms of the following variables: the absolute temperature $T(\mathbf{x}, t)$, the solute mass fraction $C(\mathbf{x}, t)$, the pore pressure $p(\mathbf{x}, t)$, and the vector of solid displacements $\mathbf{u}(\mathbf{x}, t)$.

The coupled partial differential equations supplemented by the appropriate initial and boundary conditions constitute an initial-boundary-value problem with constant coefficients defined in an open region Ω exterior to the borehole. The region is specified with a sufficiently smooth boundary Γ such that, without loss of generality, we may assume that, on the outer (far-field) boundary Γ_F , the absolute temperature, solute mass fraction, pore pressure, and displacements are time-independent and displacements and their velocities are negligibly small. Specifically, the region is an elliptical annulus in a two-dimensional case and a vertical or inclined finite cylinder in a three-dimensional case.

It is convenient to consider the thermal diffusion, solute diffusion, and fluid diffusion as a system (diffusion system). Hence, we introduce the vector $\mathbf{V} = [T \ C \ p]^T$ and, with this notation, the initial-boundary value problem describing the underlying TCP model has the form

$$M\dot{\mathbf{V}} - A\nabla^2 \mathbf{V} = -\mathbf{b}_0(\nabla \cdot \dot{\mathbf{u}}), \quad \text{in } \Omega \times (0, t_f), \tag{2.1}$$

$$\left(K + \frac{G}{3}\right)\nabla(\nabla \cdot \dot{\mathbf{u}}) + G\nabla^2 \dot{\mathbf{u}} = \nabla(\mathbf{b}_1 \cdot \dot{\mathbf{V}}), \quad \text{in } \Omega \times (0, t_f) \tag{2.2}$$

with boundary conditions

$$\mathbf{V}(\mathbf{x}, t) = \mathbf{V}_B(\mathbf{x}, t), \quad \text{on } \Gamma \times (0, t_f) \tag{2.3}$$

$$\mathbf{u}(\mathbf{x}, t) \approx \mathbf{0}, \quad \text{on } \Gamma_F \times (0, t_f) \tag{2.4}$$

$$\dot{\mathbf{u}}(\mathbf{x}, t) \approx \mathbf{0}, \quad \text{on } \Gamma_F \times (0, t_f) \tag{2.5}$$

$$\dot{\mathbf{t}}\mathbf{n} = ((\mathbf{b}_1 \cdot \dot{\mathbf{V}})I + \dot{\boldsymbol{\sigma}})\mathbf{n}, \quad \text{on } \Gamma \setminus \Gamma_F \times (0, t_f) \tag{2.6}$$

and initial conditions

$$\mathbf{V}(\mathbf{x}, 0) = \begin{cases} \mathbf{V}_I(\mathbf{x}), & \text{in } \Omega \\ \mathbf{V}_B(\mathbf{x}, 0), & \text{on } \Gamma \end{cases} \tag{2.7}$$

where M and A are 3×3 matrices of diffusion coefficients, \mathbf{b}_0 and \mathbf{b}_1 are 3×1 coupling vectors determined by physical properties and input parameters of the rock/fluid system, K and G are the bulk and shear moduli, respectively, $\boldsymbol{\tau}$ is the stress tensor, \mathbf{n} is the outward unit normal vector on the boundary, I is the $n \times n$ identity matrix, $n = 2$ or 3 is the dimension of the problem, and $\hat{\boldsymbol{\sigma}}$ is the applied boundary stress tensor. The boundary function \mathbf{V}_B is spatially independent on the inner (borehole) boundary Γ_B and time-independent on the far-field boundary Γ_F .

Remark 2.1. The matrices $M = [m_{ij}]$ and $A = [a_{ij}]$, $i, j = 1, 2, 3$, of diffusion coefficients are non-symmetric with the entries specified as follows: $m_{21} = m_{23} = 0$ and no other entries of M are zero, and $a_{ij} \neq 0$, $i, j = 1, 2, 3$. The second entry in the coupling vector \mathbf{b}_0 is zero and all the entries of the coupling vector \mathbf{b}_1 are non-zero. This implies that the TCP model is non-symmetric, and the diffusion equation (2.1) cannot be rescaled to make the coupling vectors equal. Therefore, the methods presented in [23] are not applicable to the fully coupled TCP problem (2.1)-(2.7).

We begin with a lemma that provides an alternative formulation of the Navier-type elastic system (2.2), (2.4)-(2.6). Its proof, based on the principle of minimum total potential energy, will play an important role in well-posedness analysis presented in the next section.

Lemma 2.2. *The Navier-type elastic system (2.2), (2.4)-(2.6) is equivalent to the system*

$$\left(K + \frac{G}{3}\right)\nabla(\nabla \cdot \mathbf{u}) + G\nabla^2\mathbf{u} = \nabla(\mathbf{b}_1 \cdot \mathbf{V}), \quad \text{in } \Omega \times (0, t_f) \quad (2.8)$$

$$\mathbf{u} \approx \mathbf{0}, \quad \text{on } \Gamma_F \times (0, t_f) \quad (2.4)$$

$$\boldsymbol{\tau}\mathbf{n} = ((\mathbf{b}_1 \cdot \mathbf{V})I + \hat{\boldsymbol{\sigma}})\mathbf{n}, \quad \text{on } \Gamma \setminus \Gamma_F \times (0, t_f) \quad (2.9)$$

Proof. For the sake of simplicity, we restrict the proof to the two-dimensional case $\Omega \subset \mathbb{R}^2$. Directly analogous arguments can be developed for the three-dimensional case. By the principle of minimum total potential energy, the region Ω shall displace to a position that minimizes the total potential energy of the elastic system,

$$\mathcal{V}(\mathbf{u}) = \mathcal{V}_S(\mathbf{u}) - W_b(\mathbf{u}) - W_S(\mathbf{u}) \quad (2.10)$$

where $\mathcal{V}_S(\mathbf{u})$ is the elastic energy of the system; $W_b(\mathbf{u})$ is work done by body forces due to the absolute temperature, solute mass fraction, and pore pressure; and $W_S(\mathbf{u})$ is work done by applied boundary stress.

Now we will specify $\mathcal{V}_S(\mathbf{u})$, $W_b(\mathbf{u})$, and $W_S(\mathbf{u})$. Here and in the following we will suppress the time dependence of the displacement vector \mathbf{u} for convenience. Given the displacement $\mathbf{u}(\mathbf{x}) = u(\mathbf{x})\mathbf{i} + v(\mathbf{x})\mathbf{j}$, the linearized strain is the second order symmetric tensor

$$\boldsymbol{\varepsilon}(\mathbf{u}) = \frac{1}{2}(\nabla\mathbf{u} + \nabla\mathbf{u}^T)$$

where

$$\varepsilon_{11} = u_x, \quad \varepsilon_{12} = \varepsilon_{21} = \frac{1}{2}(u_y + v_x), \quad \varepsilon_{22} = v_y$$

In linear elasticity, a relation between the stress tensor $\boldsymbol{\tau}$ and the linearized strain tensor $\boldsymbol{\varepsilon}(\mathbf{u})$ is given by the generalized Hooke's law

$$\tau_{ij}(\mathbf{u}) = a_{ijkl}\varepsilon_{kl}(\mathbf{u}) \quad (2.11)$$

where Einstein summation convention is used, indices run from 1 to n , and a_{ijkh} are the coefficients of elasticity independent of the strain tensor with the properties of symmetry

$$a_{ijkh} = a_{jihk} = a_{khij} \quad (2.12)$$

and of ellipticity: there exists a constant $\alpha > 0$ such that

$$a_{ijkh}\varepsilon_{ij}\varepsilon_{kh} \geq \alpha\varepsilon_{ij}\varepsilon_{ij}, \quad \forall \varepsilon_{ij}.$$

Specifically, for a homogeneous elastic isotropic medium, the stress-strain relation in terms of the bulk modulus K and the shear modulus G has the form

$$\boldsymbol{\tau} = 2G\boldsymbol{\varepsilon} + \left(K - \frac{2G}{3}\right)(\text{tr}\boldsymbol{\varepsilon})I \quad (2.13)$$

With this notation, the elastic strain energy of the system is

$$\mathcal{V}_S(\mathbf{u}) = \frac{1}{2} \int_{\Omega} \tau_{11}(\mathbf{u})\varepsilon_{11}(\mathbf{u}) + 2\tau_{12}(\mathbf{u})\varepsilon_{12}(\mathbf{u}) + \tau_{22}(\mathbf{u})\varepsilon_{22}(\mathbf{u})d\Omega; \quad (2.14)$$

the work done by body forces due to the absolute temperature, solute mass fraction, and pore pressure is

$$W_b(\mathbf{u}) = \int_{\Omega} (\mathbf{b}_1 \cdot \mathbf{V})(\text{tr}\boldsymbol{\varepsilon}(\mathbf{u}))d\Omega; \quad (2.15)$$

and the work done by applied boundary stress is

$$W_S(\mathbf{u}) = \int_{\Gamma} (\hat{\boldsymbol{\sigma}}\mathbf{n}) \cdot \mathbf{u}d\Gamma \quad (2.16)$$

At this point, we wish to express the elastic strain energy $\mathcal{V}_S(\mathbf{u})$ as a functional on \mathbb{V}^2 . Let $\boldsymbol{\Phi}(\mathbf{x}) = \phi(\mathbf{x})\mathbf{i} + \psi(\mathbf{x})\mathbf{j}$. Define the bilinear form $a_E : \mathbb{V}^2 \times \mathbb{V}^2 \rightarrow \mathbb{R}$ by

$$a_E(\mathbf{u}, \boldsymbol{\Phi}) = \int_{\Omega} \tau_{11}(\mathbf{u})\varepsilon_{11}(\boldsymbol{\Phi}) + 2\tau_{12}(\mathbf{u})\varepsilon_{12}(\boldsymbol{\Phi}) + \tau_{22}(\mathbf{u})\varepsilon_{22}(\boldsymbol{\Phi})d\Omega, \quad (2.17)$$

for $\mathbf{u}, \boldsymbol{\Phi} \in \mathbb{V}^2$. From (2.10) and (2.14)-(2.17), the total potential energy of the system has the form

$$\mathcal{V}(\mathbf{u}) = \frac{1}{2}a_E(\mathbf{u}, \mathbf{u}) - \int_{\Omega} (\mathbf{b}_1 \cdot \mathbf{V})\nabla \cdot \mathbf{u}d\Omega - \int_{\Gamma} (\hat{\boldsymbol{\sigma}}\mathbf{n}) \cdot \mathbf{u}d\Gamma \quad (2.18)$$

Define the following vectors:

$$\boldsymbol{\tau}_1 = [\tau_{11} \quad \tau_{12}]^T, \quad \boldsymbol{\tau}_2 = [\tau_{21} \quad \tau_{22}]^T$$

Then

$$\begin{aligned} a_E(\mathbf{u}, \boldsymbol{\Phi}) &= \int_{\Omega} \tau_{11}(\mathbf{u})\phi_x + \tau_{12}(\mathbf{u})\phi_y + \tau_{21}(\mathbf{u})\psi_x + \tau_{22}(\mathbf{u})\psi_y d\Omega \\ &= \int_{\Omega} \boldsymbol{\tau}_1(\mathbf{u}) \cdot \nabla\phi + \boldsymbol{\tau}_2(\mathbf{u}) \cdot \nabla\psi d\Omega \\ &= \int_{\Gamma} (\boldsymbol{\tau}(\mathbf{u})\mathbf{n}) \cdot \boldsymbol{\Phi}d\Gamma - \int_{\Omega} [\nabla \cdot \boldsymbol{\tau}_1 \quad \nabla \cdot \boldsymbol{\tau}_2]^T(\mathbf{u}) \cdot \boldsymbol{\Phi}d\Omega \end{aligned}$$

and we obtain the Green's formula:

$$a_E(\mathbf{u}, \boldsymbol{\Phi}) = \int_{\Gamma} (\boldsymbol{\tau}(\mathbf{u})\mathbf{n}) \cdot \boldsymbol{\Phi}d\Gamma - \int_{\Omega} [\nabla \cdot \boldsymbol{\tau}_1 \quad \nabla \cdot \boldsymbol{\tau}_2]^T(\mathbf{u}) \cdot \boldsymbol{\Phi}d\Omega, \quad \forall \mathbf{u}, \boldsymbol{\Phi} \in \mathbb{V}^2 \quad (2.19)$$

Referring to (2.4) and (2.5), the set of admissible displacements, in general, is

$$\mathcal{U}_{ad} = \{\mathbf{u} \in \tilde{\mathbb{V}}_0^n : \dot{\mathbf{u}} \in \tilde{\mathbb{V}}_0^n\}, \quad n = 2 \text{ or } 3 \quad (2.20)$$

Using the principle of minimum total potential energy, the displacement $\mathbf{u} \in \mathcal{U}_{ad}$ that the region $\Omega \subset \mathbb{R}^2$ undergoes is given by

$$D\mathcal{V}(\mathbf{u})\Phi = 0, \quad \forall \Phi \in \tilde{\mathcal{V}}_0^2 \quad (2.21)$$

where

$$D\mathcal{V}(\mathbf{u})\Phi = \left. \frac{d}{d\delta} \mathcal{V}(\mathbf{u} + \delta\Phi) \right|_{\delta=0}$$

is the Gâteaux differential of \mathcal{V} with increment Φ . From (2.18) and (2.21),

$$a_E(\mathbf{u}, \Phi) - \int_{\Omega} (\mathbf{b}_1 \cdot \mathbf{V}) \nabla \cdot \Phi d\Omega - \int_{\Gamma \setminus \Gamma_F} (\hat{\boldsymbol{\sigma}} \mathbf{n}) \cdot \Phi d\Gamma = 0, \quad \forall \Phi \in \tilde{\mathcal{V}}_0^2 \quad (2.22)$$

and applying the divergence theorem we have

$$\begin{aligned} a_E(\mathbf{u}, \Phi) - \int_{\Gamma \setminus \Gamma_F} ((\mathbf{b}_1 \cdot \mathbf{V}) I \mathbf{n}) \cdot \Phi d\Gamma \\ + \int_{\Omega} \nabla (\mathbf{b}_1 \cdot \mathbf{V}) \cdot \Phi d\Omega - \int_{\Gamma \setminus \Gamma_F} (\hat{\boldsymbol{\sigma}} \mathbf{n}) \cdot \Phi d\Gamma = 0, \quad \forall \Phi \in \tilde{\mathcal{V}}_0^2 \end{aligned} \quad (2.23)$$

Green's formula (2.19) and (2.23) yield

$$[\nabla \cdot \boldsymbol{\tau}_1 \quad \nabla \cdot \boldsymbol{\tau}_2]^T = \nabla (\mathbf{b}_1 \cdot \mathbf{V}), \quad \text{in } \Omega \quad (2.24)$$

$$\boldsymbol{\tau} \mathbf{n} = ((\mathbf{b}_1 \cdot \mathbf{V}) I + \hat{\boldsymbol{\sigma}}) \mathbf{n}, \quad \text{on } \Gamma \setminus \Gamma_F, \quad (2.25)$$

and the stress-strain relation (2.13) gives

$$[\nabla \cdot \boldsymbol{\tau}_1 \quad \nabla \cdot \boldsymbol{\tau}_2]^T = \left(K + \frac{G}{3} \right) \nabla (\nabla \cdot \mathbf{u}) + G \nabla^2 \mathbf{u} \quad (2.26)$$

From (2.24)-(2.26) and (2.4), we obtain the system (2.8), (2.9), and (2.4).

On the other hand, since $\dot{\mathbf{u}} \in \tilde{\mathcal{V}}_0^2$, differentiating (2.22) with respect to time, we have

$$a_E(\dot{\mathbf{u}}, \Phi) - \int_{\Omega} (\mathbf{b}_1 \cdot \dot{\mathbf{V}}) \nabla \cdot \Phi d\Omega - \int_{\Gamma \setminus \Gamma_F} (\dot{\boldsymbol{\sigma}} \mathbf{n}) \cdot \Phi d\Gamma = 0, \quad \forall \Phi \in \tilde{\mathcal{V}}_0^2$$

Using the same argument as above leads to the equivalence of the systems (2.8), (2.9), (2.4) and (2.2), (2.4)-(2.6). \square

Remark 2.3. From the condition of mechanical equilibrium

$$\nabla \cdot \boldsymbol{\sigma} = 0$$

where $\boldsymbol{\sigma} = \boldsymbol{\tau} - (\mathbf{b}_1 \cdot \mathbf{V}) I$ is the poroelastic stress tensor, and from the stress-strain relation (2.13), it follows that (2.8) constitutes the equation of equilibrium.

3. WELL-POSEDNESS OF THE DIFFUSION AND ELASTIC SYSTEMS

In this section, we discuss the well-posedness of the parabolic initial-boundary value problem for the diffusion system and the elliptic initial-boundary value problem for the elastic system considering coupling terms as data.

3.1. Diffusion system. The following assumptions will be made on the matrices of diffusion coefficients and the boundary and initial functions.

Assumption 3.1. The matrix M is non-singular and all the eigenvalues of the matrix $M^{-1}A$ are positive: $\lambda_i > 0$, $i = 1, 2, 3$.

Under Assumption 3.1, the matrix $M^{-1}A$ admits the eigendecomposition

$$M^{-1}A = PDP^{-1} \quad (3.1)$$

where $P = [\mathbf{e}_1 \ \mathbf{e}_2 \ \mathbf{e}_3]$, \mathbf{e}_i is the eigenvector of $M^{-1}A$ corresponding to the eigenvalue λ_i , $\|\mathbf{e}_i\| = 1$, $i = 1, 2, 3$, and $D = \text{diag}(\lambda_1, \lambda_2, \lambda_3)$.

Assumption 3.2. $\mathbf{V}_B \in L^2(0, t_f; H^{1/2}(\Gamma)^3)$, $\dot{\mathbf{V}}_B \in L^2(0, t_f; H^{1/2}(\Gamma)^3)$, and $\mathbf{V}_I \in \mathbb{V}^3$.

We transform the parabolic diffusion system (2.1), (2.3), and (2.7) to the equivalent diagonalized system with homogeneous boundary conditions. The inverse trace theorem [24] yields that there exists a continuous mapping

$$\gamma_0^- : L^2(0, t_f; H^{1/2}(\Gamma)^3) \rightarrow L^2(0, t_f; \mathbb{V}^3)$$

such that $\gamma_0 \gamma_0^- w = w$, $w \in L^2(0, t_f; H^{1/2}(\Gamma)^3)$, where

$$\begin{aligned} \gamma_0 : L^2(0, t_f; \mathbb{V}^3) &\rightarrow L^2(0, t_f; H^{1/2}(\Gamma)^3) \\ \gamma_0(v) &= v|_{\Gamma}, \quad v \in L^2(0, t_f; \mathbb{V}^3) \end{aligned}$$

is a linear and continuous trace mapping. We define the vector $\mathbf{W} \in L^2(0, t_f; \mathbb{V}^3)$ by

$$\mathbf{W}(\mathbf{x}, t) = \gamma_0^- \mathbf{V}_B(\mathbf{x}, t) \quad (3.2)$$

and the vector

$$\mathbf{U}(\mathbf{x}, t) = P^{-1}(\mathbf{V}(\mathbf{x}, t) - \mathbf{W}(\mathbf{x}, t)) \quad (3.3)$$

The transformation given by (3.1)-(3.3) leads to the following initial-boundary value problem equivalent to the parabolic diffusion system (2.1), (2.3), and (2.7):

$$\dot{\mathbf{U}} - D\nabla^2 \mathbf{U} = -P^{-1} \dot{\mathbf{W}} + DP^{-1} \nabla^2 \mathbf{W} - P^{-1} M^{-1} \mathbf{b}_0(\nabla \cdot \dot{\mathbf{u}}), \quad \text{in } \Omega \times (0, t_f) \quad (3.4)$$

$$\mathbf{U}(\mathbf{x}, t) = \mathbf{0}, \quad \text{on } \Gamma \times (0, t_f) \quad (3.5)$$

$$\mathbf{U}(\mathbf{x}, 0) = \mathbf{U}_0(\mathbf{x}), \quad \text{in } \bar{\Omega} \quad (3.6)$$

where

$$\mathbf{U}_0(\mathbf{x}) = \begin{cases} P^{-1}(\mathbf{V}_I(\mathbf{x}) - \mathbf{W}(\mathbf{x}, 0)), & \text{in } \Omega \\ \mathbf{0}, & \text{on } \Gamma \end{cases} \quad (3.7)$$

Remark 3.3. From (2.20), $\mathbf{u} \in \mathcal{U}_{ad}$ satisfies $\nabla \cdot \dot{\mathbf{u}} \in L^2(0, t_f; L^2(\Omega))$. Also, Assumption 3.2 and (3.2) yield $\dot{\mathbf{W}} \in L^2(0, t_f; \mathbb{V}^3)$ and, with the use of (3.7), it follows that $\mathbf{U}_0 \in \mathbb{V}_0^3$.

Define a bilinear form $a : \mathbb{V}_0^3 \times \mathbb{V}_0^3 \rightarrow \mathbb{R}$ by

$$a(\mathbf{v}, \mathbf{w}) = \sum_{k=1}^3 \lambda_k \int_{\Omega} \nabla v_k \cdot \nabla w_k d\Omega,$$

where $\mathbf{v} = [v_1 \ v_2 \ v_3]^T$ and $\mathbf{w} = [w_1 \ w_2 \ w_3]^T$. The form $a(\cdot, \cdot)$ defines a scalar product on \mathbb{V}_0^3 with the norm equivalent to the standard norm on \mathbb{V}_0^3 and the following inequalities hold:

$$|a(\mathbf{v}, \mathbf{w})| \leq 3\lambda_{\max} \|\mathbf{v}\|_{\mathbb{V}_0^3} \|\mathbf{w}\|_{\mathbb{V}_0^3}, \quad \mathbf{v}, \mathbf{w} \in \mathbb{V}_0^3 \quad (3.8)$$

$$\lambda_{\min} \|\mathbf{v}\|_{\mathbb{V}_0^3}^2 \leq a(\mathbf{v}, \mathbf{v}) \leq \lambda_{\max} \|\mathbf{v}\|_{\mathbb{V}_0^3}^2, \quad \mathbf{v} \in \mathbb{V}_0^3 \quad (3.9)$$

We denote this scalar product by $(\cdot, \cdot)_a$; that is,

$$a(\mathbf{v}, \mathbf{w}) = (\mathbf{v}, \mathbf{w})_a \quad (3.10)$$

and the corresponding norm is $\|\cdot\|_a = (\cdot, \cdot)_a^{1/2}$.

We consider the following weak formulation of the problem (3.4)-(3.7).

Given $\mathbf{W} \in L^2(0, t_f; \mathbb{V}^3)$, $\dot{\mathbf{W}} \in L^2(0, t_f; \mathbb{V}^3)$, $\nabla \cdot \dot{\mathbf{u}} \in L^2(0, t_f; L^2(\Omega))$, and $\mathbf{U}_0 \in \mathbb{H}^3$, find $\mathbf{U} \in L^2(0, t_f; \mathbb{V}_0^3)$ such that, for all $\boldsymbol{\varphi} \in \mathbb{V}_0^3$,

$$\begin{aligned} \frac{d}{dt}(\mathbf{U}, \boldsymbol{\varphi})_{\mathbb{H}^3} + (\mathbf{U}, \boldsymbol{\varphi})_a &= (-P^{-1}\dot{\mathbf{W}}, \boldsymbol{\varphi})_{\mathbb{H}^3} + (DP^{-1}\nabla^2\mathbf{W}, \boldsymbol{\varphi})_{\mathbb{H}^3} \\ &+ (-P^{-1}M^{-1}\mathbf{b}_0(\nabla \cdot \dot{\mathbf{u}}), \boldsymbol{\varphi})_{\mathbb{H}^3} \end{aligned} \quad (3.11)$$

and

$$\mathbf{U}(0) = \mathbf{U}_0 \quad (3.12)$$

The next lemma establishes the Hadamard well-posedness of the problem (3.11) and (3.12).

Lemma 3.4. *Problem (3.11) and (3.12) admits a unique solution $\mathbf{U} \in \mathcal{W}(0, t_f; \mathbb{V}_0^3)$ and this solution depends continuously on the data, that is the mapping*

$$\mathbf{W}, \dot{\mathbf{W}}, \nabla \cdot \dot{\mathbf{u}}, \mathbf{U}_0 \mapsto \mathbf{U}$$

from $L^2(0, t_f; \mathbb{V}^3) \times L^2(0, t_f; \mathbb{V}^3) \times L^2(0, t_f; L^2(\Omega)) \times \mathbb{H}^3$ to $\mathcal{W}(0, t_f; \mathbb{V}_0^3)$ is continuous.

Proof. The result is based on the standard Faedo-Galerkin approximation technique and we omit the proof of the existence and uniqueness of a solution, as it essentially follows the proofs of [17, Theorem 1.2, Chapter III] and [25, Theorem 1.1, Chapter III].

We will show that the solution $\mathbf{U} \in \mathcal{W}(0, t_f; \mathbb{V}_0^3)$ depends continuously on the data $\mathbf{W}, \dot{\mathbf{W}}, \nabla \cdot \dot{\mathbf{u}}$ and \mathbf{U}_0 . For each $m \in \mathbb{Z}_+$, define an approximate solution \mathbf{U}_m as

$$\mathbf{U}_m(\mathbf{x}, t) = \sum_{i=1}^m g_{im}(t) \mathbf{w}_i(\mathbf{x})$$

where $g_{im}(t)$, $1 \leq i \leq m$, are scalar functions defined on $[0, t_f]$, $\mathbf{w}_1, \dots, \mathbf{w}_m, \dots$ is a countable set of functions which is dense in \mathbb{V}_0^3 , $\mathbf{U}_m(\mathbf{x}, t)$ satisfies

$$\begin{aligned} &(\dot{\mathbf{U}}_m(t), \mathbf{w}_i)_{\mathbb{H}^3} + (\mathbf{U}_m(t), \mathbf{w}_i)_a \\ &= (-P^{-1}\dot{\mathbf{W}}(t), \mathbf{w}_i)_{\mathbb{H}^3} + (\nabla(-DP^{-1}\mathbf{W}(t)), \nabla\mathbf{w}_i)_{\mathbb{H}^3} \\ &+ (-P^{-1}M^{-1}\mathbf{b}_0(\nabla \cdot \dot{\mathbf{u}}(t)), \mathbf{w}_i)_{\mathbb{H}^3}, \quad 1 \leq i \leq m \end{aligned} \quad (3.13)$$

and

$$\mathbf{U}_m(\mathbf{x}, 0) = \mathbf{U}_{0m} \xrightarrow{m \rightarrow \infty} \mathbf{U}_0 \quad \text{in } \mathbb{H}^3 \quad (3.14)$$

Multiplying (3.13) by $g_{jm}(t)$, adding for j , and using

$$(\dot{\mathbf{U}}_m(t), \mathbf{U}_m(t))_{\mathbb{H}^3} = \frac{1}{2} \|\dot{\mathbf{U}}_m(t)\|_{\mathbb{H}^3}^2$$

we have

$$\begin{aligned} & \frac{d}{dt} \|\mathbf{U}_m(t)\|_{\mathbb{H}^3}^2 + 2\|\mathbf{U}_m(t)\|_a^2 \\ & \leq 2|(P^{-1}\dot{\mathbf{W}}(t), \mathbf{U}_m(t))_{\mathbb{H}^3}| + 2|(\nabla(DP^{-1}\mathbf{W}(t)), \nabla\mathbf{U}_m(t))_{\mathbb{H}^3}| \\ & \quad + 2|(P^{-1}M^{-1}\mathbf{b}_0(\nabla \cdot \dot{\mathbf{u}}(t)), \mathbf{U}_m(t))_{\mathbb{H}^3}| \end{aligned}$$

Applying (3.9), (3.10), and the Poincaré inequality gives

$$\begin{aligned} & \frac{d}{dt} \|\mathbf{U}_m(t)\|_{\mathbb{H}^3}^2 + 2\lambda_{\min} \|\mathbf{U}_m(t)\|_{\mathbb{V}_0^3}^2 \\ & \leq \frac{3c}{\lambda_{\min}} \|P^{-1}\|^2 \|\dot{\mathbf{W}}\|_{\mathbb{V}^3}^2 + \frac{\lambda_{\min}}{3c} c \|\mathbf{U}_m(t)\|_{\mathbb{V}_0^3}^2 \\ & \quad + \frac{3}{\lambda_{\min}} \|DP^{-1}\|^2 \|\mathbf{W}\|_{\mathbb{V}^3}^2 + \frac{\lambda_{\min}}{3} \|\mathbf{U}_m(t)\|_{\mathbb{V}_0^3}^2 \\ & \quad + \frac{3c}{\lambda_{\min}} |P^{-1}M^{-1}\mathbf{b}_0|^2 \|\nabla \cdot \dot{\mathbf{u}}(t)\|_{L^2(\Omega)}^2 + \frac{\lambda_{\min}}{3c} c \|\mathbf{U}_m(t)\|_{\mathbb{V}_0^3}^2 \end{aligned}$$

and therefore

$$\begin{aligned} \frac{d}{dt} \|\mathbf{U}_m(t)\|_{\mathbb{H}^3}^2 + \lambda_{\min} \|\mathbf{U}_m(t)\|_{\mathbb{V}_0^3}^2 & \leq \frac{3c}{\lambda_{\min}} \|P^{-1}\|^2 \|\dot{\mathbf{W}}\|_{\mathbb{V}^3}^2 + \frac{3\lambda_{\max}^2}{\lambda_{\min}} \|P^{-1}\|^2 \|\mathbf{W}\|_{\mathbb{V}^3}^2 \\ & \quad + \frac{3c}{\lambda_{\min}} |P^{-1}M^{-1}\mathbf{b}_0|^2 \|\nabla \cdot \dot{\mathbf{u}}(t)\|_{L^2(\Omega)}^2 \end{aligned}$$

Integrating the last inequality over $(0, t_f)$ and applying (3.14), we obtain

$$\begin{aligned} \|\mathbf{U}_m\|_{L^2(0, t_f; \mathbb{V}_0^3)}^2 & \leq \frac{1}{\lambda_{\min}} \|\mathbf{U}_0\|_{\mathbb{H}^3}^2 + \frac{3c}{\lambda_{\min}^2} \|P^{-1}\|^2 \|\dot{\mathbf{W}}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 \\ & \quad + \frac{3\lambda_{\max}^2}{\lambda_{\min}^2} \|P^{-1}\|^2 \|\mathbf{W}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 \\ & \quad + \frac{3c}{\lambda_{\min}^2} |P^{-1}M^{-1}\mathbf{b}_0|^2 \|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(0, t_f; L^2(\Omega))}^2 \end{aligned} \tag{3.15}$$

The inequality (3.15) holds for each $m \in \mathbb{Z}_+$ and therefore,

$$\begin{aligned} \|\mathbf{U}\|_{L^2(0, t_f; \mathbb{V}_0^3)}^2 & \leq \frac{1}{\lambda_{\min}} \|\mathbf{U}_0\|_{\mathbb{H}^3}^2 + \frac{3c}{\lambda_{\min}^2} \|P^{-1}\|^2 \|\dot{\mathbf{W}}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 \\ & \quad + \frac{3\lambda_{\max}^2}{\lambda_{\min}^2} \|P^{-1}\|^2 \|\mathbf{W}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 \\ & \quad + \frac{3c}{\lambda_{\min}^2} |P^{-1}M^{-1}\mathbf{b}_0|^2 \|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(0, t_f; L^2(\Omega))}^2 \end{aligned} \tag{3.16}$$

To estimate $\dot{\mathbf{U}}_m(t)$, note that $\dot{\mathbf{U}}_m(t) = \sum_{i=1}^m \dot{g}_{im}(t) \mathbf{w}_i(\mathbf{x}) \in \text{span}\{\mathbf{w}_1, \dots, \mathbf{w}_m\} = E_m \subset \mathbb{V}_0^3$ and

$$\|\dot{\mathbf{U}}_m(t)\|_{H^{-1}(\Omega)^3} = \sup_{\mathbf{v} \in E_m \setminus \{0\}} \frac{|(\dot{\mathbf{U}}_m(t), \mathbf{v})_{\mathbb{H}^3}|}{\|\mathbf{v}\|_{\mathbb{V}_0^3}} \tag{3.17}$$

Using (3.8), (3.10), (3.13), the Poincaré inequality, and linearity argument yields

$$\begin{aligned} |(\dot{\mathbf{U}}_m(t), \boldsymbol{\varphi})_{\mathbb{H}^3}| &\leq |a(\mathbf{U}_m(t), \boldsymbol{\varphi})| + |(P^{-1}\dot{\mathbf{W}}, \boldsymbol{\varphi})_{\mathbb{H}^3}| + |(\nabla(DP^{-1}\mathbf{W}), \nabla\boldsymbol{\varphi})_{\mathbb{H}^3}| \\ &\quad + |(P^{-1}M^{-1}\mathbf{b}_0(\nabla \cdot \dot{\mathbf{u}}), \boldsymbol{\varphi})_{\mathbb{H}^3}| \\ &\leq \left(3\lambda_{\max}\|\mathbf{U}_m(t)\|_{\mathbb{V}_0^3} + c\|P^{-1}\|\|\dot{\mathbf{W}}\|_{\mathbb{V}^3} + \lambda_{\max}\|P^{-1}\|\|\mathbf{W}\|_{\mathbb{V}^3} \right. \\ &\quad \left. + c|P^{-1}M^{-1}\mathbf{b}_0|\|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(\Omega)} \right) \|\boldsymbol{\varphi}\|_{\mathbb{V}_0^3}, \quad \forall \boldsymbol{\varphi} \in \mathbb{V}_0^3 \end{aligned} \quad (3.18)$$

From (3.17) and (3.18), we have

$$\begin{aligned} \|\dot{\mathbf{U}}_m(t)\|_{H^{-1}(\Omega)^3} &\leq 3\lambda_{\max}\|\mathbf{U}_m(t)\|_{\mathbb{V}_0^3} + c\|P^{-1}\|\|\dot{\mathbf{W}}\|_{\mathbb{V}^3} + \lambda_{\max}\|P^{-1}\|\|\mathbf{W}\|_{\mathbb{V}^3} \\ &\quad + c|P^{-1}M^{-1}\mathbf{b}_0|\|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(\Omega)} \end{aligned}$$

Squaring the last inequality and integrating it over $(0, t_f)$ gives

$$\begin{aligned} \|\dot{\mathbf{U}}_m\|_{L^2(0, t_f; H^{-1}(\Omega)^3)}^2 &\leq 36\lambda_{\max}^2\|\mathbf{U}_m\|_{L^2(0, t_f; \mathbb{V}_0^3)}^2 \\ &\quad + 4c^2\|P^{-1}\|^2\|\dot{\mathbf{W}}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 + 4\lambda_{\max}^2\|P^{-1}\|^2\|\mathbf{W}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 \\ &\quad + 4c^2|P^{-1}M^{-1}\mathbf{b}_0|^2\|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(0, t_f; L^2(\Omega))}^2 \end{aligned}$$

Applying (3.15) to the first term on the right-hand side of this inequality, we obtain

$$\begin{aligned} \|\dot{\mathbf{U}}_m\|_{L^2(0, t_f; H^{-1}(\Omega)^3)}^2 &\leq \frac{36\lambda_{\max}^2}{\lambda_{\min}}\|\mathbf{U}_0\|_{\mathbb{H}^3}^2 \\ &\quad + \left(\frac{108c\lambda_{\max}^2}{\lambda_{\min}^2} + 4c^2 \right) \|P^{-1}\|^2\|\dot{\mathbf{W}}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 \\ &\quad + \left(\frac{108\lambda_{\max}^4}{\lambda_{\min}^2} + 4\lambda_{\max}^2 \right) \|P^{-1}\|^2\|\mathbf{W}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 \\ &\quad + \left(\frac{108c\lambda_{\max}^2}{\lambda_{\min}^2} + 4c^2 \right) |P^{-1}M^{-1}\mathbf{b}_0|^2\|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(0, t_f; L^2(\Omega))}^2 \end{aligned}$$

The last inequality holds for each $m \in \mathbb{Z}_+$ and therefore,

$$\begin{aligned} &\|\dot{\mathbf{U}}\|_{L^2(0, t_f; H^{-1}(\Omega)^3)}^2 \\ &\leq \frac{36\lambda_{\max}^2}{\lambda_{\min}}\|\mathbf{U}_0\|_{\mathbb{H}^3}^2 + \left(\frac{108c\lambda_{\max}^2}{\lambda_{\min}^2} + 4c^2 \right) \|P^{-1}\|^2\|\dot{\mathbf{W}}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 \\ &\quad + \left(\frac{108\lambda_{\max}^4}{\lambda_{\min}^2} + 4\lambda_{\max}^2 \right) \|P^{-1}\|^2\|\mathbf{W}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 \\ &\quad + \left(\frac{108c\lambda_{\max}^2}{\lambda_{\min}^2} + 4c^2 \right) |P^{-1}M^{-1}\mathbf{b}_0|^2\|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(0, t_f; L^2(\Omega))}^2 \end{aligned} \quad (3.19)$$

The result follows from (3.16) and (3.19). \square

Lemma 3.4 yields the Hadamard well-posedness in the weak sense of the diffusion system (2.1), (2.3), and (2.7), namely, the existence and uniqueness of a weak solution to the model and its continuous dependence on the boundary data, initial data, and the divergence of the rock deformation velocity field.

Theorem 3.5 (Well-posedness of the diffusion system). *Given boundary data $\mathbf{V}_B \in L^2(0, t_f; H^{1/2}(\Gamma)^3)$ with $\dot{\mathbf{V}}_B \in L^2(0, t_f; H^{1/2}(\Gamma)^3)$, initial data $\mathbf{V}_I \in \mathbb{H}^3$, and $\nabla \cdot \dot{\mathbf{u}} \in L^2(0, t_f; L^2(\Omega))$, the diffusion system (2.1), (2.3), and (2.7) admits*

a unique weak solution $\mathbf{V} \in L^2(0, t_f; \mathbb{V}^3)$ with $\dot{\mathbf{V}} \in L^2(0, t_f; H^{-1}(\Omega)^3)$ and this solution depends continuously on the data $\mathbf{V}_B, \dot{\mathbf{V}}_B, \mathbf{V}_B(0), \mathbf{V}_I$, and $\nabla \cdot \dot{\mathbf{u}}$. That is, the mapping

$$\mathbf{V}_B, \dot{\mathbf{V}}_B, \mathbf{V}_B(0), \mathbf{V}_I, \nabla \cdot \dot{\mathbf{u}} \mapsto \mathbf{V}, \dot{\mathbf{V}}$$

from $L^2(0, t_f; H^{1/2}(\Gamma)^3) \times L^2(0, t_f; H^{1/2}(\Gamma)^3) \times H^{1/2}(\Gamma)^3 \times \mathbb{H}^3 \times L^2(0, t_f; L^2(\Omega))$ to $L^2(0, t_f; \mathbb{V}^3) \times L^2(0, t_f; H^{-1}(\Omega)^3)$ is continuous.

Proof. The existence and uniqueness of a weak solution follow immediately from the transformation (3.3). It remains to prove the continuous dependence on data. From (3.2), (3.3) and the Poincaré inequality, we have

$$\|\mathbf{V}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 \leq 2(1+c)\|P\|^2\|\mathbf{U}\|_{L^2(0, t_f; \mathbb{V}_0^3)}^2 + 2\hat{C}\|\mathbf{V}_B\|_{L^2(0, t_f; H^{1/2}(\Gamma)^3)}^2 \quad (3.20)$$

$$\|\dot{\mathbf{V}}\|_{L^2(0, t_f; H^{-1}(\Omega)^3)}^2 \leq 2\|P\|^2\|\dot{\mathbf{U}}\|_{L^2(0, t_f; H^{-1}(\Omega)^3)}^2 + 2\hat{C}\|\dot{\mathbf{V}}_B\|_{L^2(0, t_f; H^{1/2}(\Gamma)^3)}^2 \quad (3.21)$$

where $\hat{C} > 0$ is a constant. Also, the transformation (3.7) gives

$$\|\mathbf{U}_0\|_{\mathbb{H}^3}^2 \leq 2\|P^{-1}\|^2\|\mathbf{V}_I\|_{\mathbb{H}^3}^2 + 2\|P^{-1}\|^2\hat{C}\|\mathbf{V}_B(0)\|_{H^{1/2}(\Gamma)^3}^2 \quad (3.22)$$

Applying (3.22) to (3.16) and (3.19) and substituting the results together with

$$\|\mathbf{W}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 \leq \hat{C}\|\mathbf{V}_B\|_{L^2(0, t_f; H^{1/2}(\Gamma)^3)}^2$$

$$\|\dot{\mathbf{W}}\|_{L^2(0, t_f; \mathbb{V}^3)}^2 \leq \hat{C}\|\dot{\mathbf{V}}_B\|_{L^2(0, t_f; H^{1/2}(\Gamma)^3)}^2$$

into (3.20) and (3.21), respectively, complete the proof. \square

The next lemma further characterizes the solution \mathbf{V} to the diffusion system (2.1), (2.3), and (2.7) dealing with the initial data $\mathbf{V}_I \in \mathbb{V}^3$. This result will be needed in the later discussion of the coupled TCP model.

Lemma 3.6. *Under assumption $\mathbf{V}_I \in \mathbb{V}^3$, the weak solution \mathbf{V} of the diffusion system (2.1), (2.3), and (2.7) satisfies $\dot{\mathbf{V}} \in L^2(0, t_f; \mathbb{H}^3)$ and the following a priori estimate holds*

$$\|\dot{\mathbf{V}}\|_{L^2(0, t_f; \mathbb{H}^3)}^2 \leq \left(4 + \frac{6}{\lambda_{\min}^2}\right)\|P\|^2|P^{-1}M^{-1}\mathbf{b}_0|^2\|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(0, t_f; L^2(\Omega))}^2 + \theta \quad (3.23)$$

where $\theta > 0$ is a constant that does not depend on $\nabla \cdot \dot{\mathbf{u}}$.

Proof. As before, we use the standard Faedo-Galerkin approximation method applied to the problem (3.11) and (3.12). Note that the assumption $\mathbf{V}_I \in \mathbb{V}^3$ and (3.7) imply $\mathbf{U}_0 \in \mathbb{V}_0^3$. Multiplying (3.13) by $\dot{g}_{jm}(t)$, adding for j , and using $(\mathbf{U}_m(t), \dot{\mathbf{U}}_m(t))_a = \frac{1}{2}\|\dot{\mathbf{U}}_m(t)\|_a$ yield the inequality

$$\begin{aligned} & 2\|\dot{\mathbf{U}}_m(t)\|_{\mathbb{H}^3}^2 + \frac{d}{dt}\|\mathbf{U}_m(t)\|_a^2 \\ & \leq 2|(P^{-1}\dot{\mathbf{W}}, \dot{\mathbf{U}}_m(t))_{\mathbb{H}^3}| + 2\frac{d}{dt}(\nabla(-DP^{-1}\mathbf{W}), \nabla\mathbf{U}_m(t))_{\mathbb{H}^3} \\ & \quad + 2|(\nabla(DP^{-1}\dot{\mathbf{W}}), \nabla\mathbf{U}_m(t))_{\mathbb{H}^3}| + 2|(P^{-1}M^{-1}\mathbf{b}_0(\nabla \cdot \dot{\mathbf{u}}), \dot{\mathbf{U}}_m(t))_{\mathbb{H}^3}| \end{aligned}$$

which, in turn, gives

$$\begin{aligned} & 2\|\dot{\mathbf{U}}_m(t)\|_{\mathbb{H}^3}^2 + \frac{d}{dt}\|\mathbf{U}_m(t)\|_a^2 \\ & \leq 2\|P^{-1}\|^2\|\dot{\mathbf{W}}\|_{\mathbb{V}^3}^2 + \frac{1}{2}\|\dot{\mathbf{U}}_m(t)\|_{\mathbb{H}^3}^2 \end{aligned}$$

$$\begin{aligned}
 &+ 2 \frac{d}{dt} (\nabla(-DP^{-1}\mathbf{W}), \nabla\mathbf{U}_m(t))_{\mathbb{H}^3} + c\|DP^{-1}\|^2\|\dot{\mathbf{W}}\|_{\mathbb{V}^3}^2 \\
 &+ \frac{1}{c}\|\mathbf{U}_m(t)\|_{\mathbb{V}_0^3}^2 + 2|P^{-1}M^{-1}\mathbf{b}_0|^2\|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(\Omega)}^2 + \frac{1}{2}\|\dot{\mathbf{U}}_m(t)\|_{\mathbb{H}^3}^2
 \end{aligned}$$

and then

$$\begin{aligned}
 \|\dot{\mathbf{U}}_m(t)\|_{\mathbb{H}^3}^2 + \frac{d}{dt}\|\mathbf{U}_m(t)\|_a^2 &\leq (2 + c\lambda_{\max}^2)\|P^{-1}\|^2\|\dot{\mathbf{W}}\|_{\mathbb{V}^3}^2 + \frac{1}{c}\|\mathbf{U}_m(t)\|_{\mathbb{V}_0^3}^2 \\
 &+ 2 \frac{d}{dt} (\nabla(-DP^{-1}\mathbf{W}), \nabla\mathbf{U}_m(t))_{\mathbb{H}^3} \\
 &+ 2|P^{-1}M^{-1}\mathbf{b}_0|^2\|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(\Omega)}^2
 \end{aligned}$$

Integrating the above inequality over $(0, t_f)$ and using (3.9) and (3.10), we have

$$\begin{aligned}
 &\|\dot{\mathbf{U}}_m\|_{L^2(0,t_f;\mathbb{H}^3)}^2 + \lambda_{\min}\|\mathbf{U}_m(t_f)\|_{\mathbb{V}_0^3}^2 \\
 &\leq \lambda_{\max}\|\mathbf{U}_m(0)\|_{\mathbb{V}_0^3}^2 + (2 + c\lambda_{\max}^2)\|P^{-1}\|^2\|\dot{\mathbf{W}}\|_{L^2(0,t_f;\mathbb{V}^3)}^2 \\
 &\quad + \lambda_{\max}\|P^{-1}\|^2\|\mathbf{W}(0)\|_{\mathbb{V}^3}^2 + \lambda_{\max}\|\mathbf{U}_m(0)\|_{\mathbb{V}_0^3}^2 \\
 &\quad + \frac{\lambda_{\max}^2}{\lambda_{\min}}\|P^{-1}\|^2\|\mathbf{W}(t_f)\|_{\mathbb{V}^3}^2 + \lambda_{\min}\|\mathbf{U}_m(t_f)\|_{\mathbb{V}_0^3}^2 \\
 &\quad + \frac{1}{c}\|\mathbf{U}_m\|_{L^2(0,t_f;\mathbb{V}_0^3)}^2 + 2|P^{-1}M^{-1}\mathbf{b}_0|^2\|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(0,t_f;L^2(\Omega))}^2
 \end{aligned}$$

Applying (3.15) and the Poincaré inequality to the last inequality, we obtain

$$\begin{aligned}
 \|\dot{\mathbf{U}}_m\|_{L^2(0,t_f;\mathbb{H}^3)}^2 &\leq \left(\frac{1}{\lambda_{\min}} + 2\lambda_{\max}\right)\|\mathbf{U}_0\|_{\mathbb{V}_0^3}^2 + \lambda_{\max}\|P^{-1}\|^2\|\mathbf{W}(0)\|_{\mathbb{V}^3}^2 \\
 &\quad + \frac{\lambda_{\max}^2}{\lambda_{\min}}\|P^{-1}\|^2\|\mathbf{W}(t_f)\|_{\mathbb{V}^3}^2 + \frac{3\lambda_{\max}^2}{c\lambda_{\min}^2}\|P^{-1}\|^2\|\mathbf{W}\|_{L^2(0,t_f;\mathbb{V}^3)}^2 \\
 &\quad + \left(2 + c\lambda_{\max}^2 + \frac{3}{\lambda_{\min}^2}\right)\|P^{-1}\|^2\|\dot{\mathbf{W}}\|_{L^2(0,t_f;\mathbb{V}^3)}^2 \\
 &\quad + \left(2 + \frac{3}{\lambda_{\min}^2}\right)|P^{-1}M^{-1}\mathbf{b}_0|^2\|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(0,t_f;L^2(\Omega))}^2, \quad \forall m \in \mathbb{Z}_+
 \end{aligned}$$

This shows that the sequence $\dot{\mathbf{U}}_m$ ranges in a bounded set of $L^2(0, t_f; \mathbb{H}^3)$ and

$$\|\dot{\mathbf{U}}\|_{L^2(0,t_f;\mathbb{H}^3)}^2 \leq \left(2 + \frac{3}{\lambda_{\min}^2}\right)|P^{-1}M^{-1}\mathbf{b}_0|^2\|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(0,t_f;L^2(\Omega))}^2 + \hat{\theta} \tag{3.24}$$

where $\hat{\theta} > 0$ is a constant independent of $\nabla \cdot \dot{\mathbf{u}}$. From (3.2) and (3.3),

$$\|\dot{\mathbf{V}}\|_{L^2(0,t_f;\mathbb{H}^3)}^2 \leq 2\|P\|^2\|\dot{\mathbf{U}}\|_{L^2(0,t_f;\mathbb{H}^3)}^2 + 2\hat{C}\|\dot{\mathbf{V}}_B\|_{L^2(0,t_f;H^{1/2}(\Gamma)^3)}$$

where $\hat{C} > 0$ is a constant, and applying (3.24) the result follows. □

3.2. Elastic system. From now on, we will consider the Navier-type elastic system (2.8), (2.9), and (2.4). The following assumption will be made on the applied boundary stress tensor $\hat{\sigma}$.

Assumption 3.7.

$$\hat{\sigma} \in L^2(0, t_f; L^2(\Gamma)^{n \times n}) \quad \text{and} \quad \dot{\hat{\sigma}} \in L^2(0, t_f; L^2(\Gamma)^{n \times n})$$

where $n = 2$ or 3 is the dimension of the problem.

In solid mechanics the weak or variational formulation of a boundary value problem is equivalent to the principle of minimum total potential energy. Utilizing the notation introduced in Section 2, let $a_E : \mathbb{V}^n \times \mathbb{V}^n \rightarrow \mathbb{R}$, $n = 2$ or 3 , be a bilinear form defined as

$$a_E(\mathbf{u}, \Phi) = \int_{\Omega} \sum_{i,j=1}^n \tau_{ij}(\mathbf{u}) \varepsilon_{ij}(\Phi) d\Omega, \quad \mathbf{u}, \Phi \in \mathbb{V}^n \tag{3.25}$$

where $\boldsymbol{\tau} = [\tau_{ij}]$ and $\boldsymbol{\varepsilon} = [\varepsilon_{ij}]$, $i, j = 1, \dots, n$, are the stress and strain tensors, respectively. From (2.18), (2.20), and (2.21), under Assumption 3.7, we obtain the following weak formulation of problem (2.8), (2.9), and (2.4):

Given $\mathbf{V} \in L^2(0, t_f; \mathbb{V}^3)$ with $\dot{\mathbf{V}} \in L^2(0, t_f; \mathbb{H}^3)$ and $\hat{\boldsymbol{\sigma}} \in L^2(0, t_f; L^2(\Gamma)^{n \times n})$ with $\dot{\hat{\boldsymbol{\sigma}}} \in L^2(0, t_f; L^2(\Gamma)^{n \times n})$, find $\mathbf{u} \in L^2(0, t_f; \tilde{\mathbb{V}}_0^n)$ such that $\dot{\mathbf{u}} \in L^2(0, t_f; \tilde{\mathbb{V}}_0^n)$ and, for each $t \in [0, t_f)$,

$$a_E(\mathbf{u}, \Phi) - \int_{\Omega} (\mathbf{b}_1 \cdot \mathbf{V}) \nabla \cdot \Phi d\Omega - \int_{\Gamma \setminus \Gamma_F} (\hat{\boldsymbol{\sigma}} \mathbf{n}) \cdot \Phi d\Gamma = 0, \quad \forall \Phi \in \tilde{\mathbb{V}}_0^n \tag{3.26}$$

where \mathbf{n} is the outward unit normal vector on the boundary.

Next we will show that problem (3.26) is well-posed in the sense of Hadamard. Define continuous linear functionals F and \dot{F} on \mathbb{V}^n by the pairings

$$\Phi \mapsto \langle F, \Phi \rangle \quad \text{and} \quad \Phi \mapsto \langle \dot{F}, \Phi \rangle$$

where

$$\langle F, \Phi \rangle = \int_{\Omega} (\mathbf{b}_1 \cdot \mathbf{V}) \nabla \cdot \Phi d\Omega + \int_{\Gamma \setminus \Gamma_F} (\hat{\boldsymbol{\sigma}} \mathbf{n}) \cdot \Phi d\Gamma, \tag{3.27}$$

$$\langle \dot{F}, \Phi \rangle = \int_{\Omega} (\mathbf{b}_1 \cdot \dot{\mathbf{V}}) \nabla \cdot \Phi d\Omega + \int_{\Gamma \setminus \Gamma_F} (\dot{\hat{\boldsymbol{\sigma}}} \mathbf{n}) \cdot \Phi d\Gamma, \quad \forall \Phi \in \mathbb{V}^n \tag{3.28}$$

From (3.26)-(3.28), it follows that, for each $t \in [0, t_f)$,

$$a_E(\mathbf{u}, \Phi) = \langle F, \Phi \rangle, \quad a_E(\dot{\mathbf{u}}, \Phi) = \langle \dot{F}, \Phi \rangle, \quad \forall \Phi \in \tilde{\mathbb{V}}_0^n$$

and, using the trace theorem, for each $\Phi \in \tilde{\mathbb{V}}_0^n$ and $t \in [0, t_f)$,

$$\begin{aligned} |a_E(\mathbf{u}, \Phi)| &\leq \| \mathbf{b}_1 \cdot \mathbf{V} \|_{L^2(\Omega)} \| \nabla \cdot \Phi \|_{L^2(\Omega)} + \| \hat{\boldsymbol{\sigma}} \mathbf{n} \|_{L^2(\Gamma)^n} \| \Phi \|_{L^2(\Gamma)^n} \\ &\leq (\sqrt{n} | \mathbf{b}_1 | \| \mathbf{V} \|_{\mathbb{V}^n} + \hat{C} \| \hat{\boldsymbol{\sigma}} \|_{L^2(\Gamma)^{n \times n}}) \| \Phi \|_{\mathbb{V}^n} \end{aligned} \tag{3.29}$$

$$\begin{aligned} |a_E(\dot{\mathbf{u}}, \Phi)| &\leq \| \mathbf{b}_1 \cdot \dot{\mathbf{V}} \|_{L^2(\Omega)} \| \nabla \cdot \Phi \|_{L^2(\Omega)} + \| \dot{\hat{\boldsymbol{\sigma}}} \mathbf{n} \|_{L^2(\Gamma)^n} \| \Phi \|_{L^2(\Gamma)^n} \\ &\leq (\sqrt{n} | \mathbf{b}_1 | \| \dot{\mathbf{V}} \|_{\mathbb{H}^n} + \hat{C} \| \dot{\hat{\boldsymbol{\sigma}}} \|_{L^2(\Gamma)^{n \times n}}) \| \Phi \|_{\mathbb{V}^n} \end{aligned} \tag{3.30}$$

where $\hat{C} > 0$ is a constant.

To prove the existence and uniqueness of a solution to the problem (3.26), we will use of the Korn inequality [10, Theorem 3.1, Chapter III] and its consequence [10, Theorem 3.3, Chapter III] which imply that there exists a constant $c_E > 0$ such that

$$a_E(\mathbf{v}, \mathbf{v}) \geq c_E \| \mathbf{v} \|_{\tilde{\mathbb{V}}_0^n}^2, \quad \forall \mathbf{v} \in \tilde{\mathbb{V}}_0^n, \quad t \in [0, t_f) \tag{3.31}$$

Also, from (2.11), (2.12), and (3.25), we have

$$a_E(\mathbf{u}, \mathbf{v}) = a_E(\mathbf{v}, \mathbf{u}), \quad \forall \mathbf{u}, \mathbf{v} \in \mathbb{V}^n, \quad \forall \mathbf{v} \in \tilde{\mathbb{V}}_0^n \tag{3.32}$$

and, for each $t \in [0, t_f)$,

$$\begin{aligned}
 |a_E(\mathbf{u}, \mathbf{v})| &\leq \max_{i,j,k,l} \{a_{ijkl}\} \sum_{i,j,k,l=1}^n \left| \int_{\Omega} \varepsilon_{ij}(\mathbf{u}) \varepsilon_{kl}(\mathbf{v}) d\Omega \right| \\
 &\leq n \cdot \max_{i,j,k,l} \{a_{ijkl}\} \|\mathbf{u}\|_{\mathbb{V}^n} \|\mathbf{v}\|_{\mathbb{V}^n}
 \end{aligned}
 \tag{3.33}$$

Equations (3.31)-(3.33) yield that $a_E(\cdot, \cdot)$ is symmetric and continuous on $\tilde{\mathbb{V}}_0^n$ and there exist constants $c_E > 0$ and $\beta > 0$ such that

$$c_E \|\mathbf{u}\|_{\mathbb{V}^n}^2 \leq a_E(\mathbf{u}, \mathbf{u}) \leq \beta \|\mathbf{u}\|_{\mathbb{V}^n}^2, \quad \forall \mathbf{u} \in \tilde{\mathbb{V}}_0^n, t \in [0, t_f)$$

Thus, for each $t \in [0, t_f)$, $a_E(\cdot, \cdot)$ is an inner product on $\tilde{\mathbb{V}}_0^n$ that is associated with a topology equivalent to the standard topology on $\tilde{\mathbb{V}}_0^n$. From the Riesz representation theorem and (3.27)-(3.31) we have the following result.

Lemma 3.8. *For each $t \in [0, t_f)$, there exists a unique function $\mathbf{u} \in \tilde{\mathbb{V}}_0^n$ with $\dot{\mathbf{u}} \in \tilde{\mathbb{V}}_0^n$ such that (3.26) holds for all $\Phi \in \tilde{\mathbb{V}}_0^n$. Furthermore,*

$$\begin{aligned}
 \|\mathbf{u}\|_{\mathbb{V}^n} &\leq \frac{\sqrt{n}}{c_E} |\mathbf{b}_1| \|\mathbf{V}\|_{\mathbb{V}^n} + \frac{\hat{C}}{c_E} \|\hat{\sigma}\|_{L^2(\Gamma)^{n \times n}} \\
 \|\dot{\mathbf{u}}\|_{\mathbb{V}^n} &\leq \frac{\sqrt{n}}{c_E} |\mathbf{b}_1| \|\dot{\mathbf{V}}\|_{\mathbb{H}^n} + \frac{\hat{C}}{c_E} \|\dot{\hat{\sigma}}\|_{L^2(\Gamma)^{n \times n}}
 \end{aligned}$$

As a consequence, we obtain the Hadamard well-posedness in the weak sense of the Navier-type elastic system (2.8), (2.9), and (2.4):

Corollary 3.9 (Well-posedness of the elastic system). *Assume $\mathbf{V} \in L^2(0, t_f; \mathbb{V}^3)$ with $\dot{\mathbf{V}} \in L^2(0, t_f; \mathbb{H}^3)$ and $\hat{\sigma} \in L^2(0, t_f; L^2(\Gamma)^{n \times n})$ with $\dot{\hat{\sigma}} \in L^2(0, t_f; L^2(\Gamma)^{n \times n})$. Then (2.8), (2.9), and (2.4) admits a unique weak solution $\mathbf{u} \in L^2(0, t_f; \tilde{\mathbb{V}}_0^n)$ with $\dot{\mathbf{u}} \in L^2(0, t_f; \tilde{\mathbb{V}}_0^n)$ and this solution depends continuously on the data. That is, the mapping*

$$\mathbf{V}, \dot{\mathbf{V}}, \hat{\sigma}, \dot{\hat{\sigma}} \mapsto \mathbf{u}, \dot{\mathbf{u}}$$

from $L^2(0, t_f; \mathbb{V}^3) \times L^2(0, t_f; \mathbb{H}^3) \times L^2(0, t_f; L^2(\Gamma)^{n \times n}) \times L^2(0, t_f; L^2(\Gamma)^{n \times n})$ to $L^2(0, t_f; \tilde{\mathbb{V}}_0^n) \times L^2(0, t_f; \tilde{\mathbb{V}}_0^n)$ is continuous.

4. MAIN RESULTS

In this section, we present the main result of this paper, Theorem 4.1, and an example of its application. Theorem 4.1 gives a sufficient condition for the well-posedness in the weak sense of the coupled parabolic-elliptic initial-boundary value problem (2.1)-(2.7) describing the fully coupled TCP model. This condition depends on physical parameters of the system, coupling vectors, and the Korn constant which in turn depends only on the shape of a domain. Evaluation of the Korn constant becomes necessary for each specific domain and is essential for the proposed well-posedness theory. For instance, in the case of a two-dimensional annular region that does not experience pure rotation, the Korn constant can be expressed explicitly in terms of the ratio of the inner and outer radii of the domain. As a consequence, for such regions, the sufficient condition for the well-posedness of the problem can be formulated in terms of physical parameters, coupling vectors, and the ratio of the inner and outer radii. This result is presented in Corollary 4.2.

The sufficient conditions presented in Theorem 4.1 and Corollary 4.2 also provide compatibility conditions on the coupling vectors.

Theorem 4.1. *Under Assumptions 3.1, 3.2, and 3.7, the coupled TCP system (2.1)-(2.7) admits a unique weak solution*

$$\begin{aligned}
 (\mathbf{V}, \mathbf{u}) &\in L^2(0, t_f; \mathbb{V}^3) \times L^2(0, t_f; \tilde{\mathbb{V}}_0^n) \\
 \text{with } (\dot{\mathbf{V}}, \dot{\mathbf{u}}) &\in L^2(0, t_f; \mathbb{H}^3) \times L^2(0, t_f; \tilde{\mathbb{V}}_0^n)
 \end{aligned}
 \tag{4.1}$$

if

$$c_E - 2\left(2n + \frac{3n}{\lambda_{\min}^2}\right)^{1/2} \|P\| |P^{-1}M^{-1}\mathbf{b}_0| |\mathbf{b}_1| > 0
 \tag{4.2}$$

where $c_E = c_E(\Omega) > 0$ is the Korn constant, n is the dimension of the problem, $M^{-1}A = PDP^{-1}$, $D = \text{diag}(\lambda_1, \lambda_2, \lambda_3)$, and \mathbf{b}_0 and \mathbf{b}_1 are coupling vectors. Moreover, this solution depends continuously on the boundary data \mathbf{V}_B and $\dot{\mathbf{V}}_B$, the initial data \mathbf{V}_I and $\mathbf{V}_B(0)$, the applied boundary stress $\hat{\boldsymbol{\sigma}}$, and its time derivative $\dot{\hat{\boldsymbol{\sigma}}}$.

Proof. Property (4.1) of the solution (\mathbf{V}, \mathbf{u}) and the continuous dependence of the solution on the data follow immediately from Theorem 3.5, Lemma 3.6, and Corollary 3.9. We only need to prove that condition (4.2) is sufficient for well-posedness of the system.

Since, for each $t \in [0, t_f)$, $\dot{\mathbf{u}} \in \tilde{\mathbb{V}}_0^n$, from (3.26), the trace theorem, and the Poincaré inequality we have

$$\begin{aligned}
 a_E(\dot{\mathbf{u}}, \dot{\mathbf{u}}) &= \int_{\Omega} (\mathbf{b}_1 \cdot \dot{\mathbf{V}}) \nabla \cdot \dot{\mathbf{u}} d\Omega + \int_{\Gamma \setminus \Gamma_F} (\dot{\hat{\boldsymbol{\sigma}}}\mathbf{n}) \cdot \dot{\mathbf{u}} d\Gamma \\
 &\leq \|\mathbf{b}_1 \cdot \dot{\mathbf{V}}\|_{L^2(\Omega)} \|\nabla \cdot \dot{\mathbf{u}}\|_{L^2(\Omega)} + \|\dot{\hat{\boldsymbol{\sigma}}}\mathbf{n}\|_{L^2(\Gamma)^n} \|\dot{\mathbf{u}}\|_{L^2(\Gamma)^n} \\
 &\leq \sqrt{n} |\mathbf{b}_1| \|\dot{\mathbf{V}}\|_{\mathbb{H}^3} \|\nabla \dot{\mathbf{u}}\|_{\mathbb{H}^n} + \hat{C} \|\dot{\hat{\boldsymbol{\sigma}}}\|_{L^2(\Gamma)^{n \times n}} \|\nabla \dot{\mathbf{u}}\|_{\mathbb{H}^n}
 \end{aligned}$$

where $\hat{C} > 0$ is a constant. Utilizing (3.31), the last inequality yields

$$c_E \|\nabla \dot{\mathbf{u}}\|_{\mathbb{H}^n} \leq \sqrt{n} |\mathbf{b}_1| \|\dot{\mathbf{V}}\|_{\mathbb{H}^3} + \hat{C} \|\dot{\hat{\boldsymbol{\sigma}}}\|_{L^2(\Gamma)^{n \times n}}, \quad \forall t \in [0, t_f)$$

and therefore,

$$c_E \|\nabla \dot{\mathbf{u}}\|_{L^2(0, t_f; \mathbb{H}^n)} \leq \sqrt{2n} |\mathbf{b}_1| \|\dot{\mathbf{V}}\|_{L^2(0, t_f; \mathbb{H}^3)} + \sqrt{2}\hat{C} \|\dot{\hat{\boldsymbol{\sigma}}}\|_{L^2(0, t_f; L^2(\Gamma)^{n \times n})}
 \tag{4.3}$$

Substituting the a priori estimate (3.23) into (4.3), we have

$$\begin{aligned}
 &c_E \|\nabla \dot{\mathbf{u}}\|_{L^2(0, t_f; \mathbb{H}^n)} \\
 &\leq 2\left(2n + \frac{3n}{\lambda_{\min}^2}\right)^{1/2} \|P\| |P^{-1}M^{-1}\mathbf{b}_0| |\mathbf{b}_1| \|\nabla \dot{\mathbf{u}}\|_{L^2(0, t_f; \mathbb{H}^n)} + \Theta
 \end{aligned}
 \tag{4.4}$$

where $\Theta > 0$ is a constant independent of \mathbf{u} and $\dot{\mathbf{u}}$. Applying the Hadamard well-posedness condition for operator equations [13] to (4.4) yields the sufficient condition (4.2). \square

Corollary 4.2. *For an annular domain $\Omega \subset \mathbb{R}^2$ of inner radius R_B and outer radius R_F with $\delta = \frac{R_B}{R_F} \leq e^{-1}$, under the side condition on the displacement $\mathbf{u} = u\mathbf{i} + v\mathbf{j}$,*

$$\int_{\Omega} \left(\frac{\partial u}{\partial y} - \frac{\partial v}{\partial x}\right) d\Omega = 0
 \tag{4.5}$$

the coupled TCP system (2.1)-(2.7) has a unique weak solution if

$$\frac{G}{2} \left[1 - \left(\frac{3}{\delta^{-2} + 1 + \delta^2} \right)^{1/2} \right] - 2 \left(4 + \frac{6}{\lambda_{\min}^2} \right)^{1/2} \|P\| \|P^{-1} M^{-1} \mathbf{b}_0\| |\mathbf{b}_1| > 0 \quad (4.6)$$

where G is the shear modulus, $M^{-1}A = PDP^{-1}$, $D = \text{diag}(\lambda_1, \lambda_2, \lambda_3)$, and \mathbf{b}_0 and \mathbf{b}_1 are coupling vectors.

Proof. Under the given conditions, the following inequality holds [12]:

$$\frac{D(\dot{\mathbf{u}})}{S(\dot{\mathbf{u}})} \leq 4 \left[1 - \left(\frac{3}{\delta^{-2} + 1 + \delta^2} \right)^{1/2} \right]^{-1}$$

where $D(\dot{\mathbf{u}}) = \|\nabla \dot{\mathbf{u}}\|_{\mathbb{H}^2}^2$ and $S(\dot{\mathbf{u}}) = \int_{\Omega} \sum_{i,j=1}^2 \varepsilon_{ij}(\dot{\mathbf{u}})^2 d\Omega$. Therefore,

$$S(\dot{\mathbf{u}}) \geq \frac{1}{4} \left[1 - \left(\frac{3}{\delta^{-2} + 1 + \delta^2} \right)^{1/2} \right] \|\nabla \dot{\mathbf{u}}\|_{\mathbb{H}^2}^2 \quad (4.7)$$

For homogeneous and isotropic medium, the coefficients of elasticity a_{ijkl} , $i, j, k, l = 1, 2$, in terms of the Lamé parameters λ and μ are $a_{1111} = a_{2222} = 2\mu + \lambda$, $a_{1122} = a_{2211} = \lambda$, $a_{1212} = a_{2121} = a_{1221} = a_{2112} = \mu$ and, from (2.11) and (2.12), we have

$$\begin{aligned} a_E(\dot{\mathbf{u}}, \dot{\mathbf{u}}) &= \int_{\Omega} 2\mu(\varepsilon_{11}^2 + \varepsilon_{12}^2 + \varepsilon_{21}^2 + \varepsilon_{22}^2) + \lambda(\varepsilon_{11} + \varepsilon_{22})^2 d\Omega \\ &\geq 2\mu \int_{\Omega} \varepsilon_{11}^2 + \varepsilon_{12}^2 + \varepsilon_{21}^2 + \varepsilon_{22}^2 d\Omega \\ &= 2\mu S(\dot{\mathbf{u}}) \end{aligned} \quad (4.8)$$

where $\mu = G$ is the shear modulus. Equations (3.31), (4.7), and (4.8) give the Korn constant

$$c_E = \frac{G}{2} \left[1 - \left(\frac{3}{\delta^{-2} + 1 + \delta^2} \right)^{1/2} \right]$$

Applying this constant to the sufficient condition (4.2) yields (4.6). \square

Remark 4.3. The purpose of the condition (4.5) is to eliminate a pure rotation.

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TETYANA MALYSHEVA

DEPARTMENT OF NATURAL & APPLIED SCIENCES, UNIVERSITY OF WISCONSIN-GREEN BAY, GREEN BAY, WI 54311-7001, USA

E-mail address: malyshet@uwgb.edu

LUTHER W. WHITE

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF OKLAHOMA, NORMAN, OK 73019-3103, USA

E-mail address: lwhite@ou.edu