

## BIFURCATIONS FOR A PHYTOPLANKTON MODEL WITH TIME DELAY

CHANGJIN XU

ABSTRACT. Applying a frequency domain approach, we investigate a phytoplankton model with time delay. We use the delay as a bifurcation parameter; as it passes through a sequence of critical values, Hopf bifurcation occurs. A family of periodic solutions bifurcate from the equilibrium when the bifurcation parameter exceeds a critical value. Some numerical simulations illustrate our theoretical results.

### 1. INTRODUCTION

Phytoplankton and Zooplankton are a single celled organisms that drift with the currents on the surface of open oceans. They play an important role in stabilizing the environment; for example, they are the staple item for the food web and they are recyclers of most of energy that flows through the ocean ecosystem. Phytoplankton systems have received much attention from biologists and mathematicians [2, 4]. In 2010, Dhar and Sharma [3] investigated the stability of the phytoplankton system

$$\begin{aligned}\frac{dP_s(t)}{dt} &= rP_s\left[1 - \frac{P_s}{K}\right] - \alpha P_s P_i + \gamma P_i, \\ \frac{dP_i(t)}{dt} &= \alpha P_s P_i - \beta P_i,\end{aligned}\tag{1.1}$$

where  $P_s, P_i$  are the population densities of susceptible and infected phytoplankton at any instant of time  $t$ .  $r$  is the intrinsic growth rate of the population of susceptible phytoplankton,  $K$  is the carrying capacity of the population of susceptible phytoplankton,  $\alpha$  is the disease contact rate of the disease phytoplankton population,  $\beta$  is the removal rate of the disease phytoplankton population, out of which  $\gamma$  fraction of infected phytoplankton rejoin the susceptible phytoplankton population. In details, one can see [3]. Taking into account that there is a certain time delay during the process of reproduction (for example, egg formation will take  $\tau$  units of time before hatching), the dynamic behavior of the system not only is affected by the current state of the system, but also the past state of the system, i.e., there

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exists inherent lag in the system. Based on this view of point, in this paper, we will revise system (1.1) as follows:

$$\begin{aligned}\frac{dP_s(t)}{dt} &= rP_s\left[1 - \frac{P_s(t-\tau)}{K}\right] - \alpha P_s P_i + \gamma P_i, \\ \frac{dP_i(t)}{dt} &= \alpha P_s P_i - \beta P_i.\end{aligned}\tag{1.2}$$

In this article, we investigate Hopf bifurcation for system (1.2). It is worth pointing out that many early work on Hopf bifurcation of the delayed differential equations is used the state-space formulation for delayed differential equations, known as the “time domain” approach. But there exists another approach that comes from the theory of feedback systems known as frequency domain method which was initiated and developed by Allwright [1], Mees and Chua [7] and Muiola and Chen [8, 9] and is familiar to control engineers. This alternative representation applies the engineering feedback systems theory and methodology: an approach described in the “frequency domain”—the complex domain after the standard Laplace transforms having been taken on the state-space system in the time domain. This new methodology has some advantages over the classical time-domain methods [5, 6]. A typical one is its pictorial characteristic that utilizes advanced computer graphical capabilities thereby bypassing quite a lot of profound and difficult mathematical analysis.

In this paper, we will devote our attention to finding the Hopf bifurcation point for models (1.2) by means of the frequency-domain approach. We found that if the coefficient  $\tau$  is used as a bifurcation parameter, then Hopf bifurcation occurs for the model (1.2). That is, a family of periodic solutions bifurcates from the equilibrium when the bifurcation parameter exceeds a critical value. Some numerical simulations are carried out to illustrate the theoretical analysis. We believe that it is the first time to investigate Hopf bifurcation of the model (1.2) using the frequency-domain approach.

The remainder of the paper is organized as follows: in Section 2, applying the frequency-domain approach formulated by Muiola and Chen [9], the existence of Hopf bifurcation parameter is determined and shown that Hopf bifurcation occurs when the bifurcation parameter exceeds a critical value. In Section 3, some numerical simulation are carried out to verify the correctness of theoretical analysis result. Finally, some conclusions and discussions are included in Section 4.

## 2. STABILITY OF THE EQUILIBRIUM AND LOCAL HOPF BIFURCATIONS

In model (1.2), we assume that the condition

$$(H1) \quad (K\alpha - \beta)(\beta - \gamma) > 0.$$

It is obvious that system (1.2) has a unique positive equilibrium  $E_*(P_s^*, P_i^*)$ , where,

$$P_s^* = \frac{\beta}{\alpha}, \quad P_i^* = \frac{r\beta(K\alpha - \beta)}{K\alpha^2(\beta - \gamma)}.$$

We can rewrite the nonlinear system (1.2) as a matrix form

$$\frac{dx(t)}{dt} = Ax(t) + H(x),\tag{2.1}$$

where  $x = (P_s(t), P_i(t))^T$ ,

$$A = \begin{pmatrix} r & 1 \\ 0 & -\beta \end{pmatrix}, \quad H(x) = \begin{pmatrix} -\frac{rP_sP_s(t-\tau)}{K} - \alpha P_sP_i \\ \alpha P_sP_i \end{pmatrix}.$$

Choosing the coefficient  $\tau$  as a bifurcation and introducing a “state-feedback control”  $u = g[y(t - \tau); \tau]$ , where  $y(t) = (y_1(t), y_2(t))^T$ , we obtain a linear system with a non-linear feedback as follows

$$\begin{aligned} \frac{dx}{dt} &= Ax + Bu, \\ y &= -Cx, \\ u &= g[y(t - \tau); \tau], \end{aligned} \tag{2.2}$$

where

$$B = C = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \quad u = g[y(t - \tau), \tau] = \begin{pmatrix} -\frac{ry_1y_1(t-\tau)}{K} - \alpha y_1y_2 \\ \alpha y_1y_2 \end{pmatrix}.$$

Next, taking Laplace transform on (2.2), we obtain the standard transfer matrix of the linear part of the system:

$$G(s; \tau) = C[sI - A]^{-1}B.$$

Then

$$G(s; \tau) = \begin{pmatrix} \frac{1}{s-r} & \frac{r}{(s-r)(s+\beta)} \\ 0 & \frac{1}{s+\beta} \end{pmatrix}. \tag{2.3}$$

If this feedback system is linearized about the equilibrium  $y = -C(P_s^*, P_i^*)^T$ , then the Jacobian of (2.3) is

$$J(\tau) = \frac{\partial g}{\partial y} \Big|_{y=\bar{y}=-C(P_s^*, P_i^*)^T} = \begin{pmatrix} \frac{r(P_s^* + P_s^* e^{-s\tau})}{K} + \alpha x_2^* & \alpha x_1^* \\ -\alpha x_2^* & -\alpha x_1^* \end{pmatrix}.$$

Let

$$\begin{aligned} h(\lambda, s; \tau) &= \det |\lambda I - G(s; \tau)J(\tau)| \\ &= \lambda^2 + \left[ \frac{\alpha P_s^*}{s+\beta} + \frac{r\alpha P_i^*}{(s-r)(s+\beta)} - \frac{1}{s-r} \left( \frac{rP_s^*(1+e^{-s\tau})}{K} + \alpha P_i^* \right) \right] \lambda \\ &\quad + \frac{\alpha P_i^*}{s+\beta} \left[ \frac{\alpha P_s^*}{s-r} - \frac{r\alpha P_s^*}{(s-r)(s+\beta)} \right] = 0. \end{aligned}$$

Applying the generalized Nyquist stability criterion with  $s = i\omega$ , we obtain the following results.

**Lemma 2.1** (Moiola, Chen, 1996). *If an eigenvalue of the corresponding Jacobian of the nonlinear system, in the time domain, assumes a purely imaginary value  $i\omega_0$  at a particular  $\tau = \tau_0$ , then the corresponding eigenvalue of the constant matrix  $G(i\omega_0; \tau_0)J(\tau_0)$  in the frequency domain must assume the value  $-1 + i0$  at  $\tau = \tau_0$ .*

To apply Lemma 2.1, let  $\hat{\lambda} = \hat{\lambda}(i\omega; \tau)$  be the eigenvalue of  $G(i\omega; \tau)J(\tau)$  that satisfies  $\hat{\lambda}(i\omega_0; \tau_0) = -1 + 0i$ . Then

$$h(-1, i\omega_0; \tau_0) = 0.$$

Separating the real and imaginary parts and rearranging, we obtain

$$E \cos \omega_0 \tau_0 + F \sin \omega_0 \tau_0 = S, \tag{2.4}$$

$$F \cos \omega_0 \tau_0 + E \sin \omega_0 \tau_0 = T, \quad (2.5)$$

where

$$\begin{aligned} E &= rP_s^*(\beta^2 - \omega_0^2), \quad F = -2r\beta\omega_0P_s^*, \\ S &= Kr(\omega_0^2 - \beta^2) - 2K\beta\omega_0^2 + K\alpha P_s^*(\beta r + \omega_0^2) \\ &\quad - rK\alpha\beta P_i^* - (\beta^2 - \omega_0^2)(rP_s^* + K\alpha P_i^*), \\ T &= Kr\alpha\omega_0P_i^* - K\omega_0(\beta^2 - \omega_0^2) + 2Kr\beta\omega_0 \\ &\quad + K\alpha P_s^*\omega_0(\beta - r) + 2\beta\omega_0(rP_s^* + \alpha K P_i^*). \end{aligned} \quad (2.6)$$

It follows from (2.4) and (2.5) that

$$E^2 + F^2 = S^2 + T^2. \quad (2.7)$$

Then

$$K^2\omega_0^6 + \theta_1\omega_0^4 + \theta_2\omega_0^3 + \theta_3\omega_0^2 + \theta_4 = 0, \quad (2.8)$$

where

$$\begin{aligned} \theta_1 &= (Kr - 2K\beta + K\alpha P_s^* + rP_s^* + K\alpha P_i^*)^2 - r^2(P_s^*)^2, \\ \theta_2 &= 2K[Kr\alpha P_i^* - K\beta^2 + 2Kr\beta + K\alpha P_s^*(\beta - r) + 2\beta(rP_s^* + K\alpha P_i^*)], \\ \theta_3 &= 2(Kr - 2K\beta + K\alpha P_s^* + rP_s^* + K\alpha P_i^*)(Kr\alpha\beta P_s^* - Kr\beta^2 - Kr\alpha\beta P_i^*), \\ &\quad - 2r^2\beta^2(P_s^*)^2, \\ \theta_4 &= (Kr\alpha\beta P_s^* - Kr\beta^2 - Kr\alpha\beta P_i^*)^2 + \theta_3^2 - r^2\beta^4(P_s^*)^2 + [Kr\alpha P_i^* - K\beta^2 \\ &\quad + 2Kr\beta + K\alpha P_s^*(\beta - r) + 2\beta(rP_s^* + K\alpha P_i^*)]^2 - r^2\beta^*(P_s^*)^2. \end{aligned} \quad (2.9)$$

It is easy to see that if  $\theta_4 < 0$ , then (2.8) has at least one positive root (say  $\omega_0$ ). By (2.8), we can compute the value of  $\omega_0$  by means of Matlab software. Then from (2.4) and (2.5), we obtain

$$\tau_0 = \frac{1}{\omega_0} [\arccos \frac{ES - FT}{E^2 - F^2} + 2k\pi] (k = 0, 1, 2, \dots). \quad (2.10)$$

In the sequel, we will consider the transversality condition for Hopf bifurcation of system (1.2).

In view of the definition of  $h(\lambda, s; \tau)$ , we have

$$\begin{aligned} \left[\frac{d\lambda}{d\tau}\right]^{-1} &= \frac{2K^3(s-r)\lambda + rP_s^*(1 + e^{-s\tau})(s + \beta) + K\alpha P_i^*(s + \beta)}{K^2rP_s^*(e^{-s\tau})\lambda(s + \beta)} \\ &\quad - \frac{K^3(s-r)\alpha P_s^* + Kr\alpha P_i^*}{K^2rP_s^*(e^{-s\tau})\lambda(s + \beta)}. \end{aligned}$$

Thus we obtain

$$\left[\frac{d\lambda}{d\tau}\right]^{-1} = \frac{C + iD}{A + iB},$$

which leads to

$$\operatorname{Re}\left[\frac{d\lambda}{d\tau}\right]^{-1}_{\tau=\tau_0} = \frac{AC + BD}{A^2 + B^2},$$

where

$$A = -K^2rP_s^*(\beta\omega_0 \sin \omega_0 \tau_0 - \omega_0^2 \cos \omega_0 \tau_0), \quad (2.11)$$

$$B = -K^2rP_s^*(\omega_0^2 \sin \omega_0 \tau_0 + \beta\omega_0 \cos \omega_0 \tau_0), \quad (2.12)$$

$$C = 2K^3(r\beta - \omega_0^2) - K^3r\alpha P_s^* + Kr\alpha P_i^* - K\alpha\beta P_i^* - rP_s^*[\beta(1 + \cos \omega_0\tau_0) + \omega_0 \sin \omega_0\tau_0], \quad (2.13)$$

$$D = 2K^3(\beta + r) + K^3\omega_0 - rP_s^*[\omega_0(1 + \cos \omega_0\tau_0) - \beta \sin \omega_0\tau_0] - K\omega_0\alpha P_i^*. \quad (2.14)$$

To obtain our main result, we assume

$$(H2) \quad AC + BD \neq 0.$$

**Theorem 2.2** (Existence of Hopf bifurcation parameter). *Let  $\theta_4, A, B, C, D$  be defined by (2.9), (2.11), (2.12), (2.13), (2.14), respectively. For (1.2), if  $\theta_4 < 0$  and conditions (H1)–(H2) hold, then Hopf bifurcation point of system (1.2) is*

$$\tau_0 = \frac{1}{\omega_0} \left[ \arccos \frac{ES - FT}{E^2 - F^2} + 2k\pi \right] (k = 0, 1, 2, \dots),$$

where  $\omega_0$  is positive real roots of (2.8), and  $E, F, S, T$  are defined by (2.6).

### 3. NUMERICAL EXAMPLES

In this section, we shall carry out numerical simulations for supporting our theoretical analysis. As an example, We consider system (1.2) with  $r = 0.5$ ,  $K = 2$ ,  $\alpha = 2$ ,  $\gamma = 0.4\beta = 2$ ; that is,

$$\begin{aligned} \frac{dP_s(t)}{dt} &= 0.5P_s \left[ 1 - \frac{P_s(t-\tau)}{2} \right] - 2P_sP_i + 0.4P_i, \\ \frac{dP_i(t)}{dt} &= 2P_sP_i - 2P_i. \end{aligned} \quad (3.1)$$

By Theorem 2.2, we obtain  $\tau_0 \approx 2.17$ . Numerical simulations for  $\tau = 2.1$  are shown in Figure 1. Thus we conclude that when  $\tau < \tau_0 \approx 2.17$ , system (3.1) is asymptotically stable. Numerical simulations for  $\tau = 2.4$  are shown in Figure 2. Thus we conclude that when  $\tau > \tau_0 \approx 2.17$ , system (3.1) undergoes a Hopf bifurcation that occurs near the positive equilibrium. Therefore  $\tau_0 \approx 2.17$  is a supercritical Hopf bifurcation point.

**Conclusions and discussions.** In this paper, we investigated a class of phytoplankton model with time delay. By choosing the coefficient  $\tau$  as a bifurcating parameter and analyzing the associating characteristic equation. It is found that a Hopf bifurcation occurs when the bifurcating parameter  $\tau$  passes through a critical value. Considering computational complexity, the direction and the stability of the bifurcating periodic orbits for system (1.2) have not been investigated. It is beyond the scope of the present paper and will be further investigated elsewhere in the near future.

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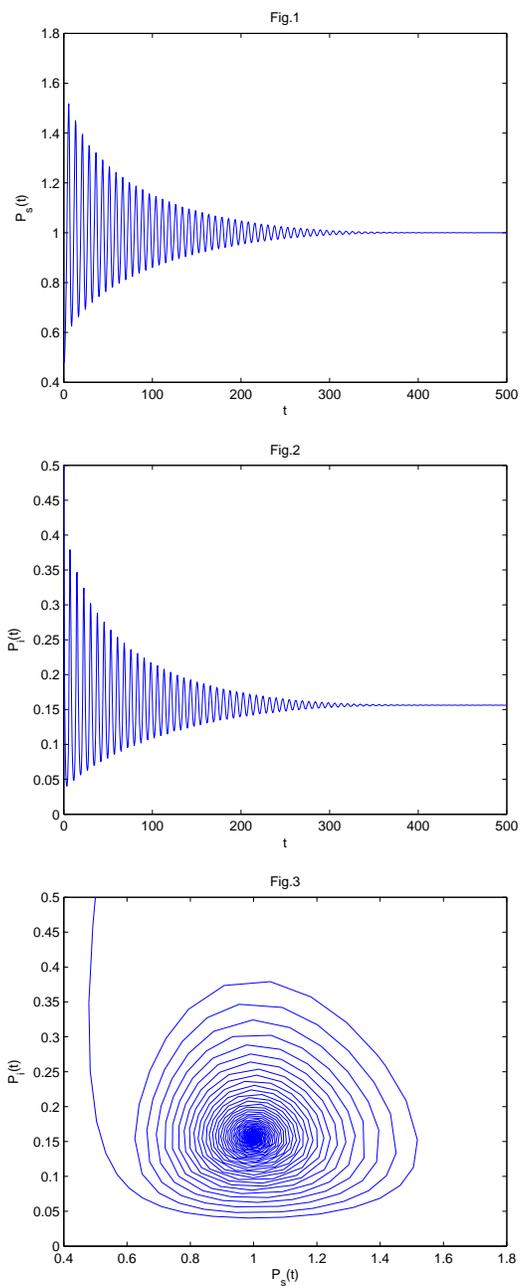


FIGURE 1. Trajectories and phase for (3.1) with  $\tau = 2.1 < \tau_0 \approx 2.17$ . The positive equilibrium is asymptotically stable. The initial value is  $(0.5, 0.5)$ .

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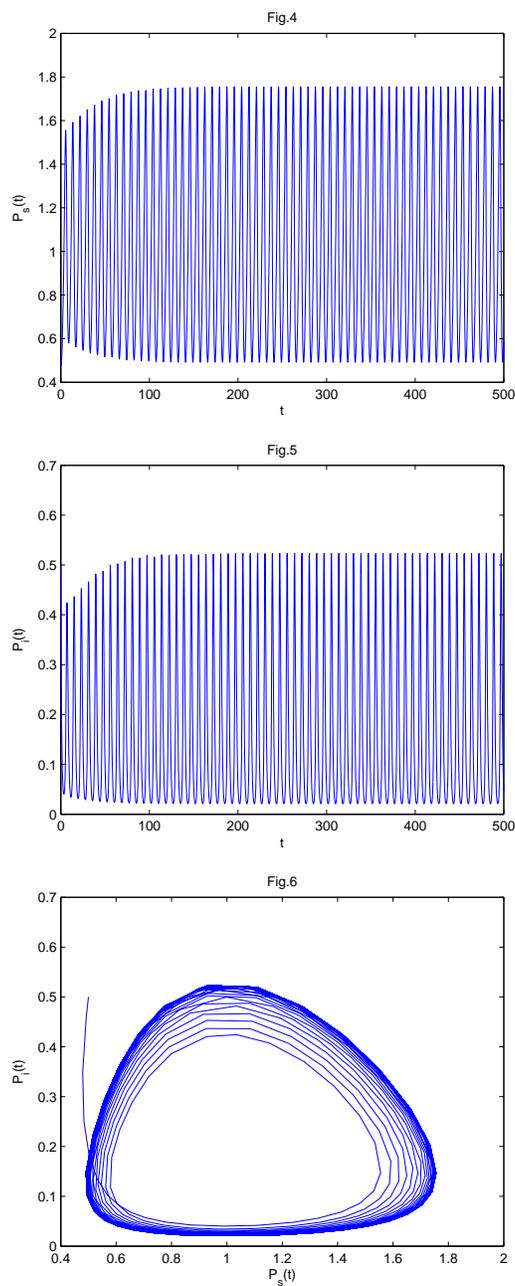


FIGURE 2. Trajectories and phase for (3.1) with  $\tau = 2.4 > \tau_0 \approx 2.17$ . Hopf bifurcation occurs from the positive equilibrium. The initial value is  $(0.5, 0.5)$

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CHANGJIN XU

GUIZHOU KEY LABORATORY OF ECONOMICS SYSTEM SIMULATION, SCHOOL OF MATHEMATICS AND STATISTICS, GUIZHOU COLLEGE OF FINANCE AND ECONOMICS, GUIYANG 550004, CHINA

*E-mail address:* `xcj403@126.com`