

LOWER BOUNDS FOR THE BLOWUP TIME OF SOLUTIONS TO A NONLINEAR PARABOLIC PROBLEM

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ABSTRACT. In this short article, we study the blow-up properties of solutions to a parabolic problem with a gradient nonlinearity under homogeneous Dirichlet boundary conditions. By constructing an auxiliary function and by modifying the first order differential inequality technique introduced by Payne et al., we obtain a lower bound for the blow-up time of solutions in a bounded domain $\Omega \subset \mathbb{R}^n$ for any $n \geq 3$. This article generalizes a result in [16].

1. INTRODUCTION

When dealing with a parabolic problem there are several interesting features to analyze, one of which is the so called finite time blow-up. The question of blow-up of solutions to nonlinear parabolic equations and systems has received considerable attention since the elegant work of Fujita [6]. We refer to the interested readers the survey papers [2, 7, 10] and the book [17].

In practical situations, one would like to know, among other things, whether the solutions blow up, and if so, at what time T blow-up occurs. However, when the solution does blow up at some finite T , this time can seldom be determined explicitly, and much effort has been devoted to the calculation of bounds for T . Most of the methods used until recently can only yield upper bounds for T , which are of little value in particular situations when blow-up has to be avoided. By using the first-order differential inequality technique, lower bounds for the blow-up time of solutions to semilinear heat equations under different boundary conditions and suitable constraint on the data were obtained by Payne et al. [12, 13, 14, 16]. Thereafter, the differential inequality technique was successfully employed to derive lower bounds for the blow-up time of solutions to other parabolic problems, see [1, 3, 5, 11, 15].

In this article, we shall study a parabolic problem with a gradient nonlinearity of the following form

$$\begin{aligned} u_t &= \Delta u + u^p - |\nabla u|^q, & (x, t) &\in \Omega \times (0, T), \\ u(x, t) &= 0, & (x, t) &\in \partial\Omega \times (0, T), \\ u(x, 0) &= u_0(x) \geq 0, & x &\in \Omega, \end{aligned} \tag{1.1}$$

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where Ω is a bounded domain in \mathbb{R}^n with smooth boundary $\partial\Omega$, Δ and ∇ are the Laplace and gradient operator with respect to x , respectively, T is the possible blow-up time and $p, q > 1$ are fixed (finite) parameters. In [4, 9], conditions on p, q and $u_0(x)$ were given for which the solutions to (1.1) would blow up in finite time. In fact the restrictions on p and q were

$$1 < p < \frac{n+2}{n-2}, \quad 1 < q < \frac{2p}{p+1}, \quad \text{for } n \geq 2,$$

or

$$p \text{ is large enough and } q = \frac{2p}{p+1}, \text{ for } n = 1.$$

In a recent paper Payne et al. [16] obtained lower bounds of the blow-up time of solutions to (1.1) when $n = 3$. Naturally, we hope to obtain the lower bounds for blow-up time of solutions to (1.1) with any smooth bounds $\Omega \subset \mathbb{R}^n$ and any $n \geq 3$. That is what we will do in this article.

As indicated in [18] it is well known that if $p \leq q$ the solution will not blow up in finite time. Also it is well known that if the initial data are small enough the solution will actually decay exponentially as $t \rightarrow \infty$ (see e.g.[14, 19]). Since we are interested in a lower bound for the blow-up time T , only the case $p > q$ is considered.

2. A LOWER BOUND FOR THE BLOW-UP TIME

In this section we seek a lower bound for the blow-up time T of solutions to (1.1) in some appropriate measure. The idea of the proof of the following theorem is inspired by that in [1].

Theorem 2.1. *Let $u(x, t)$ be the nonnegative classical solution of problem (1.1) for $p > q > 1$ in a smooth bounded domain $\Omega \subset \mathbb{R}^n$ with $n \geq 3$. Define*

$$\varphi(t) = \int_{\Omega} u^k dx,$$

where k is a parameter restricted by the condition

$$k > \max \left\{ 1, \frac{(7n-16)(p-1)}{2}, (q-1)(3n-8) \right\}. \quad (2.1)$$

If $u(x, t)$ blows up in the measure φ at the finite time T , then T is bounded from below as

$$T \geq \int_{\varphi(0)}^{+\infty} \frac{1}{C_1 + C_2 \xi^{\frac{3n-6}{3n-8}}} d\xi, \quad (2.2)$$

where C_1 and C_2 are positive constants which will be determined in the proof.

Proof. Applying the divergence theorem to the first equation in (1.1), we have

$$\begin{aligned}
 \frac{d\varphi}{dt} &= k \int_{\Omega} u^{k-1} u_t dx \\
 &= k \int_{\Omega} u^{k-1} (\Delta u + u^p - |\nabla u|^q) dx \\
 &= k \int_{\Omega} u^{k-1} \Delta u dx + k \int_{\Omega} u^{k+p-1} dx - k \int_{\Omega} u^{k-1} |\nabla u|^2 dx \quad (2.3) \\
 &= -\frac{4(k-1)}{k} \int_{\Omega} |\nabla u^{k/2}|^2 dx + k \int_{\Omega} u^{k+p-1} dx \\
 &\quad - \frac{kq^q}{(k+q-1)^q} \int_{\Omega} |\nabla u^{\frac{k+q-1}{q}}|^q dx.
 \end{aligned}$$

Moreover, from [12, (2.10)] it follows that

$$\int_{\Omega} |\nabla u^{\frac{k+q-1}{q}}|^q dx \geq \left(\frac{2\sqrt{\lambda}}{q}\right)^q \int_{\Omega} u^{k+q-1} dx, \quad (2.4)$$

where the positive constant λ is the first eigenvalue of the problem

$$\begin{aligned}
 \Delta w + \lambda w &= 0 \quad \text{in } \Omega, \\
 w &= 0 \quad \text{on } \partial\Omega.
 \end{aligned} \quad (2.5)$$

Thus by combining (2.3) with (2.4) we obtain

$$\frac{d\varphi}{dt} \leq -\frac{4(k-1)}{k} \int_{\Omega} |\nabla u^{k/2}|^2 dx + k \int_{\Omega} u^{k+p-1} dx - \frac{k(2\sqrt{\lambda})^q}{(k+q-1)^q} \int_{\Omega} u^{k+q-1} dx. \quad (2.6)$$

Noticing (2.1), we can apply first Hölder's inequality and then Young's inequality to the second term on the right hand side of (2.3) to obtain

$$\begin{aligned}
 \int_{\Omega} u^{k+p-1} dx &\leq |\Omega|^{m_1} \left(\int_{\Omega} u^{\frac{k(7n-14)}{7n-16}} dx \right)^{m_2} \\
 &\leq m_1 |\Omega| + m_2 \int_{\Omega} u^{\frac{k(7n-14)}{7n-16}} dx,
 \end{aligned} \quad (2.7)$$

where

$$m_1 = 1 - \frac{(k+p-1)(7n-16)}{k(7n-14)} \in (0, 1), \quad m_2 = \frac{(k+p-1)(7n-16)}{k(7n-14)} \in (0, 1).$$

Combining (2.7) and (2.6) yields

$$\begin{aligned}
 \frac{d\varphi}{dt} &\leq -\frac{4(k-1)}{k} \int_{\Omega} |\nabla u^{k/2}|^2 dx + km_1 |\Omega| + km_2 \int_{\Omega} u^{\frac{k(7n-14)}{7n-16}} dx \\
 &\quad - \frac{k(2\sqrt{\lambda})^q}{(k+q-1)^q} \int_{\Omega} u^{k+p-1} dx.
 \end{aligned} \quad (2.8)$$

We now use Hölder's inequality in the third term on the right hand side of (2.8):

$$\int_{\Omega} u^{\frac{k(7n-14)}{7n-16}} dx \leq \left(\int_{\Omega} u^k dx \right)^{\alpha} \left(\int_{\Omega} u^{\frac{k}{2} \frac{2n}{n-2}} dx \right)^{1-\alpha}, \quad (2.9)$$

where $0 < \alpha = \frac{2(3n-7)}{7n-16} < 1$. Next, using the Sobolev inequality for $W_0^{1,2} \hookrightarrow L^{\frac{2n}{n-2}}$ ($n \geq 3$) [20]), we obtain

$$\|u^{k/2}\|_{L^{\frac{2n}{n-2}}}^{\frac{2n(1-\alpha)}{n-2}} \leq C_s^{\frac{2n(1-\alpha)}{n-2}} \|\nabla u^{k/2}\|_{L^2}^{\frac{2n(1-\alpha)}{n-2}}, \quad (2.10)$$

where $C_s = \left(\frac{1}{n(n-2)\pi}\right)^{1/2} \left(\frac{n!}{2\Gamma(\frac{n}{2}+1)}\right)^{1/n}$ is the best imbedding constant (see [8, Chap. 7]). By substituting (2.10) into (2.9), we arrive at

$$\int_{\Omega} u^{\frac{k(7n-14)}{7n-16}} dx \leq C_s^{\frac{2n(1-\alpha)}{n-2}} \left(\int_{\Omega} u^k dx\right)^{\alpha} \left(\int_{\Omega} |\nabla u^{k/2}|^2 dx\right)^{\frac{n(1-\alpha)}{n-2}}, \quad (2.11)$$

which, with the help of Young's inequality, gives

$$\int_{\Omega} u^{\frac{k(7n-14)}{7n-16}} dx \leq \frac{C_s^{\frac{n}{3n-8}} (6n-16)}{(7n-16)\varepsilon_1^{\frac{2}{2(3n-8)}}} \left(\int_{\Omega} u^k dx\right)^{\frac{3n-7}{3n-8}} + \frac{n(1-\alpha)\varepsilon_1}{n-2} \int_{\Omega} |\nabla u^{k/2}|^2 dx. \quad (2.12)$$

Here ε_1 is a positive constant to be determined later. By Hölder's inequality, we have

$$\int_{\Omega} u^{q+k-1} dx \geq |\Omega|^{-\frac{q-1}{k}} \left(\int_{\Omega} u^k dx\right)^{1+\frac{q-1}{k}}. \quad (2.13)$$

Combining (2.12) and (2.13) with (2.8) gives

$$\begin{aligned} \frac{d\varphi}{dt} &\leq km_1|\Omega| + \left[\frac{n(1-\alpha)\varepsilon_1 km_2}{n-2} - \frac{4(k-1)}{k}\right] \int_{\Omega} |\nabla u^{k/2}|^2 dx \\ &+ \frac{km_2 C_s^{\frac{n}{3n-8}} (6n-16)}{(7n-16)\varepsilon_1^{\frac{2}{2(3n-8)}}} \varphi^{\frac{3n-7}{3n-8}} - \frac{k(2\sqrt{\lambda})^q}{(k+q-1)^q} |\Omega|^{-\frac{q-1}{k}} \varphi^{1+\frac{q-1}{k}}. \end{aligned} \quad (2.14)$$

Next, we apply Young's inequality to the third term on the right-hand side of (2.14) to conclude that

$$\varphi^{\frac{3n-7}{3n-8}} \leq \frac{\varepsilon_2}{m_3} \varphi^{1+\frac{q-1}{k}} + \frac{1}{m_4} \varepsilon_2^{-\frac{m_4}{m_3}} \varphi^{\frac{3n-6}{3n-8}}, \quad (2.15)$$

where

$$m_3 = \frac{2k - (q-1)(3n-8)}{k}, \quad m_4 = \frac{2k - (q-1)(3n-8)}{k - (q-1)(3n-8)},$$

and ε_2 is a positive constant to be fixed. Combining (2.15) and (2.14), we obtain

$$\begin{aligned} \frac{d\varphi}{dt} &\leq C_1 + \left[\frac{n(1-\alpha)\varepsilon_1 km_2}{n-2} - \frac{4(k-1)}{k}\right] \int_{\Omega} |\nabla u^{k/2}|^2 dx + C_2 \varphi^{\frac{3n-6}{3n-8}} \\ &+ \left[\frac{\varepsilon_2 km_2 C_s^{\frac{n}{3n-8}} (6n-16)}{(7n-16)\varepsilon_1^{\frac{2}{2(3n-8)}} m_3} - \frac{k(2\sqrt{\lambda})^q |\Omega|^{-\frac{q-1}{k}}}{(k+q-1)^q}\right] \varphi^{1+\frac{q-1}{k}}, \end{aligned} \quad (2.16)$$

where

$$C_1 = km_1|\Omega|, \quad C_2 = \frac{km_2 C_s^{\frac{n}{3n-8}} (6n-16)\varepsilon_2^{-\frac{m_4}{m_3}}}{(7n-16)\varepsilon_1^{\frac{2}{2(3n-8)}} m_4}.$$

Therefore, by choosing

$$\varepsilon_1 = \frac{4(k-1)(n-2)}{nk^2 m_2 (1-\alpha)}$$

first and

$$\varepsilon_2 = \frac{(7n-16)m_3k(2\sqrt{\lambda})^q|\Omega|^{-\frac{q-1}{k}}\varepsilon_1^{\frac{n}{2(3n-8)}}}{km_2(6n-16)C_s^{\frac{n}{3n-8}}(k+q-1)^q}$$

next, we obtain the differential inequality

$$\frac{d\varphi}{dt} \leq C_1 + C_2\varphi^{\frac{3n-6}{3n-8}}, \quad (2.17)$$

or equivalently

$$\frac{d\varphi}{C_1 + C_2\varphi^{\frac{3n-6}{3n-8}}} \leq dt. \quad (2.18)$$

Integrating of the differential inequality (2.18) from 0 to t leads to

$$\int_{\varphi(0)}^{\varphi(t)} \frac{1}{C_1 + C_2\xi^{\frac{3n-6}{3n-8}}} d\xi \leq t. \quad (2.19)$$

Passing to the limit as $t \rightarrow T^-$, we obtain

$$\int_{\varphi(0)}^{+\infty} \frac{1}{C_1 + C_2\xi^{\frac{3n-6}{3n-8}}} d\xi \leq T. \quad (2.20)$$

Thus, the proof is complete. \square

Remark 2.2. It is easy to see that when $n = 3$, the lower bound for the blow-up time derived here is consistent with the one obtained by Payne et al. [16].

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